

Uniform Coercivity, Pointwise Large-Field Suppression, and Conditional Closure of the Lattice Yang–Mills Mass Gap at Weak Coupling in $d = 4$

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Abstract

We address the remaining interface gaps in the programme [E26I]–[E26VIII] toward a uniform log-Sobolev inequality (LSI) and transfer-matrix spectral gap for lattice $SU(N_c)$ Yang–Mills in $d = 4$ at weak coupling. Four gaps are treated: (G1) the pointwise-in-background validity of Balaban’s T -operation small-factor bound — stated in v2 as the explicit hypothesis (H-PTW), since the detailed audit appendices announced in v1 were absent from the document (their references rendered as “??”); (G2) a uniform small-field coercivity estimate for the effective action; (G3) uniform analyticity of boundary terms; (G4) a quantitative bootstrap of all constants. **The central correction of v2:** the sign of the one-step coupling drift in v1’s Theorem 5.2 ($g_{k+1}^{-2} = g_k^{-2} + 2b_0 \ln L_{\text{RG}}$, coupling weakening toward the infrared) is inverted relative to asymptotic freedom, and contradicts the companions’ own use of $n_{\text{max}} \sim 1/(2b_0 g^2 \ln 2)$ (ai.viXra:2602.0032 §8, 2602.0033, 2602.0041 §7.3), a formula meaningful only if the coupling *grows* along blocking and exits the weak regime. With the corrected sign the monotone bootstrap of v1’s Theorem 5.3 reverses: the inductive conditions are guaranteed only up to the finite horizon $k^*(\beta) = (g_0^{-2} - \gamma_0^{-2})/(2b_0 \ln L_{\text{RG}}) + O(1)$, and since the multiscale construction uses $\log_2 L$ scales, the conclusion holds *on the volume window* $\log_2 L \leq k^*(\beta)$, i.e. $L \leq e^{C/g^2 + O(1)}$ with $C = 1/(2b_0) = 24\pi^2/(11N_c)$ — exactly the window of the companion papers ai.viXra:2602.0032/0033 (v2). Full volume-uniformity would additionally require a strong-coupling handoff beyond the crossover scale (Osterwalder–Seiler regime), stated as hypothesis (H-XOVER) and not established here; accordingly the “unconditional closure” of v1 is retitled to *conditional closure*. Version 2 also repairs the dangling “??” references, supplies the missing proof of Lemma 3.1, rewrites the proof of Lemma 7.4 in the series’ fundamental-trace convention (its statement $W''(0) = 1/(2N_c)$ is correct; the v1 proof mixed normalized and fundamental traces), corrects $\sum_{k \geq 0} (k+1)2^{-3k} = 64/49$ (v1: $8/49$), and fills in the companion identifiers. All corrections are verified in a companion numerical suite.

1 Introduction

1.1 Context and objective

On $\Lambda = (\mathbb{Z}/L\mathbb{Z})^d$ with Wilson action

$$A(U) := \sum_{p \in \mathcal{P}(\Lambda)} \left(1 - \frac{1}{N_c} \operatorname{Re} \operatorname{Tr} U_p\right), \quad S_\beta = \beta A(U), \quad \beta = \frac{2N_c}{g^2}, \quad (1)$$

the mass gap is $\Delta_{\text{phys}}(\beta, L) = -\log \|\widehat{T}|_{1^\perp}\|$. The series [E26I]–[E26VIII] establishes the route Uniform LSI $\Rightarrow$ DLR-LSI $\Rightarrow$ SZ mixing $\Rightarrow$ RP $\Rightarrow \Delta_{\text{phys}} > 0$. This paper treats the four interface gaps (G1)–(G4); v2 states honestly which are closed and which become explicit hypotheses.

1.2 The four gaps

(G1) pointwise large-field suppression (v2: hypothesis (H-PTW), §3); (G2) fiber-level control without exponential collapse (§2); (G3) uniform analyticity of boundary terms (§4); (G4) positivity of α_* and simultaneous compatibility of constants (§5), now on the volume window of Theorem 1.1.

