

Quaternionic Norm Response Theory (QNRT): Discrete Spacetime, Photon Basis, UV–IR Scale Duality, and Cosmological Scale Estimates

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Abstract

Quaternionic Norm Response Theory (QNRT) proposes a discrete, geometrically structured spacetime based on the Hurwitz quaternion lattice Γ_H . Through a single algebraic operation—quaternion inversion—the theory generates a \mathbb{Z}_2 UV–IR scale duality whose fixed point simultaneously determines the dark-energy scale and the neutrino mass scale.

The dynamical field $\phi : \Gamma_H \rightarrow \mathbb{H}_{\mathbb{C}}$ takes values in the complexified quaternion algebra $\mathbb{H}_{\mathbb{C}} \cong M_2(\mathbb{C})$ and carries the Hermitian norm $N(\phi) = \phi\phi^\dagger$, giving the logarithmic scale field $\sigma = \log(N(\phi)/\ell_P^2)$. Inversion $\phi \mapsto \phi^{-1} = \phi^\dagger/N(\phi)$ sends $\sigma \mapsto -\sigma$, exchanging ultraviolet and infrared regimes. The norm-response scaling (NRS) theorem shows that this map has a unique fixed point $r_c = \sqrt{\ell_P R_\star}$.

The fixed-point scale $r_c = 84 \mu\text{m}$ (with $R_\star \approx 4.4 \times 10^{26} \text{ m}$, $N_\star = R_\star/\ell_P \approx 2.70 \times 10^{61}$) defines a reference wavenumber $k_0 = 1/r_c$. This is precisely the wavenumber of the self-dual lattice mode ($n = n_c = \sqrt{N_\star}$ in the harmonic expansion over R_\star), an algebraic identity verified to 0.000% numerical error. The corresponding energy $E_0 = \hbar ck_0 = 2.354 \text{ meV}$ and temperature $T_c = 27.27 \text{ K}$ equal the geometric mean of the Planck and de Sitter temperatures, and coincide with the CMB temperature at redshift $z_c \approx 9$ (cosmic reionisation epoch).

The self-dual scale gives $\rho_\Lambda = \hbar c/(\ell_P R_\star)^2 = c^4/(GR_\star^2) \approx 6.3 \times 10^{-10} \text{ J m}^{-3}$, within 21% of the Planck 2018 value $5.3 \times 10^{-10} \text{ J m}^{-3}$. The identity $\rho_\Lambda = c^4/(GR_\star^2)$ shows that \hbar cancels algebraically, so the cosmological constant problem reduces to the purely classical question of why $R_\star/\ell_P \approx 2.7 \times 10^{61}$. The fermionic counterpart yields $M_c c^2 = \sqrt{M_P c^2 \cdot M_H c^2} \approx 2.4 \text{ meV}$, within the active-neutrino mass window. The lattice-derived kinetic coefficient $c_{\text{lat}} = 6$ fixes the continuum action; the duality-symmetric potential $V(\sigma) = 2\rho_\Lambda \cosh \sigma$ produces thawing-quintessence dark energy with $w_0 \in [-0.990, -0.908]$ and $w_a \in [-0.082, -0.014]$ for $\sigma_i \in [0.5, 3.0]$. All predictions are consistent with ΛCDM .

The principal limitations are as follows. The value $R_\star \approx 4.4 \times 10^{26} \text{ m}$ is set by self-consistency of the numerical outputs; its epoch-independent determination is Open Problem OP1 and is not resolved in this paper. The QNRT predictions $w_0 \in [-0.990, -0.908]$ are inconsistent with the DESI DR2 (2025) central value $w_0 \approx -0.75$ at 2.7–4.0 σ , while remaining consistent with ΛCDM . The Euclid forecasted precision $\sigma(w_0) \sim 0.02$ will provide a decisive test. The Standard Model gauge

embedding and the physical determination of σ_i are the remaining principal open problems.

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1 Introduction

Two deep mysteries of modern physics—the cosmological constant problem [1] and the smallness of neutrino masses—lack algebraic explanations within the Standard Model. The observed dark-energy density $\rho_\Lambda \approx 6 \times 10^{-10} \text{ J m}^{-3}$ is 10^{122} times smaller than Planck-scale estimates, while neutrino masses at the meV scale are separated from other fermion masses by many orders of magnitude.

QNRT proposes that both scales have a common geometric origin: spacetime is discrete at the Planck scale, with the Hurwitz quaternion lattice Γ_H adopted as its fundamental structure. This choice is motivated by the algebraic properties of the complexified quaternion algebra $\mathbb{H}_\mathbb{C} \cong M_2(\mathbb{C})$, which serves as the field algebra of the theory (Section 3). Within this adopted structure, the cosmological hierarchy is a consequence of a UV–IR duality generated by quaternion inversion on the field algebra. The proposal is not a complete quantum theory of gravity; it is a framework in which the Planck and Hubble scales are related by an algebraic fixed-point condition.

This revised version (v3.3) incorporates the following corrections relative to v3.1 and v3.2:

- (C1) **Algebraic basis of the Minkowski interval.** The element $q = ict + x\mathbf{i} + y\mathbf{j} + z\mathbf{k}$ appearing in the spacetime motivation belongs to $\mathbb{H}_\mathbb{C}$, not to the real quaternion algebra \mathbb{H} . Real quaternion norms are positive-definite (Euclidean); the Minkowski signature $(-, +, +, +)$ arises from the Hermitian norm on $\mathbb{H}_\mathbb{C}$. The text has been corrected accordingly (Section 2).
- (C2) **Asymptotic formula for $r_4(N)$.** The saddle-point argument in Section 6.3 previously cited $r_4(N) \sim 8\pi^2 N$; the correct average asymptotic is $r_4(N) \sim \pi^2 N$ (confirmed numerically to 0.05%).
- (C3) **Colour gauge group.** The identification of $G_2 \supset \text{SU}(3)$ as the lattice automorphism group for colour is withdrawn; explicit calculation shows that the Weyl group $W(G_2)$ generates components $\pm 1/3, \pm 2/3$ not in Γ_H , ruling out this route. The current status of the $\text{SU}(3)$ embedding is stated as open (Section 10).
- (C4) **E_8 and anomaly cancellation.** The claim that $E_8 \cong \Gamma_H \oplus \Gamma_H$ provides the context for anomaly cancellation is withdrawn. The E_8 extension is blocked by the non-associativity of the Cayley integers and gives $c_{\text{lat}}(E_8) = 60 \neq 6$, inconsistent with the lattice action. The anomaly cancellation question is left as an open problem (Section 10).
- (C5) **R_\star numerical value.** Earlier versions stated $R_\star \approx 1.3 \times 10^{26} \text{ m} \approx R_H(t_0)$ (the Hubble radius, c/H_0). However, all specific numerical outputs— $r_c = 84 \mu\text{m}$, k_0 , E_0 , T_c , N_\star , and $\rho_\Lambda/\rho_\Lambda^{\text{obs}} = 1.209$ —are mutually consistent only for $R_\star \approx 4.4 \times 10^{26} \text{ m}$ ($N_\star \approx 2.70 \times 10^{61}$), which is approximately $3.2 R_H$, close to the comoving particle horizon. Setting $R_\star = R_H$ gives $r_c \approx 47 \mu\text{m}$, inconsistent with all other quoted values. This version uses $R_\star \approx 4.4 \times 10^{26} \text{ m}$ throughout; its physical determination remains Open Problem OP1.
- (C6) **ρ_Λ numerical value.** With $R_\star \approx 4.4 \times 10^{26} \text{ m}$: $\rho_\Lambda = c^4/(GR_\star^2) \approx 6.3 \times 10^{-10} \text{ J m}^{-3}$, within 21% of $\rho_\Lambda^{\text{obs}} \approx 5.3 \times 10^{-10} \text{ J m}^{-3}$ (Planck 2018). The earlier value $7 \times 10^{-9} \text{ J m}^{-3}$ was obtained with $R_\star = R_H$ and is superseded.

(C7) “12% discrepancy” corrected. $\rho_\Lambda^{\text{QNRT}}/\rho_\Lambda^{\text{obs}} = 1.209$ is a 20.9% excess, not 12% as stated in earlier versions.

(C8) Neutrino mass estimate revised. With $R_\star = 4.4 \times 10^{26}$ m used consistently, $M_c c^2 = \sqrt{M_P c^2 \cdot M_H c^2} \approx 2.4$ meV (revised from 4.2 meV, which used $R_\star = R_H$). The value remains within the active-neutrino mass window.

The dark-energy equation-of-state predictions (w_0, w_a) are unchanged from v3.1. The values of ρ_Λ and $M_c c^2$ are revised in v3.3 (corrections C6 and C8).

Derivation chain.

$$\text{Lorentz invariance} \Rightarrow N_{\mathbb{H}_\mathbb{C}}(q) = -c^2 t^2 + |\mathbf{x}|^2 \Rightarrow \mathbb{H}_\mathbb{C} \cong M_2(\mathbb{C})$$

Hurwitz lattice $\Gamma_H \subset \mathbb{H}$ integer norm, 24 unit elements

$$\text{NRS map: } \phi \mapsto \phi^{-1} = \phi^\dagger / N(\phi) \Rightarrow N(\phi) \mapsto 1/N(\phi), \sigma \mapsto -\sigma$$

$$\text{NRS theorem: unique fixed point } r_c = \sqrt{\ell_P R_\star}$$

Photon basis: $k_0 = 1/r_c$ self-dual lattice mode (algebraic identity)

$$\rho_\Lambda = c^4 / (G R_\star^2) \approx 6.3 \times 10^{-10} \text{ J m}^{-3}, \quad M_c c^2 \approx 2.4 \text{ meV}, \quad \text{dark-energy dynamics } w(z)$$

2 The Hurwitz Lattice as Discrete Spacetime

2.1 Algebraic motivation: complexified quaternions and the Minkowski interval

The quaternion algebra \mathbb{H} over \mathbb{R} has basis $\{1, \mathbf{i}, \mathbf{j}, \mathbf{k}\}$ with $\mathbf{i}^2 = \mathbf{j}^2 = \mathbf{k}^2 = \mathbf{ijk} = -1$. For a real quaternion $q = a + b\mathbf{i} + c\mathbf{j} + d\mathbf{k}$ with $a, b, c, d \in \mathbb{R}$, the norm is $N(q) = q\bar{q} = a^2 + b^2 + c^2 + d^2 \geq 0$, which is positive-definite (Euclidean) and cannot yield the Minkowski signature directly.

The Minkowski interval $-c^2 t^2 + x^2 + y^2 + z^2$ arises from the *complexified* quaternion algebra $\mathbb{H}_\mathbb{C} = \mathbb{H} \otimes_\mathbb{R} \mathbb{C}$. Consider the element

$$q(x) = ict + x\mathbf{i} + y\mathbf{j} + z\mathbf{k} \in \mathbb{H}_\mathbb{C}, \quad (1)$$

whose Hermitian norm is

$$N(q) = qq^\dagger = -c^2 t^2 + x^2 + y^2 + z^2. \quad (2)$$

The Minkowski interval thus emerges from the Hermitian norm on $\mathbb{H}_\mathbb{C}$, not from the Euclidean norm on \mathbb{H} . This observation motivates $\mathbb{H}_\mathbb{C}$ as the algebra for the dynamical field ϕ (Section 3), while the lattice Γ_H is defined within \mathbb{H} using the Euclidean norm.

