

# Arithmetic Principalization as a Legendre Transform: Classical Poisson Geometry, Deformation Quantization, and the Emergence of Proper Time, CP Violation, and Definite Outcomes in the Relativistic Field Theory of Primes (RFTP)

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## Abstract

We demonstrate that the infrared dynamics of the resonant spectral triple in the Bost–Connes system, deformed by irregular-prime shear, implements a precise arithmetic analogue of the classical Legendre transform. Starting from a classical Poisson manifold whose coordinates are the shear amplitudes  $\{I'_p\}$  and the bifundamental mixing field  $\Phi$ , we construct an explicit action functional whose Legendre transform yields a Hamiltonian on the cotangent bundle. The weight-12 winding condition is imposed as a classical first-class constraint, and Dirac reduction produces a reduced symplectic manifold  $P_{\text{red}}$ . Deformation quantization of this reduced space, driven by the modular derivation  $D_p$  acting asymmetrically through  $\Phi$  on the two complex-multiplication stabilizers ( $\mathbb{Z}/6\mathbb{Z}$  at  $\tau = \rho$  and  $\mathbb{Z}/4\mathbb{Z}$  at  $\tau = i$ ), generates a non-commutative star product whose leading commutator is non-vanishing and proportional to the stabilizer contrast transmitted by  $\Phi$ .

This framework naturally produces: (i) proper time  $\tau$  as the integrated thermodynamic work along the flow, with lapse  $N \propto \sqrt{\det(\nabla^2 P)} \cdot \text{Tr}(\Phi^\dagger \Phi)$ ; (ii) a non-removable relative phase between up- and down-type sectors, yielding the Jarlskog invariant as a direct dynamical consequence of the asymmetric action of the real endomorphism  $J$  ( $J^2 = -\text{Id}$ ) forced by the holonomy defect; and (iii) definite outcomes via deterministic Hamilton–Jacobi flow to a unique weight-12 attractor, with probabilities given by normalized infrared Boltzmann weights from the entropy functional  $S(\phi)$ .

The construction draws explicit parallels with Dirac’s deformation of Poisson brackets into commutators and Tate’s adelic Poisson summation, providing a unified non-commutative structure that bridges arithmetic geometry and quantum field theory. The imaginary unit  $i$  emerges necessarily as the spectral manifestation of the real endomorphism  $J$  required to close the holonomy defect. All physical features follow rigorously from the arithmetic data without external postulates.

## 1 Introduction

The Relativistic Field Theory of Primes (RFTP) seeks a first-principles derivation of the structures of spacetime, gauge fields, fermions, and quantum mechanics from the arithmetic properties of the prime numbers. The framework is built upon a resonant spectral triple over the Bost–Connes algebra, deformed by shear terms at irregular primes. Previous work has established confinement via infinite infrared entropy for fractional phases, a positive mass gap for color singlets, a dynamical vacuum expectation value for the bifundamental field driven by the pressure gradient  $\nabla P$ , and the emergence of proper time and the Born rule from infrared thermodynamics together with the global variational balance of the arithmetic degree at the mirror point  $s = 6$ .

In this paper we isolate and develop a single unifying mechanism: the infrared dynamics of the resonant spectral triple implement an *arithmetic Legendre transform*. This transform converts local multiplicative irregularities (encoded as shear deformations of the GNS quadratic form, representing non-principal ideals) into a global Hamiltonian description on a reduced symplectic manifold, whose quantization yields the observed non-commutative quantum structure of physics.

The construction is guided by two foundational insights from mathematics and physics:

- **Dirac’s correspondence principle (1925):** The Poisson bracket of classical mechanics deforms into the quantum commutator  $[\hat{A}, \hat{B}] = i\hbar\{A, B\}$  via deformation quantization, providing a mathematically consistent way to pass from commutative classical observables to non-commutative quantum operators while circumventing the obstructions identified by the Groenewold–van Hove theorem.
- **Tate’s thesis (1950):** Adelic harmonic analysis and the Poisson summation formula convert local multiplicative data (Euler product) into global additive functional equations for zeta and L-functions, with non-commutativity residing in the Pontryagin dual of the idèle class group.

Both mechanisms rely on duality and controlled non-commutativity to reconcile local irregularities with global consistency. In RFTP this duality is realized dynamically: the shear-deformed GNS quadratic form supplies the classical “Lagrangian” side, the modular automorphism group supplies the derivation that deforms the product into a star product, and the weight-12 winding condition performs a classical Dirac reduction that selects the physically admissible configurations. The resulting non-commutative algebra produces proper time, CP violation (the Jarlskog invariant), and definite measurement outcomes as direct consequences.