1.3 Main result

Theorem 1.1 (Mass gap, lattice, $d = 4$, weak coupling; windowed and conditional). *For $G = \text{SU}(N_c)$, $d = 4$, there exists $\beta_0 = \beta_0(N_c) < \infty$ such that for all $\beta \geq \beta_0$ and all even $L \geq 2$ with*

$$\log_2 L \leq k^*(\beta) := \frac{g_0^{-2} - \gamma_0^{-2}}{2b_0 \ln L_{\text{RG}}} + O(1) \quad (\text{equivalently } L \leq e^{C/g^2 + O(1)}, \quad C = \frac{1}{2b_0} = \frac{24\pi^2}{11N_c}),$$

and conditionally on (H-PTW) (§3) and the companion inputs [E26I]–[E26VIII]:

$$\operatorname{Ent}_{\mu_\beta}(f^2) \leq \frac{2}{\alpha_*} \sum_{e \in E(\Lambda)} \int |\nabla_e f|^2 d\mu_\beta, \quad \alpha_* = \alpha_*(N_c, \beta_0) > 0, \quad (2)$$

and

$$\Delta_{\text{phys}}(\beta, L) \geq c(N_c, \beta_0) > 0. \quad (3)$$

Removing the volume window would additionally require hypothesis (H-XOVER) below. (v1 asserted (2)–(3) unconditionally and uniformly in all L ; see Remark 5.4 for why the corrected coupling flow forces the window.)

(H-XOVER) Strong-coupling handoff. *Beyond the crossover scale $k^*(\beta)$ the effective theories are at $O(1)$ or strong coupling; volume-uniformity for all L requires establishing the conditional inequalities of the multiscale decomposition at those scales by strong-coupling methods (Osterwalder–Seiler cluster expansion) and matching them to the weak-coupling regime. This is not carried out in this paper or, to our knowledge, in the companions.*

1.4 Notation and conventions

As in v1: $\langle X, Y \rangle = -2 \operatorname{Tr}(XY)$ with Tr the *fundamental* (unnormalized) trace ($\operatorname{Tr} \mathbf{1} = N_c$) — v2 note: this convention is used consistently, repairing the mixed usage in v1’s §7.2; $\beta = 2N_c/g^2$; $p_0(g) \geq c_0 |\log g|^{1+\epsilon_0}$; $\varepsilon_k = c_* g_k^{1-\delta}$; $R^{(k)}, B^{(k)}, T_k$ as in [Bal89b].

2 Small-field coercivity

As in v1 (this section survives the audit unchanged in substance). Setup: conditional fiber measure

$$d\mu_k^{\bar{U}}(U^{\text{fast}}) = \frac{1}{Z_k(\bar{U})} \exp(-S_{\text{eff},k}(U^{\text{fast}}, \bar{U})) \prod_{e \in E_k(B)} dU_e. \quad (4)$$

Definition 2.1. $\Omega_k^{\text{sf}}(B; \bar{U})$: all plaquettes satisfy $\|U_p - 1\|_{\text{HS}} < \varepsilon_k$.

Definition 2.2. Exponential chart around a minimiser $U_e^{(0)}(\bar{U})$.

Assumption 2.3 (Balaban one-step inputs (published)). (B1) one-step RG map; (B2) polymer bounds $|R^{(k)}(X)| \leq e^{-p_0(g_k)|X|} e^{-\kappa d_k(X)}$; (B3) analyticity with radius $\hat{\alpha}(\gamma)$ independent of k ; (B4) T -operation small-factor bound $T_k(Y)1 \leq \exp(-c_{\text{sf}} p_0(g_k)|Y|)$. (v2: the phrase “Pointwise-in-background validity: ??” of v1 is replaced by the explicit hypothesis (H-PTW) of §3; the running-coupling control is derived in §5/§7 with the corrected sign, on the finite horizon.)

Remark 2.4. As in v1: coercivity is used only in the inductive regime; large-field configurations are handled by the T -operation mechanism.

Proposition 2.5 (Riemannian Hessian bound). In $\Omega_k^{\text{sf}}(B; \bar{U})$, under the inductive hypotheses at scale k and $g_k \leq \gamma_0$:

$$\text{Hess}_{\text{Riem}}(V_{k,B})(A, A) \geq \frac{\beta_k \gamma_{\text{Bal}}}{4} \|A\|^2, \quad (5)$$

with the proof as in v1 (quadratic part from (IC6); cubic and curvature corrections $O(g_k^{1-\delta}) + O(\beta_k^{-1})$ relative; polymer contribution $O(\beta_k^{-1})$ relative; thresholds (T4)–(T5)).

Remark 2.6. As in v1: (5) is a convexity input for the sweeping-out estimates, not a global fiber LSI.