Remark 2.1 (Two roles of \mathbb{H} and $\mathbb{H}_\mathbb{C}$). *The theory uses \mathbb{H} and $\mathbb{H}_\mathbb{C}$ in distinct roles: $\Gamma_H \subset \mathbb{H}$ defines the discrete geometry of spacetime (the lattice sites), while $\mathbb{H}_\mathbb{C} \cong M_2(\mathbb{C})$ is the algebra of field values $\phi: \Gamma_H \rightarrow \mathbb{H}_\mathbb{C}$. The Euclidean norm on \mathbb{H} governs the lattice structure; the Hermitian norm $N(\phi) = \phi\phi^\dagger$ on $\mathbb{H}_\mathbb{C}$ governs the field dynamics.*

2.2 The Hurwitz lattice

Definition 2.2 (Hurwitz lattice).

$$\Gamma_H = \left\{ n_0 + n_1 \mathbf{i} + n_2 \mathbf{j} + n_3 \mathbf{k} \mid n_\mu \in \mathbb{Z} \text{ or } n_\mu \in \mathbb{Z} + \frac{1}{2} \right\}. \quad (3)$$

Γ_H is the unique maximal order in \mathbb{H} closed under conjugation and admitting a Euclidean division algorithm [2, 3]. It corresponds to the F_4 root lattice—the densest sphere packing in four dimensions [4]—and provides the discrete substrate for the dynamical field.

Proposition 2.3 (Integer norm and unit elements). *For all $\phi \in \Gamma_H$: (i) $N(\phi) = \phi \bar{\phi} \in \mathbb{Z}_{\geq 0}$; (ii) the 24 elements with $N(u) = 1$ form the binary tetrahedral group: $\{\pm e_\mu\}$ (8 elements) and $\frac{1}{2}(\pm 1 \pm \mathbf{i} \pm \mathbf{j} \pm \mathbf{k})$ (16 elements); (iii) $N(\alpha\beta) = N(\alpha)N(\beta)$.*

Remark 2.4 (Why quaternions, and why Γ_H ?). *The norm-response operation $\phi \mapsto \phi^{-1} = \phi^\dagger / N(\phi)$ requires (a) a multiplicative norm and (b) existence of an inverse for all non-zero ϕ . By Hurwitz’s composition theorem [2], the only finite-dimensional real division algebras with a multiplicative norm are \mathbb{R} , \mathbb{C} , \mathbb{H} , and \mathbb{O} . Of these, \mathbb{H} is the unique four-dimensional choice matching the dimensionality of spacetime. Within \mathbb{H} , Γ_H is adopted because (i) it is the unique maximal order; (ii) its Euclidean algorithm guarantees unique prime factorisation; and (iii) its 24-element unit group is the binary tetrahedral group, a discrete subgroup of $SU(2)$, which generates the $SU(2)$ gauge sector (Section 9).*

2.3 Relation to other discrete-spacetime approaches

Discrete models of spacetime include spin foam models [14], causal dynamical triangulations [15], and causal set theory [16]. QNRT differs in two respects. First, the discrete structure is Γ_H , which carries the full quaternion algebra. Second, the norm-response mechanism links Planck-scale physics directly to the cosmological scale R_\star through an algebraic fixed-point condition, rather than through a coarse-graining argument.

3 The Norm-Response Mechanism

3.1 Field domain and norms on $\mathbb{H}_\mathbb{C}$

The dynamical field takes values in the complexified quaternion algebra:

$$\phi : \Gamma_H \rightarrow \mathbb{H}_\mathbb{C} = \mathbb{H} \otimes_{\mathbb{R}} \mathbb{C}. \quad (4)$$

Write $\phi = A + B\mathbf{j}$ with $A, B \in \mathbb{H}$. The Hermitian conjugate $\phi^\dagger = \bar{A} - \bar{B}\mathbf{j}$ gives $N(\phi) \equiv \phi\phi^\dagger \in \mathbb{R}_{\geq 0}$. The complexification is physically necessary: Dirac spinors and complex gauge representations cannot be accommodated within \mathbb{H} alone; $\mathbb{H}_\mathbb{C} \cong M_2(\mathbb{C})$ is the minimal algebra admitting a single Weyl spinor [5].

3.2 Quaternion inversion and the \mathbb{Z}_2 duality

Definition 3.1 (Norm-response operation). *The map*

$$\phi \mapsto \phi^{-1} = \frac{\phi^\dagger}{N(\phi)} \quad (5)$$

is called the norm-response operation. The new norm satisfies $N(\phi^{-1}) = 1/N(\phi)$.

Definition 3.2 (Logarithmic scale field). *In Planck units ($\ell_P = 1$, $N_{\text{UV}} = 1$):*

$$\sigma(\phi) = \log N(\phi) \in \mathbb{R}. \quad (6)$$

Corollary 3.3 (\mathbb{Z}_2 norm-response duality).

$$\sigma(\phi^{-1}) = -\sigma(\phi). \quad (7)$$

The inversion $\phi \mapsto \phi^{-1}$ exchanges large norms (UV, $\sigma > 0$) with small norms (IR, $\sigma < 0$). This \mathbb{Z}_2 UV–IR exchange is the single algebraic operation from which the entire theory flows.

3.3 Field decomposition

The decomposition $\phi = e^{\sigma/2}u$, $u \in \text{SU}(2)$ ($u\bar{u} = 1$), separates the scale sector σ (local Weyl rescaling $g_{\mu\nu} \rightarrow e^{2\sigma}g_{\mu\nu}$) from the directional sector u . This decomposition plays a role in the matter coupling (Section 9).

4 The Self-Dual Scale

4.1 Scale interval and the effect of inversion

The theory has UV cutoff ℓ_P and IR reference scale R_* , so $\sigma \in [0, S]$ with $S = \log(R_*/\ell_P) \approx 141$. Under inversion (7): $\sigma \mapsto -\sigma$, a reflection about $\sigma = 0$ (Planck scale), not about the midpoint.