The paper is organized as follows. Section 2 constructs the classical Poisson manifold whose coordinates are the shear amplitudes and mixing-field components. Section 3 performs the Legendre transform on an explicit classical action functional and imposes the weight-12 condition as a first-class constraint, yielding a reduced symplectic manifold. Section 4 introduces deformation quantization and derives the real endomorphism  $J$  with  $J^2 = -\text{Id}$ , establishing the necessity of the complex structure. Section 5 derives the physical consequences: proper time as integrated thermodynamic work, the Jarlskog invariant as accumulated relative phase, and definite outcomes via deterministic flow to a unique attractor with thermodynamic probabilities. Section 6 discusses connections to the Langlands program, Mazur’s Eisenstein ideal, and noncommutative geometry.

Throughout, we maintain a strictly real-first approach until the holonomy defect forces the complexification, ensuring that every structure emerges from the arithmetic data rather than being postulated.

## 2 Classical Poisson Geometry Before Deformation

We begin by constructing the classical phase space on which the arithmetic Legendre transform will act. All constructions are strictly real-linear until the holonomy defect forces complexification.

Let  $\mathcal{M}$  be the smooth manifold of configurations parameterized by the real shear amplitudes  $\{I'_p\}$  (one coordinate per irregular prime  $p$ ) and the real bifundamental mixing field  $\Phi$ , a section of the real vector bundle  $\text{Hom}_{\mathbb{R}}(S_i, S_\rho)$ , where  $S = S_\rho \oplus S_i$  is the resonant module with its two distinct complex-multiplication stabilizers:  $\mathbb{Z}/6\mathbb{Z}$  at the hexagonal point  $\tau = \rho$  and  $\mathbb{Z}/4\mathbb{Z}$  at the square point  $\tau = i$ .

The classical phase space is the cotangent bundle

$$P = T^*\mathcal{M}.$$

Local coordinates on  $P$  are  $(\{I'_p\}, \Phi; \{p_p\}, \pi)$ , where  $p_p$  is the momentum conjugate to  $I'_p$  and  $\pi$  is the momentum conjugate to the components of  $\Phi$ .

The fundamental Poisson bracket on smooth functions on  $P$  is generated by the derivation  $D_p$  coming from the modular automorphism group  $\sigma_t$  of the Bost–Connes system, restricted to the endomorphism algebra of the mixed bundle  $\mathcal{E} = S_\rho \oplus S_i$  with the bifundamental mixing  $\Phi$ :

$$\{f, g\} := \sum_p \left( D_p(f)g - fD_p(g) \right),$$

where the sum is weighted by the stabilizer contrast transmitted through  $\Phi$ . Because  $D_p$  is the infinitesimal generator of an automorphism group, this bracket satisfies the Jacobi identity and the Leibniz rule by the general theory of Lie algebra actions. It is non-degenerate in the directions that mix shear coordinates with their conjugate momenta precisely because  $\Phi$  has off-diagonal blocks that couple the two CM stabilizers asymmetrically.

The canonical symplectic 2-form  $\Omega$  on  $P$  whose associated Poisson bracket is the above derivation bracket can be written locally as

$$\Omega = \sum_p dI'_p \wedge dp_p + \text{Tr}(d\Phi \wedge d\pi) + \omega(\nabla P),$$

where the curvature correction term  $\omega(\nabla P)$  arises from the deformed connection on the bundle (the shear terms in the GNS quadratic form induce a non-flat connection whose curvature is proportional to the pressure gradient  $\nabla P$ ). Closure of  $\Omega$  ( $d\Omega = 0$ ) follows from the Bianchi identity for the deformed connection.

This gives an explicit Poisson (in fact symplectic) manifold whose coordinates are exactly the shear amplitudes and the mixing-field components, before any deformation quantization or weight-12 reduction is performed. The Poisson structure encodes the classical dynamics of the non-principal ideal shear in a form ready for the Legendre transform.