Proposition 2.7 (Conditional LSI, Holley–Stroock, honest version). For every $\beta \geq \beta_0$ and every $\bar{U}$, the fiber measure satisfies LSI(α_0) with

$$\alpha_0 := \alpha_{\text{Haar}} \exp(-2 \text{osc}(V_{k,B})) > 0, \quad (6)$$

L -independent, possibly exponentially small in β (as in v1, Remark 2.8).

Remark 2.8. Exponentially small α_0 is acceptable for the existence statement $\Delta_{\text{phys}} > 0$.

3 Large-field suppression: the (H-PTW) hypothesis

v1 claimed to resolve (G1) by “auditing Balaban’s proof”, with the detailed audit in appendices whose references rendered as “??” — the announced audit appendices are absent from the v1 document. v2 therefore states the required property as an explicit hypothesis, keeping v1’s three-point supporting sketch as a remark.

Lemma 3.1 (RCP as fibre integral). *There is a version of the regular conditional probability given for every $\bar{U}$ by the ratio of fibre integrals*

$$\mu_k(A \mid \mathcal{G}_{k+1})(\bar{U}) = \frac{\int_A e^{-S_{\text{eff},k}(U_{\text{fast}}, \bar{U})} \prod dU_e}{\int e^{-S_{\text{eff},k}(U_{\text{fast}}, \bar{U})} \prod dU_e}, \quad (7)$$

well-defined and continuous in $\bar{U}$.

Proof. (Supplied in v2; v1's proof reference rendered as “??”.) The integrand is continuous and strictly positive on the compact fibre $G^{|E_k(B)|}$, so both integrals are finite, the denominator is bounded below by $e^{-\sup S_{\text{eff},k}} > 0$, and continuity in $\bar{U}$ follows from dominated convergence (the action is continuous in $\bar{U}$ on compacts). The resulting kernel integrates correctly against $\mathcal{G}_{k+1}$ -measurable functions by Fubini, hence is a version of the conditional probability. $\square$

Corollary 3.2. *Pointwise bounds in $\bar{U}$ imply essential-supremum bounds.*

Theorem 3.3 (Uniform large-field suppression; conditional on (H-PTW)). *Assume (H-PTW) below. For $\beta \geq \beta_0$, $0 \leq k \leq n_{\text{max}}$, every block B and every $\bar{U}$:*

$$\mu_k(Z_k(B) \mid \mathcal{G}_{k+1} = \bar{U}) \leq C_{\text{blk}} e^{-c_{\text{lf}} p_0(g_k)}. \quad (8)$$

(H-PTW) Pointwise small factor. *For every real background $\bar{U}$, every connected large-field component Y , and every k : $T_k(Y)1(\bar{U}) \leq \exp(-c_{\text{sf}} p_0(g_k)|Y|)$, i.e. Balaban's small-factor estimate [Bal89b, eq. (1.89)] holds pointwise in $\bar{U}$ and not only for admissible backgrounds.*

Proposition 3.4 (Supporting sketch for (H-PTW); from v1). *Three observations support (H-PTW): (i) the T -operation is a ratio of strictly positive fibre integrals (Lemma 3.1), defined for every real $\bar{U}$, with the $O(\beta_k)$ offset cancelling in the ratio; (ii) the estimates in Balaban's derivation of (1.89) use only uniform inductive bounds whose constants do not reference the particular $\bar{U}$; (iii) the admissibility/analyticity-domain restrictions concern the complexified continuation used for the next RG step, not the real-valued ratio. (v2: this is v1's Proposition 3.4 sketch; the “detailed audit” it deferred to was not contained in the v1 document, so we do not claim (i)–(iii) constitute a proof. Discharging (H-PTW) requires writing out that audit.)*

Given (8)'s ingredients, the lattice-animal summation is as in v1:

$$\mu_k(Z_k(B) \mid \bar{U}) \leq \sum_{n \geq 1} C_d^n e^{-c_{\text{sf}} p_0(g_k)n} \leq 2C_d e^{-c_{\text{sf}} p_0(g_k)} \quad (9)$$

under (T6) ($C_d e^{-c_{\text{sf}} p_0} \leq \frac{1}{2}$; counting bound verified by brute force in the suite).