4.2 Norm-response scaling theorem

Definition 4.1 (Norm-response scaling (NRS)).

$$\tilde{\phi} = \phi^{-1} \cdot \sqrt{N_{\text{IR}}}, \quad N_{\text{IR}} = e^S = R_*/\ell_P. \quad (8)$$

Theorem 4.2 (NRS implements UV–IR duality). *Under (8):*

$$\sigma(\tilde{\phi}) = S - \sigma(\phi). \quad (9)$$

The unique fixed point is

$$\sigma(\phi_*) = \frac{S}{2} \iff N(\phi_*) = e^{S/2} \iff r_c = \sqrt{\ell_P R_*}. \quad (10)$$

Proof. $N(\tilde{\phi}) = N(\phi^{-1}) \cdot N_{\text{IR}} = N_{\text{IR}}/N(\phi)$. Thus $\sigma(\tilde{\phi}) = \log(N_{\text{IR}}/N(\phi)) = S - \sigma(\phi)$. Fixed point: $S - \sigma = \sigma \Rightarrow \sigma = S/2$. \square

5 Self-Dual Scale, Vacuum Energy, and Photon Basis

5.1 Vacuum energy

By dimensional analysis, the vacuum energy density associated with r_c is

$$\rho_\Lambda = \frac{\hbar c}{r_c^4} = \frac{\hbar c}{(\ell_P R_*)^2}. \quad (11)$$

With $\ell_P \approx 1.616 \times 10^{-35}$ m and $R_\star \approx 4.36 \times 10^{26}$ m ($N_\star = R_\star/\ell_P \approx 2.70 \times 10^{61}$): $r_c = \sqrt{\ell_P R_\star} \approx 84 \mu\text{m}$ and $\rho_\Lambda \approx 6.3 \times 10^{-10}$ J m $^{-3}$, within 21% of $\rho_\Lambda^{\text{obs}} \approx 5.3 \times 10^{-10}$ J m $^{-3}$ [6].

Remark 5.1 (Numerical value of R_\star). *The working value $R_\star \approx 4.36 \times 10^{26}$ m (≈ 46 Gly, close to the comoving particle horizon) is not the Hubble radius $R_H = c/H_0 \approx 1.37 \times 10^{26}$ m. Setting $R_\star = R_H$ would give $r_c \approx 47 \mu\text{m}$, inconsistent with all other numerical outputs of this paper ($k_0, E_0, T_c, N_\star, \rho_\Lambda/\rho_\Lambda^{\text{obs}}$). The physical identification of R_\star is Open Problem OP1; the value used here is fixed by self-consistency of the numerical outputs.*

Remark 5.2 (\hbar cancellation and the classical identity). *Substituting $\ell_P^2 = \hbar G/c^3$ into (11):*

$$\rho_\Lambda = \frac{\hbar c}{(\ell_P R_\star)^2} = \frac{\hbar c \cdot c^3}{\hbar G R_\star^2} = \frac{c^4}{G R_\star^2}. \quad (12)$$

The factor \hbar cancels algebraically (verified numerically to 0.14%, solely the accumulated rounding error in published constants). Equation (12) shows ρ_Λ depends only on c, G , and R_\star . The cosmological constant problem—why $\rho_\Lambda/\rho_{\text{Planck}} \approx 10^{-123}$ —translates into the purely classical-gravitational question: why is $R_\star/\ell_P \approx 2.7 \times 10^{61}$? This reframing is consistent with the observation [1] that a classical theory involving only G, c , and a macroscopic length carries no fine-tuning problem. QNRT provides the algebraic mechanism (the NRS fixed point) that selects $r_c = \sqrt{\ell_P R_\star}$, but the physical determination of R_\star is OP1.

Remark 5.3 (R_\star and the 21% discrepancy (OP1)). *With $R_\star \approx 4.36 \times 10^{26}$ m:*

$$\frac{\rho_\Lambda^{\text{QNRT}}}{\rho_\Lambda^{\text{obs}}} = \frac{6.3 \times 10^{-10}}{5.3 \times 10^{-10}} \approx 1.209 \quad (20.9\% \text{ excess}).$$

In terms of density parameters, this corresponds to $\Omega_\Lambda^{\text{QNRT}} \approx 0.83$ versus the observed $\Omega_\Lambda = 0.685$. The discrepancy is not resolved in this paper; it constitutes the coincidence problem within QNRT (OP1).

5.2 Photon basis and reference wavenumber

The self-dual scale r_c defines a natural reference wavenumber, which we call the *photon basis* of QNRT.

Definition 5.4 (Reference wavenumber and photon basis).

$$k_0 := \frac{1}{r_c}, \quad E_0 := \hbar c k_0 = \frac{\hbar c}{r_c}, \quad T_c := \frac{E_0}{k_B} = \frac{\hbar c}{k_B r_c}. \quad (13)$$

Numerically: $k_0 = 1.190 \times 10^4$ m $^{-1}$, $E_0 = 2.354$ meV, $T_c = 27.27$ K (all consistent with $R_\star \approx 4.36 \times 10^{26}$ m, $N_\star \approx 2.70 \times 10^{61}$).

Proposition 5.5 (Geometric mean property). *T_c is the geometric mean of the Planck and de Sitter temperatures:*

$$T_c = \sqrt{T_{\text{Planck}} \times T_{dS}}, \quad T_{\text{Planck}} = \frac{\hbar c}{k_B \ell_P}, \quad T_{dS} = \frac{\hbar c}{k_B R_\star}. \quad (14)$$

Proof. $\sqrt{T_{\text{Planck}} \cdot T_{dS}} = \hbar c / (k_B \sqrt{\ell_P R_\star}) = \hbar c / (k_B r_c) = T_c$. Verified numerically to 0.000% error with $R_\star = 4.36 \times 10^{26}$ m. \square

Proposition 5.6 (Self-dual lattice mode). *The Hurwitz lattice harmonic modes over R_\star have wavenumbers $k_n = n/R_\star$ for $n = 1, 2, \dots, N_\star$. The self-dual mode $n = n_c = \sqrt{N_\star}$ has wavenumber*

$$k_{n_c} = \frac{n_c}{R_\star} = \frac{\sqrt{N_\star}}{R_\star} = \frac{1}{\sqrt{\ell_P R_\star}} = \frac{1}{r_c} = k_0. \quad (15)$$

Proof. $n_c = \sqrt{N_\star} = \sqrt{R_\star/\ell_P}$, so $n_c/R_\star = (R_\star/\ell_P)^{1/2}/R_\star = 1/\sqrt{\ell_P R_\star} = k_0$. Verified numerically to 0.000% error. \square

Equation (15) is an algebraic identity requiring no additional assumption. The self-dual lattice mode $n = n_c$ is simultaneously: (a) the geometric midpoint of the UV–IR spectrum in logarithmic wavenumber; (b) the NRS fixed-point wavenumber $k_0 = 1/r_c$; and (c) the photon wavenumber whose energy $E_0 = \hbar c k_0$ corresponds to temperature $T_c = 27.27$ K.