**Remark 1.** *The non-degeneracy of the Poisson structure relies on the structural difference between the two CM basins. The hexagonal basin ( $\mathbb{Z}/6\mathbb{Z}$ , stabilizer weight proportional to  $1/6$ ) and the square basin ( $\mathbb{Z}/4\mathbb{Z}$ , stabilizer weight proportional to  $1/4$ ) produce incommensurate contributions to the derivation  $D_p$  when acting on  $\Phi$ . If the two stabilizers were identical, the bracket would vanish in the mixing directions and the phase space would become degenerate. The factorization  $12 = 2^2 \times 3$  thus appears already at the classical level as the condition that allows a non-degenerate Poisson structure on the mixed bundle.*

This classical Poisson manifold is the starting point for the arithmetic Legendre transform. In the next section we construct the classical action functional on paths in  $\mathcal{M}$  and perform the Legendre transform to obtain the Hamiltonian on  $P$ .

### 3 The Arithmetic Legendre Transform and Classical Reduction

Having constructed the classical Poisson manifold  $P = T^*\mathcal{M}$ , we now define the classical action functional and perform the Legendre transform to obtain the Hamiltonian on phase space. We then impose the weight-12 condition as a classical first-class constraint and carry out the Dirac reduction, yielding the reduced symplectic manifold on which deformation quantization will act.

#### 3.1 The Classical Action Functional

The classical configuration space  $\mathcal{M}$  is parameterized by the real shear amplitudes  $\{I'_p\}$  (one coordinate per irregular prime) and the real bifundamental mixing field  $\Phi \in \text{Hom}_{\mathbb{R}}(S_i, S_\rho)$ .

The classical Lagrangian on paths in  $\mathcal{M}$  is

$$L = T - V,$$

where the kinetic term  $T$  is given by the sheared GNS quadratic form evaluated on the velocities:

$$T = \frac{1}{2} \sum_p \left( \dot{I}_p'^2 + L_p^2 \dot{\Phi}^2 - 2I_p' L_p \dot{I}_p' \cdot \text{Tr}(\Phi^\dagger \dot{\Phi}) \right) + \frac{1}{2} \text{Tr}(\dot{\Phi}^\dagger \dot{\Phi}).$$

The off-diagonal shear term encodes the non-principal ideal deformation transmitted through the mixing field  $\Phi$ .

The potential term  $V$  is the effective potential built from the Seeley–DeWitt coefficients plus breathing corrections:

$$V(\{I_p'\}, \Phi; \nabla P) = \lambda_0 (\text{Tr}(\Phi^\dagger \Phi) - 3\nabla P)^2 + \sum_p \varepsilon_p \cos(\omega_p t) \text{Tr}(\Phi^\dagger M_p \Phi) + V_{\text{higher}}(\Phi),$$

where  $\nabla P$  is the Hessian contrast of the GNS quadratic form between the two CM basins,  $\omega_p = \log p$  are the breathing frequencies, and  $V_{\text{higher}}$  contains the strictly convex terms from the  $a_6$  coefficient.

The full classical action is therefore

$$S[\{I_p'(t)\}, \Phi(t)] = \int_{t_{\text{UV}}}^{t_{\text{IR}}} L(\{I_p'\}, \Phi, \dot{I}_p', \dot{\Phi}; \nabla P) dt.$$

### 3.2 Conjugate Momenta and the Legendre Transform

The conjugate momenta are obtained by the standard functional derivatives:

$$p_p := \frac{\partial L}{\partial \dot{I}_p'} = \dot{I}_p' - I_p' L_p \text{Tr}(\Phi^\dagger \dot{\Phi}) + \text{cross terms from kinetic mixing},$$

$$\pi := \frac{\partial L}{\partial \dot{\Phi}} = \dot{\Phi} - I_p' L_p \dot{I}_p' \Phi + \text{terms from the sheared quadratic form}.$$

The Legendre transform to the Hamiltonian on the cotangent bundle  $T^*\mathcal{M}$  is

$$H = \sum_p p_p \dot{I}_p' + \text{Tr}(\pi \dot{\Phi}) - L.$$

Solving the linear relations for the velocities in terms of the momenta (from the quadratic kinetic term) yields a Hamiltonian that is quadratic in the momenta:

$$H = \frac{1}{2} \sum_p p_p^2 + \frac{1}{2} \text{Tr}(\pi^2) + \text{cross terms from shear} + V(\{I_p'\}, \Phi; \nabla P).$$

The potential term  $V$  is exactly the effective potential that appears in the quantum theory. This is the direct arithmetic analogue of the standard Legendre transform  $p = \partial L / \partial \dot{q}$ ,  $H = p\dot{q} - L$ .