4 Uniform analyticity of boundary terms

Proposition 4.1 (Analyticity of $B^{(k)}$). *As in v1: $B^{(k)}(X)$ is analytic with radius $\hat{\alpha}_1(\gamma)$ independent of k , with*

$$\sup_{|z| \leq \hat{\alpha}_1} |B^{(k)}(X; U + z)| \leq C_B(k+1)|X|e^{-\kappa d_k(X)}, \quad |\nabla B^{(k)}(X)| \leq \frac{C_B(k+1)|X|}{\hat{\alpha}_1} e^{-\kappa d_k(X)}. \quad (10)$$

The proof (inductive conditions, cluster-expansion localisation, sup-norm bound from [Bal89b] eq. (1.69), Cauchy estimate) is as in v1. Note $(k+1) \leq n_{\text{max}} + 1$, which on the window of Theorem 1.1 is $O(k^(\beta))$, a β -dependent but L -independent constant.*

5 Quantitative bootstrap (corrected)

5.1 Control of the running coupling: corrected sign and finite horizon

Definition 5.1 (Wilson projection and running coupling). *As in v1, with the v2-corrected normalization of Lemma 7.4: writing $S_k^{\text{loc}} = \beta_k A(U) + S_k^{\text{loc,irr}}$,*

$$\beta_k := \frac{1}{W''(0)} \frac{d^2}{dt^2} \Big|_{t=0} S_k^{\text{loc}}(U_p = \exp(tT)) = 2N_c \frac{d^2}{dt^2} \Big|_{t=0} S_k^{\text{loc}}, \quad g_k^{-2} := \frac{\beta_k}{2N_c}. \quad (11)$$

Theorem 5.2 (One-step coupling renormalization; sign corrected in v2). *There exist $\gamma_{\text{run}} > 0$, $C_{\text{run}} < \infty$ such that if (IC1)–(IC6) hold at scale k with $g_k \leq \gamma_{\text{run}}$:*

$$g_{k+1}^{-2} = g_k^{-2} - 2b_0 \ln L_{\text{RG}} + \delta_k, \quad |\delta_k| \leq C_{\text{run}} g_k^2, \quad b_0 = \frac{11N_c}{48\pi^2} : \quad (12)$$

the coupling grows toward the infrared (asymptotic freedom). (v1 stated (12) with $+2b_0 \ln L_{\text{RG}}$; Remark 5.4.)

Theorem 5.3 (Bootstrap with finite horizon; corrected in v2). *Let γ_0 satisfy (T1)–(T7). If $g_0 \leq \gamma_0/\sqrt{2}$ (say), then g_k is (weakly) increasing and*

$$g_k \leq \gamma_0 \quad \text{for all } 0 \leq k \leq k^*(\beta), \quad k^*(\beta) \geq \frac{g_0^{-2} - \gamma_0^{-2}}{2b_0 \ln L_{\text{RG}} + C_{\text{run}} \gamma_0^2} = \frac{C}{g^2} (1 + O(1)), \quad C = \frac{1}{2b_0},$$

by iterating (12) downward from g_0^{-2} ; under (T3) ($C_{\text{run}} \gamma_0^2 \leq b_0 \ln L_{\text{RG}}$) the guaranteed horizon is at least half the one-loop value (numerically exhibited in the suite, including the noisy recursion). (v1's Theorem 5.3 claimed the monotone bound $g_{k+1} \leq g_k \leq \gamma_0$ for all k — true only with the inverted sign, under which the coupling never exits the weak regime and the “global small-coupling hypothesis” appears closable unconditionally. With the correct sign the property is a finite-horizon statement, and since the multiscale construction uses $n_{\text{max}} = \log_2 L$ scales, the constructions of §§2–4 are validated only for $\log_2 L \leq k^(\beta)$: the volume window of Theorem 1.1.)*

Remark 5.4 (The sign story; new in v2). Three independent witnesses fix the sign: (i) asymptotic freedom — the effective coupling of a nonabelian gauge theory grows at coarser scales; v1's Table 1 (“ β_k increasing”) describes infrared freedom instead; (ii) the companions' own usage: $n_{\text{max}} \sim 1/(2b_0 g^2 \ln 2)$ (2602.0032 §8, 2602.0033, 2602.0041 §7.3) is the exit time of the growing coupling, meaningless under v1's sign (the suite exhibits that with the + sign the equation $g_k = \gamma_0$ has no solution); (iii) the corrected flow reproduces exactly the volume window $L \leq e^{C/g^2 + O(1)}$, $C = 24\pi^2/(11N_c)$, of the conditional companions 2602.0032/0033 (v2) — the series becomes mutually consistent. The windowed Theorem 1.1 is therefore not a loss peculiar to this paper: it is the honest common boundary of the programme, and passing it requires (H-XOVER).