Cosmological interpretation. The CMB temperature equals T_c at redshift

$$z_c = T_c/T_{\text{CMB},0} - 1 = 27.27/2.7255 - 1 \approx 9.0, \quad (16)$$

coinciding with the epoch of cosmic reionisation ($z \approx 6$ – 10). This is a post-diction consistent with Λ CDM, not a new prediction.

6 Effective Action from the Lattice

6.1 Lattice action

The dynamical content of the scale sector σ is encoded in a nearest-neighbour action on Γ_H :

$$S_{\text{lat}} = \sum_{\langle ij \rangle} \frac{|\phi_i - \phi_j|^2}{|\phi_i|^2}. \quad (17)$$

This is the simplest quadratic lattice action invariant under multiplication by a unit element $u \in \Gamma_H$. Under the NRS map, each term transforms by exchanging the roles of i and j , so the action is \mathbb{Z}_2 -dual under $i \leftrightarrow j$.

6.2 Isotropy coefficient $c_{\text{lat}} = 6$

Proposition 6.1. *The continuum limit of (17) is*

$$S_{\text{lat}} \longrightarrow \int d^4x \frac{c_{\text{lat}}}{4} M_P^2 (\partial_\mu \sigma)^2, \quad c_{\text{lat}} = \sum_{k=1}^{24} (u_k^0)^2 = 2 + 4 = 6. \quad (18)$$

Cross-check: $4c_{\text{lat}} = \sum_{k=1}^{24} N(u_k) = 24$. \checkmark

The computation: 8 axis-aligned unit elements $\pm e_\mu$ contribute $(u^0)^2 = 1$ for $\pm e_0$ (2 terms, total 2) and 0 for $\pm e_{1,2,3}$; the 16 half-integer elements $\frac{1}{2}(\pm 1 \pm \mathbf{i} \pm \mathbf{j} \pm \mathbf{k})$ each contribute $(1/2)^2 = 1/4$ (total $16 \times 1/4 = 4$). Hence $c_{\text{lat}} = 2 + 4 = 6$.

6.3 Potential and effective action

The \mathbb{Z}_2 symmetry $\sigma \mapsto -\sigma$ restricts V to even functions of σ . The duality-symmetric potential is

$$V(\sigma) = 2\rho_\Lambda \cosh \sigma = \rho_\Lambda e^\sigma + \rho_\Lambda e^{-\sigma}, \quad (19)$$

where e^σ is the UV contribution and $e^{-\sigma}$ the IR contribution, equal at the self-dual point $\sigma = 0$. This is the unique \mathbb{Z}_2 -even potential of this exponential form [7]. Its minimum is at $\sigma = 0$ with $V''(0) = 2\rho_\Lambda > 0$.

The potential (19) follows from the lattice partition function $Z = \sum_N r_4(N) \exp(-\beta \cdot 2\rho_\Lambda \cosh \sigma_N)$ by the saddle-point condition. The relevant asymptotic of Jacobi's four-square formula is

$$r_4(N) \sim \pi^2 N \quad (N \rightarrow \infty), \quad (20)$$

confirmed numerically to 0.05%.

Remark 6.2 (Correction to earlier versions). *Earlier drafts incorrectly cited $r_4(N) \sim 8\pi^2 N$. The correct average asymptotic is $r_4(N) \sim \pi^2 N$, obtained from $\sum_{n=1}^N r_4(n) \sim \pi^2 N^2/2$ (lattice-point count in a 4D ball of radius \sqrt{N}). This correction does not affect any quantitative prediction of the paper.*

The effective action is

$$S = \int d^4x \sqrt{-g} \left[\frac{M_P^2}{2} R - \frac{3}{2} M_P^2 (\partial_\mu \sigma)^2 - \frac{1}{2g_w^2} \text{Tr}(F_{\mu\nu} F^{\mu\nu}) - V(\sigma) \right]. \quad (21)$$

The kinetic coefficient $\frac{3}{2} M_P^2$ follows from $c_{\text{lat}} = 6$. The gravitational sector $M_P^2 R/2$ is taken from the standard Einstein–Hilbert action; its derivation from the Hurwitz lattice is not addressed here.

7 Dark Energy Predictions

7.1 Equations of motion

Variation of (21) with respect to $g^{\mu\nu}$ and σ gives (Appendix B):

$$3H^2 = \kappa \left[\rho_m + \frac{3}{2} \dot{\sigma}^2 + 2\rho_\Lambda \cosh \sigma \right], \quad (22)$$

$$\ddot{\sigma} + 3H\dot{\sigma} = -\frac{2\rho_\Lambda}{3} \sinh \sigma. \quad (23)$$

The dark-energy equation of state is

$$w = \frac{\frac{3}{2} \dot{\sigma}^2 - 2\rho_\Lambda \cosh \sigma}{\frac{3}{2} \dot{\sigma}^2 + 2\rho_\Lambda \cosh \sigma}. \quad (24)$$

The coefficients $3/2$ and $2\rho_\Lambda/3$ originate from $c_{\text{lat}} = 6$.

7.2 Numerical results

Remark 7.1 (Normalisation). *The system (22)–(23) is integrated in dimensionless units with $\rho_{\text{crit},0} \equiv 1$, $\Omega_{m,0} = 0.315$, and $\rho_\Lambda = \Omega_{\Lambda,\text{obs}}/2 = 0.3425$.*

Remark 7.2 (w_a estimation method). *The CPL parameter w_a is extracted using the two-point formula $w_a \approx 2[w(z=1) - w_0]$, computed from dense-output numerical integration at $a = 0.5$ ($z = 1$). All w_a values are negative (thawing behaviour), and $w_0 + w_a \approx -1$.*

Table 1: CPL dark-energy parameters from numerical integration of (22)–(23) with $c_{\text{lat}} = 6$. No entry falls within the DESI DR2 (2025) 1σ range $w_0 \in [-0.79, -0.11]$.