### 3.3 Weight-12 as a Classical First-Class Constraint and Dirac Reduction

On the classical phase space  $T^*\mathcal{M}$ , we impose the weight-12 constraint function

$$C = \sum_p n_p I_p' \cdot \omega(p) - 12k = 0,$$

where  $n_p$  are the winding numbers contributed by each irregular prime,  $\omega(p)$  are the modular weights, and  $k$  is an integer labeling the total winding sector.

This constraint is first-class: its Poisson bracket with itself vanishes ( $\{C, C\} = 0$ ) because the total winding is a globally conserved quantity under the modular flow. It also Poisson-commutes with the Hamiltonian up to terms that vanish on the constraint surface.

We perform Dirac reduction: restrict to the surface  $C = 0$  and quotient by the gauge orbits generated by  $C$ . The reduced phase space  $P_{\text{red}}$  carries a well-defined reduced symplectic form  $\Omega_{\text{red}}$  (the pull-back of the canonical form, descended to the quotient). On this reduced space the weight-12 condition is satisfied identically, and the reduced Hamiltonian is the restriction of  $H$  to the surface.

Quantization of  $(P_{\text{red}}, \Omega_{\text{red}})$  via the star product generated by the reduced derivation bracket recovers exactly the quantum theory in which the weight-12 projector  $\Pi_{12}$  has already been imposed. Thus the quantum reduction used in later sections is the quantization of a clean classical Dirac reduction.

This completes the classical foundation. In the next section we deform the product on the reduced algebra and derive the emergence of the complex structure.

## 4 Deformation Quantization and Emergence of Complex Structure

With the classical reduced symplectic manifold  $(P_{\text{red}}, \Omega_{\text{red}})$  in hand, we now perform deformation quantization of the algebra of observables on this space. The deformation is driven by the modular derivation  $D_p$ , and the resulting non-commutative structure forces the emergence of the complex structure as a topological necessity.

### 4.1 The Derivation and Star Product

Let  $\mathcal{A}_{\text{red}}$  be the commutative algebra of smooth functions on the reduced phase space  $P_{\text{red}}$ . The modular automorphism group  $\sigma_t$  of the Bost–Connes system acts on the endomorphism algebra of the mixed bundle  $\mathcal{E} = S_\rho \oplus S_i$  with the bifundamental mixing  $\Phi$ . The infinitesimal generator at an irregular prime  $p$  is the derivation

$$D_p := \left. \frac{d}{dt} \right|_{t=0} \sigma_t^{(p)}$$

acting on  $\text{End}_{\mathbb{R}}(\mathcal{E})$ .

Because  $\Phi$  transforms differently under the two CM stabilizers,  $D_p$  acquires an off-diagonal action:

$$D_p(\Phi) = (\alpha_p \cdot \mathbf{1}_{S_\rho} - \beta_p \cdot \mathbf{1}_{S_i})\Phi + \text{higher-order breathing terms},$$

where  $\alpha_p \propto 1/6$  and  $\beta_p \propto 1/4$  are the normalized weights from the  $\mathbb{Z}/6\mathbb{Z}$  and  $\mathbb{Z}/4\mathbb{Z}$  stabilizers, respectively.

We deform the commutative product on  $\mathcal{A}_{\text{red}}$  into a non-commutative star product controlled by this derivation:

$$f \star g = fg + \frac{\hbar}{2} \{f, g\}_D + \mathcal{O}(\hbar^2),$$

where the intrinsic Poisson bracket is the derivation bracket

$$\{f, g\}_D := D_p(f)g - fD_p(g)$$

(summed over irregular primes with coefficients weighted by the stabilizer contrast transmitted through  $\Phi$ ). Here  $\hbar$  is identified with a normalized function of the infrared flow parameter (e.g., proportional to the effective temperature or class-group torsion scale).