Definition 5.5 (Master threshold). $\gamma_0 > 0$ *largest satisfying (T1)–(T7) as in v1 (with (T3) now read as “remainder subdominant to the drift”, guaranteeing the horizon of Theorem 5.3); $\beta_0 := 2N_c/\gamma_0^2$. Existence: suite, Block 7.*

Remark 5.6. As in v1: the weak-dependence threshold δ_* is a fixed positive number; $\delta_k \leq C_{\text{geom}} e^{-p_0(g_k)} \leq \delta_*$ under (T7).

Proposition 5.7 (Simultaneous compatibility, on the window). *For $\beta \geq \beta_0$ and $0 \leq k \leq \min(n_{\max}, k^*(\beta))$: (a) $g_k \leq \gamma_0$ (Theorem 5.3); (b) the Hessian bound (5); (c) the large-field bound (8) under (H-PTW); (d) $\delta_k \leq \delta_*$; (e) $\sum_k D_k \leq D_\infty < \infty$ with $D_k \leq C_D(k+1)L_{\text{RG}}^{-(d-1)k}$ and*

$$\sum_{k \geq 0} (k+1)2^{-3k} = \left(\frac{8}{7}\right)^2 = \frac{64}{49} \quad (13)$$

(v1 printed 8/49; the value, verified in the suite, is 64/49 — finiteness, which is all that is used, is unaffected); (f) $\alpha_ = (C_E/2 + C_B C_P)^{-1} > 0$, L -independent (Rothaus; the identity $(C_E/2 + C_B C_P)^{-1} = 2/(C_E + 2C_B C_P)$ is machine-checked). The non-circular order of determination is as in v1's §5.2.*

6 Assembly: proof of Theorem 1.1

As in v1's §6, with the v2 statuses: Step 1 (uniform LSI on the window) uses Theorem 5.3 (finite horizon), Proposition 2.7, the cross-scale bounds of [E26III]–[E26V], Theorem 3.3 (under (H-PTW)), and Proposition 5.7. Step 2 (DLR-LSI) uses [E26VI]–[E26VII] with the pointwise (in $\bar{U}$) large-field bound. Step 3 (Stroock–Zegarlinski) and Step 4 (reflection positivity $\Rightarrow$ gap) are as in v1. Every step is confined to the volume window; beyond it, (H-XOVER) would be needed.

7 Coupling constant bootstrap: complete proof (corrected)

7.1 Inductive conditions at a single scale

Definition 7.1. (IC1)–(IC6) as in v1.

Assumption 7.2 (One-step RG map (published input)). *As in v1.*

Remark 7.3. As in v1: strictly one-step.

7.2 Normalizations and the Wilson direction

Lemma 7.4 (Second variation of the Wilson density; proof rewritten in v2). *Let $T \in \mathfrak{su}(N_c)$ with $\|T\|^2 = -2 \text{Tr}(T^2) = 1$ (Tr fundamental) and $U(t) = \exp(tT)$. Then*

$$W(U(t)) = \frac{t^2}{4N_c} + O(t^4), \quad \left. \frac{d^2}{dt^2} W(U(t)) \right|_{t=0} = \frac{1}{2N_c}. \quad (14)$$

Proof. (v2; v1's proof mixed the normalized trace “tr” with the fundamental Tr, writing “ $\text{Tr}(T^2) = N_c \text{tr}(T^2) = -N_c/2$ ” alongside “ $\text{Re Tr} = N_c - t^2/4$ ”, which are mutually inconsistent; the statement (14) is nevertheless correct in the fundamental convention, as verified numerically for $\text{SU}(2/3/4)$ in the suite.) With $\|T\|^2 = -2 \text{Tr}(T^2) = 1$: $\text{Tr}(T^2) = -\frac{1}{2}$. Then $\text{Re Tr} \exp(tT) = N_c + \frac{t^2}{2} \text{Re Tr}(T^2) + O(t^4) = N_c - \frac{t^2}{4} + O(t^4)$, so $W = 1 - \frac{1}{N_c} \text{Re Tr} = \frac{t^2}{4N_c} + O(t^4)$. $\square$

Definition 7.5 (Wilson projection). *As in (11), with $W''(0) = 1/(2N_c)$; well-definedness and consistency as in Lemma 7.6.*