σ_i	w_0	w_a
0.5	-0.990	-0.014
1.0	-0.968	-0.044
1.5	-0.947	-0.069
2.0	-0.930	-0.082
3.0	-0.908	-0.076

Table 2: Redshift evolution $w(z)$ for the lattice-derived action.

σ_i	$w(z=0)$	$w(z=0.5)$	$w(z=1)$	$w(z=2)$	$w(z=3)$
0.5	-0.990	-0.995	-0.997	-0.999	-1.000
1.0	-0.968	-0.982	-0.990	-0.997	-0.999
1.5	-0.947	-0.967	-0.981	-0.993	-0.997
2.0	-0.930	-0.954	-0.971	-0.988	-0.995
3.0	-0.908	-0.928	-0.947	-0.973	-0.986

7.3 Physical meaning of σ_i

The scale field $\sigma \approx 2 \log(r/r_c)$ is dimensionless, with $\sigma = 0$ at r_c . In the photon-basis representation, $\sigma_i = \log(T_{\text{init}}/T_c)$ where $T_{\text{init}} = T_c e^{\sigma_i}$ is the CMB temperature at the epoch when Hubble friction froze σ . Numerically, with $T_c = 27.27$ K: the entries in Table 3 follow directly (e.g. $\sigma_i = 1.0 \Rightarrow T_{\text{init}} = 27.27 e^{1.0} = 74.1$ K). The field began rolling toward $\sigma = 0$ only after H fell below the effective scalar mass $m_\sigma = \sqrt{2\rho_\Lambda/3}$. The determination of σ_i from first principles is Open Problem OP2.

7.4 Observational comparison

7.4.1 Current status of w_0 and w_a

DESI DR2 (2025; BAO combined with CMB and SNIa [9]) reports dynamical dark energy at 2.8 – 4.2σ significance, with favoured parameters $w_0 \approx -0.75$ and $w_a \approx -0.60$ (DESI+CMB+DESY5). The QNRT predictions (Table 1) give $w_0 \in [-0.990, -0.908]$ for $\sigma_i \in [0.5, 3.0]$. These values lie outside the DESI DR2 1σ range: the predictions are 2.7 – 4.0σ from the DESI central value in w_0 and 2.5 – 2.9σ in w_a . QNRT therefore predicts a thawing quintessence signal [8] that is consistent with Λ CDM but inconsistent with the current DESI central value. Both outcomes are observationally viable pending Euclid.

7.4.2 Falsifiability conditions

(a) DESI confirmed at high significance, $w_0 \approx -0.75$: strongly disfavours the present QNRT framework. Extensions would be required.

(b) Future surveys find w_0 consistent with -1 : Λ CDM remains viable; QNRT with any $\sigma_i \in [0.5, 3]$ is consistent.

(c) $w_a > 0$ measured at high significance: freezing-quintessence behaviour, incompatible with the thawing dynamics of QNRT; this would falsify the model.

Table 3: Physical scales corresponding to σ_i .

σ_i	r_i/r_c	T_{init} (K)	z_{init}	σ_i/S
0.5	1.28	45.0	15.5	0.0035
1.0	1.65	74.1	26.2	0.0071
2.0	2.72	201.5	72.9	0.014
3.0	4.48	547.8	200.0	0.021

(d) Euclid joint (w_0, w_a) measurement: QNRT predicts a one-parameter curve in the (w_0, w_a) plane. Euclid’s forecasted precision $\sigma(w_0) \sim 0.02$ [10] would distinguish $w_0 = -0.95$ from $w_0 = -1$ at $\sim 2.5\sigma$.

7.4.3 Sensitivity to the kinetic coefficient

Table 4: Dark-energy predictions for lattice-derived and Brans–Dicke kinetic coefficients.

σ_i	Lattice ($3M_P^2/2$)		Brans–Dicke ($M_P^2/16\pi$)	
	w_0	w_a	w_0	w_a
0.5	−0.990	−0.014	−1.000	−0.000
1.0	−0.968	−0.044	−1.000	−0.000
2.0	−0.930	−0.082	−1.000	−0.001
3.0	−0.908	−0.076	−0.999	−0.000

8 Neutrino Mass Estimate

The norm-response mechanism applied to the fermionic sector gives a reference mass M_c as the geometric mean of the Planck and Hubble mass scales:

$$M_c c^2 = \sqrt{M_P c^2 \cdot M_H c^2}, \quad M_P = \sqrt{\hbar c/G}, \quad M_H = \frac{\hbar}{R_\star c}, \quad (25)$$

giving, with $R_\star \approx 4.36 \times 10^{26}$ m:

$$M_c c^2 \approx 2.4 \text{ meV}. \quad (26)$$

This lies within the active-neutrino mass window ($\lesssim 50$ meV [11, 13]) and below the KATRIN upper bound (< 450 meV [12]).

Remark 8.1. *Earlier versions quoted $M_c c^2 \approx 4.2$ meV, obtained by inserting $R_\star = R_H \approx 1.37 \times 10^{26}$ m (Hubble radius) into (25). The present version uses $R_\star \approx 4.36 \times 10^{26}$ m consistently throughout, giving $M_c c^2 \approx 2.4$ meV. Both values lie within the neutrino mass window; the qualitative conclusion is unchanged. Equation (25) is a dimensional estimate, not a precision prediction. It does not yield individual mass eigenvalues or the mass ordering; precision predictions require the gauge embedding (OP3) and the generation structure (Part II).*

9 Connection to Matter Fields

9.1 The eight degrees of freedom of $\mathbb{H}_{\mathbb{C}}$

The algebra $\mathbb{H}_{\mathbb{C}} \cong M_2(\mathbb{C})$ has real dimension 8. Every invertible $\phi \in M_2(\mathbb{C})$ admits a unique polar decomposition:

Proposition 9.1 (Polar decomposition of $\mathbb{H}_{\mathbb{C}}$).

$$\phi = P \cdot U, \quad (27)$$

where $P = \sqrt{\phi\phi^\dagger}$ is positive-definite Hermitian and $U = P^{-1}\phi$ is unitary, giving the splitting

$$P_{4 \text{ real DOF}} = \exp\left(\frac{\sigma}{2}I_2 + \frac{\xi_a\tau^a}{2}\right), \quad U_{4 \text{ real DOF}} \in U(2) = U(1) \times \text{SU}(2), \quad (28)$$

where τ^a ($a = 1, 2, 3$) are the Pauli matrices and $\xi_a \in \mathbb{R}$.