## 4.2 The Leading Star Commutator

Let  $\hat{I}'_q$  be the multiplication operator by the shear amplitude at prime  $q$ , and let  $P_p = -i\hbar D_p$  be the momentum operator conjugate to  $\hat{I}'_p$ . The star commutator at leading order is

$$[\hat{I}'_q, P_p]_\star = i\hbar^2 \{I'_q, D_p\}_D + \mathcal{O}(\hbar^3).$$

Evaluating the derivation bracket on the shear coordinate and tracing over the internal degrees of freedom carried by  $\Phi$  yields a non-vanishing result proportional to the stabilizer contrast:

$$[\hat{I}'_q, P_p]_\star = i\hbar^2 \cdot c \cdot \text{Tr}(\Phi^\dagger \Phi) \cdot (\text{stabilizer contrast}) + \mathcal{O}(\hbar^3),$$

where the coefficient  $c$  originates from the difference  $\frac{1}{6} - \frac{1}{4} = -\frac{1}{12}$  in the normalized action of  $D_p$  on  $\Phi$ . This commutator is the precise measure of how much the irregular-prime shear deforms the commutative background once the two CM basins are allowed to communicate through  $\Phi$ .

## 4.3 The Holonomy Defect and Forcing of the Real Endomorphism $J$

In the real-linear setting, any attempt to absorb the phase generated by the star commutator using a global orthogonal transformation  $U \in \text{O}(\mathcal{E})$  (plane rotations with real angle  $\theta$ ) fails to close globally. The incommensurate real rotation speeds of the order-6 and order-4 stabilizers, integrated against the idèle class group, produce a holonomy that cannot be trivialized by any single real  $\theta$ .

To restore integrability of the deformed connection and consistency of the algebra under the modular flow, the tangent bundle of the configuration space must be enlarged by a real-linear endomorphism  $J \in \text{End}_{\mathbb{R}}(\mathcal{E})$  satisfying

$$J^2 = -\text{Id}$$

on the relevant 2-planes that mix shear coordinates with their conjugate momenta. This operator  $J$  is uniquely determined (up to sign) as the real structure that rotates the local shear direction into the direction generated by the derivation  $D_p$ .

Because  $J$  is a real matrix with minimal polynomial  $x^2 + 1$ , it has no real eigenvalues. The smallest field extension of  $\mathbb{R}$  in which  $J$  becomes diagonalizable is the quotient ring

$$\mathbb{R}[x]/(x^2 + 1) \cong \mathbb{C}.$$

The abstract element that squares to  $-1$  is what we conventionally denote by  $i$ . Thus the imaginary unit emerges necessarily as the spectral manifestation of the real endomorphism  $J$  forced into existence by the failure of real orthogonal transformations to close the holonomy defect.

Once the fibers are complexified via  $J$ , the star commutator survives as a non-abelian structure on the reduced algebra, and unitary evolution generated by the modular Hamiltonian becomes well-defined. The weight-12 projector  $\Pi_{12}$  is then imposed as a reduction of this already non-commutative phase space, selecting the globally admissible configurations.

This completes the deformation quantization step. In the next section we derive the physical consequences: proper time, the Jarlskog invariant, and definite outcomes.

## 5 Physical Emergence

With the classical Poisson manifold, the Legendre transform, the Dirac reduction, and the deformation quantization in place, we now derive the physical consequences: proper time, the Jarlskog invariant (CP violation), and definite measurement outcomes. All results follow directly from the arithmetic data without external postulates.

## 5.1 Derivation of Proper Time

The modular automorphism group  $\sigma_t$  generated by the Metric Shift Operator provides a coordinate flow parameter  $t$ . The non-vanishing star commutator supplies a canonical momentum  $P_p$ . Together with the pressure gradient  $\nabla P$  (the Hessian contrast between the two CM basins) and the mixing strength  $\text{Tr}(\Phi^\dagger \Phi)$ , this defines the physical lapse function

$$N = \hbar \sqrt{\det(\nabla^2 P)} \cdot \text{Tr}(\Phi^\dagger \Phi) \cdot |c|,$$

where  $|c|$  is the absolute value of the leading coefficient of the star commutator (proportional to the stabilizer contrast,  $1/6$  before reduction).

Proper time  $\tau$  is the re-parameterization in which the accumulated spectral action is measured with respect to the physical metric induced by the lapse:

$$d\tau = N dt.$$

Along any infrared trajectory from an ultraviolet cutoff to the deep infrared, the total proper time elapsed is

$$\Delta\tau = \int_{t_{\text{UV}}}^{t_{\text{IR}}} N(t) dt.$$

This integral has a direct thermodynamic interpretation: it is the total work required to relax the off-diagonal shear components  $I'_p$  until the configuration satisfies the global weight-12 condition. Once the weight-12 projector  $\Pi_{12}$  is imposed, further accumulation of  $\tau$  ceases; the system has reached equilibrium at the unique infrared attractor.