Lemma 7.6 (Well-definedness and consistency). *As in v1 (translation covariance; conjugation invariance; the extraction recovers β_k when $S_k^{\text{loc,irr}}$ has vanishing plaquette-second-variation). Verified numerically on a one-plaquette action $\beta A + c W^2$ (the W^2 term has vanishing second variation), 12 cases, exact recovery; suite.*

7.3 Base case

Lemma 7.7. *As in v1: the Wilson action satisfies (IC1)–(IC6) at $k = 0$ for $\beta \geq \beta_0$.*

7.4 Vacuum polarization structure

Lemma 7.8 (Unique marginal operator). *As in v1 (Ward–Takahashi + hypercubic symmetry): $\Pi_{\mu\nu}^{(k)}(p) = c_k(\delta_{\mu\nu}\hat{\Delta}(p) - \hat{a}_\mu\hat{a}_\nu) + \Pi_{\mu\nu}^{(k)'}(p)$.*

Remark 7.9. As in v1: $b_k(g_k) := c_k(g_k)/(2N_c)$ is the drift of g^{-2} .

Lemma 7.10 (Cauchy bound on remainder). *As in v1: $|b_k(g_k) - b_k(0)| \leq C_{\text{an}}g_k^2$.*

Lemma 7.11 (One-loop value; sign corrected in v2). $b_k(0) = -2b_0 \ln L_{\text{RG}}$: *integrating out ultraviolet fluctuations between scales k and $k+1$ lowers g^{-2} (asymptotic freedom); equivalently the shift of β is $-4N_cb_0 \ln L_{\text{RG}}$. The magnitude is the universal one-loop coefficient [26, 27, 28, 29]. (v1 stated $+2b_0 \ln L_{\text{RG}}$, citing a normalization-compatibility argument (§7.8) which compared to Balaban’s recursion with an index convention running toward the ultraviolet; when both recursions are oriented fine→coarse, the sign is $-$. Remark 5.4 gives the independent witnesses.)*

7.5 Proof of Theorem 5.2 and normalization compatibility

As in v1’s §7.7 with Lemma 7.11’s sign, and §7.8 rewritten accordingly: Balaban’s recursion, oriented fine→coarse, has drift $-2b_0 \ln L_{\text{RG}}$ in g^{-2} ; the v1 orientation (expressing the finer coupling in terms of the coarser) accounts for the $+$ sign it quoted.

Note added in v2

Prepared with the companion suite `verify_2602.0051.py` (repository `THE-ERIKSSON-PROGRAMME, verification/2602-0051/`; 10 checks, all passing, ~ 3 s, NumPy/SciPy). Changes: (1) title retagged “Unconditional” $\rightarrow$ “Conditional” Closure; Theorem 1.1 restated on the volume window $L \leq e^{C/g^2+O(1)}$, $C = 24\pi^2/(11N_c)$, conditional on (H-PTW), with (H-XOVER) identified as what full volume-uniformity would additionally require; (2) the sign of the one-step coupling drift corrected (Theorem 5.2, Lemma 7.11, Table 1’s “ β_k increasing”): v1’s sign is anti-asymptotic-freedom and inconsistent with the companions’ n_{max} ; the corrected flow reverses the monotone bootstrap of Theorem 5.3 into a finite-horizon statement and aligns this paper’s conclusion exactly with the windows of `ai.viXra:2602.0032/0033` (v2) (Remark 5.5’=Remark 5.4; suite Blocks 3a–3c); (3) the dangling “??” references of v1 ((B4), Lemma 3.1, Proposition 3.4 — the announced audit appendices absent from the document) repaired: Lemma 3.1 now has a proof, and the pointwise small-factor property is the explicit hypothesis (H-PTW) with v1’s three-point sketch retained as supporting (not proving) material; (4) Lemma 7.4’s proof rewritten in the fundamental-trace convention (statement correct; v1’s proof

mixed conventions; verified numerically); (5) $\sum (k+1)2^{-3k} = 64/49$ (v1: 8/49); (6) companion identifiers filled in: [21]=ai.viXra:2602.0055, [22]=2602.0054, [23]=2602.0053, [24]=2602.0052; the v1 self-citation [25] removed; [17]–[20] retain their classic-viXra identifiers. All v1 numbered statements keep their numbers and positions; (H-PTW)/(H-XOVER) are unnumbered tagged hypotheses.

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