Table 5: Eight real degrees of freedom of $\mathbb{H}_{\mathbb{C}} \cong M_2(\mathbb{C})$.

DOF	Field	Type	Role in QNRT
σ	scale field	boson	dark energy (established)
ξ_1, ξ_2, ξ_3	traceless Hermitian	boson	unidentified; deferred
$\theta \in U(1)$	complex phase	boson	$U(1)_Y$ candidate
$u \in \text{SU}(2)$	directional field	boson	$\text{SU}(2)$ gauge (established)

The current paper uses only σ and u (4 of the 8 real DOF). The remaining 4 DOF (θ and ξ_a) are present in the algebra and are identified as subjects for future work.

9.2 What is established

SU(2) gauge field. The decomposition $\phi = e^{\sigma/2}u$, $u \in \text{SU}(2)$, separates the scale sector σ from the directional sector u . The kinetic term for u in the lattice action (17) gives the SU(2) field strength $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu - ig_w[A_\mu, A_\nu]$ in the continuum limit.

Weyl spinors and chirality. Since $\mathbb{H}_{\mathbb{C}} \cong M_2(\mathbb{C})$ acts on \mathbb{C}^2 by left matrix multiplication, any $\Psi \in \mathbb{C}^2$ is a Weyl spinor on which $u \in \text{SU}(2)$ acts as $\Psi \mapsto u\Psi$. This action automatically projects onto a definite chirality.

9.3 What remains to be established

The following connections are not claimed in this paper.

- (i) **Identification of lattice SU(2) with $\text{SU}(2)_L$.**
- (ii) **Physical role of the traceless Hermitian fields ξ_a .**
- (iii) **Hypercharge $U(1)_Y$.**

- (iv) **Colour gauge group** $SU(3)_c$. The colour group is not contained in $\mathbb{H}_\mathbb{C} \cong M_2(\mathbb{C})$. The F_4 root lattice contains A_2 sub-root systems whose Weyl groups are isomorphic to the $SU(3)$ Weyl group; this is explored in Part II. Note: an earlier version of this paper suggested that $G_2 \supset SU(3)$ could serve as the lattice automorphism group for colour. Explicit calculation shows that the Weyl group $W(G_2)$ generates components $\pm 1/3, \pm 2/3$ not present in Γ_H ; this route is therefore ruled out.
- (v) **Anomaly cancellation**. This requires a full Standard Model embedding including the colour sector and correct hypercharge assignments, none of which is established in this paper. Note: an earlier version suggested an E_8 -based context for anomaly cancellation; this is withdrawn. The E_8 extension is blocked by the non-associativity of the Cayley integers and gives $c_{\text{lat}}(E_8) = 60 \neq 6$, inconsistent with the lattice action derived in Section 6.
- (vi) **Three fermion generations**. The D_4 sub-root system of F_4 has a triality automorphism of order 3, a natural candidate for generating three generations. The concrete field identification is deferred to Part II.
- (vii) **Fine-structure constant** α . The value $\alpha = 1/137$ is not derived from the Hurwitz lattice in this paper; its lattice origin is part of OP3.

10 Open Problems

OP1. Physical identification and determination of R_\star . This is the central unsolved problem of QNRT. All physical outputs are conditioned on $R_\star \approx 4.4 \times 10^{26}$ m ($N_\star \approx 2.70 \times 10^{61}$), set by self-consistency of the numerical outputs (Remark 5.1). This value is approximately 3.2 times the Hubble radius $R_H = c/H_0 \approx 1.37 \times 10^{26}$ m and close to the comoving particle horizon (≈ 46 Gly). Three approaches to deriving R_\star from first principles have been investigated: (i) de Sitter attractor gives $R_\star \approx 0.24$ m (sub-metre); (ii) dynamical identification $R_\star(t) = R_H(t)$ gives $\Omega_{DE} \approx 0.27$ (inconsistent with observations); (iii) lattice-compatibility gives only trivial constraints. None gives a complete resolution. Via $\rho_\Lambda = c^4/(GR_\star^2)$, OP1 is a purely classical-gravitational question.

OP2. Physical determination of σ_i . The initial displacement σ_i is a free parameter. Determination requires analysing QNRT dynamics during inflation or reheating.

OP3. Standard Model gauge embedding. Identify the lattice $SU(2)$ with $SU(2)_L$; provide $SU(3)_c$ (likely via the A_2 sub-root system of F_4 , investigated in Part II); verify anomaly cancellation; derive $\alpha = 1/137$ from lattice structure.

OP4. Kinetic coefficient. The Voronoi–Regge correspondence on the F_4 lattice confirms $\alpha_{\text{lat}} = c_{\text{lat}} M_P^2/4 = (3/2)M_P^2$. The coefficient is understood geometrically; it is not an undetermined free parameter.

OP5. Connection to CMB anisotropies. Computing the connection of the σ field to CMB observables requires perturbation theory for spatial fluctuations $\delta\sigma(x, t)$, reserved for Part II.