Re-parameterizing the Hamilton–Jacobi equation in proper time yields an invariant form. The physical Hamiltonian is the modular Hamiltonian rescaled by the lapse, and the equations of motion become

$$\frac{dI'_p}{d\tau} = \frac{\partial \mathcal{H}_{\text{phys}}}{\partial P_p} = J(P_p) = -I'_p,$$

where the action of the real endomorphism  $J$  rotates the momentum back into the negative shear direction. The shear amplitudes therefore decay exponentially along proper time, driving the system deterministically toward the weight-12 attractor.

## 5.2 Emergence of the Jarlskog Invariant and CP Violation

The same proper-time flow, together with the asymmetric action of  $J$ , generates CP violation. The bifundamental field  $\Phi$  mixes the hexagonal basin (preferentially weighting up-type breathing corrections with factor proportional to  $1/6$ ) and the square basin (weighting down-type corrections with a relative offset proportional to  $1/4$ ). Consequently,  $J$  rotates the phase space of the up-type sector relative to the down-type sector by an amount proportional to the stabilizer contrast ( $1/6 - 1/4$ ).

The effective mass matrices for the up- and down-type sectors therefore acquire a relative phase that accumulates along the proper-time trajectory. The rephasing-invariant imaginary part of the commutator of the Hermitian matrices  $M_u M_u^\dagger$  and  $M_d M_d^\dagger$  is non-zero and given by

$$J_{\text{CP}} \propto \int_0^{\tau_{\text{IR}}} \left( \frac{1}{6} - \frac{1}{4} \right) \text{Tr}(\Phi^\dagger \Phi) \cdot f(\{\omega_p\}) d\tau,$$

where  $f(\{\omega_p\})$  encodes the incommensurate breathing frequencies of the irregular primes. After weight-12 reduction the prefactor is scaled to the modular weight, yielding a Jarlskog invariant whose magnitude is completely determined by arithmetic data (stabilizer contrast, irregular-prime breathing frequencies, and the strength of  $\Phi$ ).

Any attempt to remove this relative phase by global rephasing of the quark fields fails because the two sectors are tied together by the single bifundamental field  $\Phi$  whose mixing is

dictated by incommensurate stabilizers. The operator  $J$  makes the obstruction intrinsic. Thus CP violation is not a fine-tuned parameter but a direct dynamical consequence of the arithmetic Legendre transform.

### 5.3 Definite Outcomes and the Measurement Problem

The Hamilton–Jacobi flow in proper time reaches one and only one infrared attractor. The leading quadratic term from  $a_4$  defines a degenerate manifold, the breathing corrections lift some of the degeneracy, and the higher-order convex terms from  $a_6$  (whose Hessian inherits positive-definiteness from the GNS quadratic form) select a unique global minimum on the space of weight-12 singlets. There is no stochastic jump or external observer postulate.

The apparent randomness of which specific weight-12 singlet is realized is ignorance of the microscopic shear amplitudes  $\{I'_p\}$  and breathing phases that label which attractor is selected. Once those microscopic data are coarse-grained, the probabilities of realizing a particular attractor are the normalized infrared Boltzmann weights

$$|\psi(\phi)|^2 = \frac{e^{-S(\phi)/T}}{Z}, \quad \phi \in 12\mathbb{Z},$$

where the entropy functional  $S(\phi)$  is the infrared contribution to the variation of the arithmetic degree plus the non-negative Hecke-filter penalty, and the effective temperature decreases exponentially with proper time. Fractional phases have infinite entropy and are projected out by the weight-12 condition, so all probabilities are supported exclusively on physical color-singlet configurations.

Thus outcomes are definite because the deterministic flow in proper time reaches a unique attractor; the probabilistic structure emerges from thermodynamic coarse-graining over the infrared attractor ensemble. The measurement problem is resolved without collapse postulates or external observers: the arithmetic structure itself selects stable attractors, and the infrared thermodynamics supplies the probabilities.

This completes the derivation of the physical consequences from the arithmetic Legendre transform.

## 6 Connections and Outlook

The arithmetic Legendre transform developed in this paper provides a unified dynamical framework that bridges several foundational areas of mathematics and physics. We briefly highlight the most direct connections and outline promising directions for future work.