11 Conclusion

QNRT derives observable consequences through a single algebraic operation: $\phi \mapsto \phi^{-1} = \phi^\dagger/N(\phi)$ on the Hurwitz lattice. This revised paper (v3.3) incorporates eight corrections

(C1–C8) to earlier versions:

- (C1–C4) Algebraic corrections (v3.2): HC-based Minkowski metric; $r_4(N) \sim \pi^2 N$; withdrawal of $G_2 \supset \text{SU}(3)$ automorphism; withdrawal of E_8 anomaly-cancellation claim.
- (C5) $R_\star \approx 4.4 \times 10^{26}$ m (not $R_H \approx 1.37 \times 10^{26}$ m): all numerical outputs are consistent only with the larger value.
- (C6) $\rho_\Lambda \approx 6.3 \times 10^{-10} \text{ J m}^{-3}$, within 21% of observed (not $7 \times 10^{-9} \text{ J m}^{-3}$, “one order of magnitude”).
- (C7) $\rho_\Lambda^{\text{QNRT}}/\rho_\Lambda^{\text{obs}} = 1.209$ is a 20.9% excess (not “12%”).
- (C8) $M_c c^2 \approx 2.4 \text{ meV}$ with $R_\star = 4.4 \times 10^{26}$ m (revised from 4.2 meV with $R_\star = R_H$; both within neutrino window).

The QNRT prediction $\rho_\Lambda \approx 6.3 \times 10^{-10} \text{ J m}^{-3}$ is within 21% of the Planck 2018 value. The dark-energy predictions $w_0 \in [-0.990, -0.908]$ are inconsistent with DESI DR2 (2025) at 2.7–4.0 σ , while remaining consistent with ΛCDM . If Euclid measures $w_0 \in [-0.990, -0.908]$ at high significance, QNRT is supported; if $w_0 \approx -0.75$ is confirmed, the framework is challenged. If $w_a > 0$ is established at high significance, the model is falsified.

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A Continuum Limit of the Lattice Action

Step 1: Taylor expansion. Let $\phi_i = \phi(x)$ and $\phi_j = \phi(x + a\hat{e}_k)$ with $a = \ell_P$. Using $\phi = e^{\sigma/2}u$:

$$\frac{|\phi_i - \phi_j|^2}{|\phi_i|^2} \approx a^2 \left[\frac{1}{4} (\hat{e}_k^\mu \partial_\mu \sigma)^2 + N(u^{-1} \hat{e}_k^\mu \partial_\mu u) \right] + O(a^4). \quad (29)$$

Step 2: Sum over 24 neighbours. By F_4 symmetry, $\sum_k u_k^\mu u_k^\nu = c_{\text{lat}} \delta^{\mu\nu}$ with $c_{\text{lat}} = T_{00} = \sum_k (u_k^0)^2$. Computing from the 24 unit elements: the 8 axis-aligned units $\pm e_\mu$ contribute $(u^0)^2 = 1$ for $\pm e_0$ (total 2) and 0 for $\pm e_{1,2,3}$; the 16 half-integer units $\frac{1}{2}(\pm 1 \pm \mathbf{i} \pm \mathbf{j} \pm \mathbf{k})$ each contribute $(1/2)^2 = 1/4$ (total 4). Hence $c_{\text{lat}} = 2 + 4 = 6$.

Step 3: Continuum integral. Converting $\sum_{\text{sites}} \rightarrow (1/a^4) \int d^4x$ with $a = \ell_P = 1/M_P$:

$$S_{\text{lat}} \longrightarrow \int d^4x \frac{c_{\text{lat}}}{4} M_P^2 (\partial_\mu \sigma)^2 = \int d^4x \frac{3}{2} M_P^2 (\partial_\mu \sigma)^2. \quad (30)$$

B Derivation of the FLRW Equations

Starting from the effective action (21) with the SU(2) sector constant on cosmological scales:

Einstein equations. $T_{\mu\nu}^\sigma = 3M_P^2\partial_\mu\sigma\partial_\nu\sigma - g_{\mu\nu}[\frac{3}{2}M_P^2(\partial\sigma)^2 + V(\sigma)]$. In flat FLRW with $\sigma = \sigma(t)$: $\rho_\sigma = \frac{3}{2}M_P^2\dot{\sigma}^2 + V(\sigma)$, $p_\sigma = \frac{3}{2}M_P^2\dot{\sigma}^2 - V(\sigma)$. The (00) component gives (22) with $\kappa = 1/M_P^2$.

Klein–Gordon equation. Varying with respect to σ :

$$-V'(\sigma) + 3M_P^2\Box\sigma = 0, \quad \Box\sigma = -\ddot{\sigma} - 3H\dot{\sigma}, \quad V'(\sigma) = 2\rho_\Lambda \sinh \sigma, \quad (31)$$

which yields (23).

C Supplementary Analysis of the R_\star Problem

C.1 Entropy maximisation and the lattice partition function

The Hurwitz lattice partition function

$$Z = \sum_{N=1}^{N_\star} r_4(N) \exp(-\beta \cdot 2\rho_\Lambda \cosh \sigma_N), \quad (32)$$

with $r_4(N) = 8 \sum_{d|N, 4 \nmid d} d$, converges for any $\beta\rho_\Lambda > 0$. The saddle point occurs near $N = n_c$ when $\beta\rho_\Lambda \sim 1$. Setting $T_{\text{lattice}} = T_{dS} = \hbar H / (2\pi k_B)$ and substituting $\rho_\Lambda = \hbar c / (\ell_P R_\star)^2$ gives $R_\star \sim 2\pi\ell_P$ (Planck-scale). Simple entropy maximisation does not select the observed value $R_\star \approx 4.4 \times 10^{26}$ m.

C.2 Summary: current understanding of the R_\star problem

- (i) No dynamical argument within the current framework selects $R_\star \approx 4.4 \times 10^{26}$ m. The value is fixed by self-consistency of the numerical outputs (Remark 5.1); it is notably close to the comoving particle horizon (≈ 46.5 Gly) but this coincidence is not derived.
- (ii) The partition function Z converges—a non-trivial algebraic property.
- (iii) R_\star is frozen by Hubble friction during matter domination, exactly as σ is frozen at σ_i .
- (iv) OP1 and OP2 are connected: both R_\star and σ_i are frozen at the same epoch. A unified treatment of QNRT during inflation and reheating would simultaneously determine both. This is the central goal of Part II.
- (v) The value $R_\star \approx 4.4 \times 10^{26}$ m is the working hypothesis. Via $\rho_\Lambda = c^4 / (GR_\star^2)$, OP1 asks why the Universe has the observed dark-energy density today. This is a classical gravitational question; no quantum input (\hbar) is required.