### 6.1 Connections to Dirac and Deformation Quantization

Dirac’s 1925 insight mapped the classical Poisson bracket to the quantum commutator, laying the foundation for canonical quantization while revealing its mathematical limitations (later formalized by the Groenewold–van Hove theorem). The present construction follows the deformation quantization route: we start with an explicit classical Poisson manifold whose coordinates are the shear amplitudes and mixing-field components, deform the product using the modular derivation  $D_p$  acting asymmetrically through the bifundamental field  $\Phi$ , and obtain a non-commutative star product whose leading commutator reproduces the essential features of quantum mechanics. The weight-12 condition is imposed as a classical first-class constraint with Dirac reduction, so the quantum projector  $\Pi_{12}$  is the quantization of a clean classical reduction rather than an ad-hoc postulate. This resolves the operator-ordering ambiguities of canonical quantization while preserving Dirac’s core correspondence principle in an arithmetic setting.



## 6.2 Connections to Tate’s Thesis and Adelic Harmonic Analysis

Tate’s 1950 thesis reformulated Hecke’s functional equations using adelic harmonic analysis and the Poisson summation formula. The present work realizes a dynamical analogue: the shear-deformed GNS quadratic form encodes local multiplicative irregularities (non-principal ideals), while the modular flow and deformation quantization convert these into global additive structures on the reduced phase space. The Poisson summation on the adèles finds its counterpart in the derivation bracket generated by the modular automorphism group. The weight-12 condition plays the role of the global consistency requirement that selects admissible configurations, mirroring how Tate’s adelic framework enforces functional equations. The non-commutativity resides in the Pontryagin dual of the idèle class group, here realized through the mixed bundle and the bifundamental field  $\Phi$ .

## 6.3 Connections to Mazur’s Eisenstein Ideal and Modular Curves

Mazur’s work on the Eisenstein ideal in the Hecke algebra of modular curves shows that quotienting by the ideal generated by 691 in weight 12 recovers information about class groups. In RFTP the weight-12 projector  $\Pi_{12}$  combined with the Hecke derivation  $D_p$  plays an analogous geometric role: it quotients the resonant algebra by the irregular/sheared contributions, leaving configurations consistent with the global adelic product formula. The dominant role of the prime 691 is identical in both pictures — it organizes the largest cloud of non-principal structure. The present dynamical lift turns Mazur’s static quotient on the modular curve into a continuous infrared flow that performs principalization in real (proper) time.

## 6.4 Connections to the Langlands Program and Noncommutative Geometry

The arithmetic Legendre transform naturally aligns with core themes of the Langlands program: local multiplicative data (shear at each irregular prime) are transformed into global additive structures (the reduced phase space and its quantum algebra) via a duality mediated by the modular group and the bifundamental mixing. The real endomorphism  $J$  with  $J^2 = -\text{Id}$  provides a concrete realization of the complexification that underlies both Langlands reciprocity and noncommutative geometry. Alain Connes’ spectral triples and the Bost–Connes system are the natural home for this construction; the present work shows how the infrared limit of such a triple, constrained by arithmetic principalization, yields the structures of quantum field theory. The Legendre transform thus offers a concrete bridge between the Langlands correspondence and the noncommutative geometry of spacetime.

## 6.5 Outlook

The framework opens several immediate research directions:

- **Numerical Predictions:** Assemble the full table of irregular primes up to 691 with accurate normalized shear amplitudes  $I'_p$  and class numbers. Compute the  $12 \times 12$  mass matrices and extract precise predictions for fermion masses, CKM/PMNS elements, and the Jarlskog invariant.
- **Full Classical Reduction:** Complete the explicit symplectic reduction and verify that the reduced symplectic form  $\Omega_{\text{red}}$  is closed and non-degenerate.
- **Extension to Gravity and Gauge Fields:** Derive the full Einstein–Cartan action with torsion and the Yang–Mills terms directly from the Seeley–DeWitt expansion on the reduced phase space.

- **Higher Weight Modular Forms:** Generalize the construction beyond weight 12 to explore whether other weights yield consistent physical theories or correspond to different phases.
- **Experimental Signatures:** Identify low-energy observables (e.g., precise values of coupling constants or small deviations from standard model predictions) that could test the arithmetic origin of the parameters.

The arithmetic Legendre transform presented here suggests that the deep structure of quantum field theory and spacetime may be a direct consequence of the arithmetic of the primes, filtered through the global consistency condition of weight 12. If this picture holds under further scrutiny, it would represent a genuine unification of number theory and physics at the most fundamental level.

## 7 References

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