

# Emergence of Spacetime and a Big-Bang–Like Singularity from Octonionic Algebra

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## Abstract

We present a conceptual framework in which the octonionic division algebra is taken as a pre-geometric fundamental structure. We demonstrate that the intrinsic non-associativity of the octonions excludes the existence of a fundamental global time and enforces a symmetry reduction in order to admit physically interpretable dynamics. This reduction occurs through an algebraically singular transition, naturally interpreted as a Big-Bang–like origin. Four-dimensional spacetime emerges as a stable associative subalgebra, while cosmological expansion arises as a necessary consequence of residual non-associative degrees of freedom. The framework is mathematically consistent but remains speculative and currently lacks experimental support.

## 1 Introduction

Modern cosmology is formulated on the assumption that spacetime geometry is a fundamental structure. Cosmological singularities, most notably the Big Bang, signal the breakdown of this description and suggest that spacetime itself may be emergent.

This work explores a framework in which geometry is not postulated but arises from a more primitive algebraic structure. The octonionic division algebra is a natural candidate, being the largest real division algebra and the only one that is fundamentally non-associative. We investigate whether the emergence of spacetime and a Big-Bang–like origin follow necessarily from algebraic consistency alone.

## 2 Octonionic Algebra

The octonions  $\mathbb{O}$  form an eight-dimensional real division algebra

$$\mathbb{O} = \text{span}_{\mathbb{R}}\{e_0, e_1, \dots, e_7\}, \quad e_0 = 1, \quad (1)$$

with multiplication

$$e_i e_j = -\delta_{ij} + C_{ijk} e_k, \quad (2)$$

where  $C_{ijk}$  is totally antisymmetric and encoded by the Fano plane.

Unlike  $\mathbb{R}$ ,  $\mathbb{C}$ , and  $\mathbb{H}$ , the octonions are non-associative. The associator is defined as

$$[a, b, c] := (ab)c - a(bc), \quad (3)$$

and is generically non-vanishing.

## 3 Absence of Fundamental Time

Associativity is a prerequisite for defining a globally consistent composition of operations. Due to the non-vanishing associator, the octonionic algebra admits no canonical ordering of products.

**Proposition.** Within the full octonionic algebra, no globally defined notion of time evolution exists.

**Consequence.** Time cannot be fundamental in an octonionic framework. Any physically meaningful temporal structure must emerge from associative subalgebras.

## 4 Automorphisms and Symmetry Reduction

The automorphism group of the octonions is the exceptional Lie group

$$\text{Aut}(\mathbb{O}) = G_2. \quad (4)$$

A fully  $G_2$ -symmetric configuration preserves all non-associative degrees of freedom. However, physically interpretable dynamics require stable associative composition laws. This enforces a symmetry reduction of the form

$$G_2 \longrightarrow SU(3) \longrightarrow SU(2) \times U(1). \quad (5)$$

This reduction is not imposed externally but arises as a consistency requirement for dynamics.

## 5 The Big Bang as an Algebraic Singularity

We introduce an order parameter measuring non-associativity,

$$\Phi := \langle [e_i, e_j, e_k] \rangle. \quad (6)$$

In a maximally symmetric phase,  $\Phi \neq 0$ . As associative subalgebras become dynamically selected,

$$\Phi \rightarrow 0 \quad (7)$$

on the emergent physical sector.

The transition is singular in the sense that

$$\frac{d\Phi}{dt} \rightarrow \infty, \quad (8)$$

even though no spacetime metric has yet been defined.

**Interpretation.** The Big Bang corresponds to an algebraic singularity: the transition from maximal non-associativity to effective associativity.

## 6 Emergence of Spacetime

Any triple of compatible imaginary units generates a quaternionic subalgebra

$$\mathbb{H} \subset \mathbb{O}, \quad (9)$$

on which the associator vanishes identically.

The octonionic norm

$$N(x) = x\bar{x} \quad (10)$$

induces a quadratic form on  $\mathbb{H}$ , which may be identified with a Lorentzian metric,

$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2. \quad (11)$$

Thus, four-dimensional spacetime emerges as a stable associative subalgebra rather than being fundamental.

## 7 Octonionic Action Principle

Since spacetime geometry is not assumed a priori, the fundamental action must be purely algebraic. The associator provides the natural dynamical measure.

We introduce an octonionic field

$$\Psi(\tau) \in \mathbb{O}, \quad (12)$$

where  $\tau$  is a pre-geometric ordering parameter.

The fundamental action is defined as

$$S_{\mathbb{O}} = \int d\tau \left[ \frac{1}{2} \langle D_{\tau} \Psi, D_{\tau} \Psi \rangle - \frac{\lambda}{4} \langle [\Psi, \Psi, \Psi], [\Psi, \Psi, \Psi] \rangle \right], \quad (13)$$

with  $\lambda > 0$ .

Variation yields the equation of motion

$$D_{\tau}^2 \Psi + \lambda \frac{\delta}{\delta \Psi} \langle [\Psi, \Psi, \Psi]^2 \rangle = 0. \quad (14)$$

The action is minimized when

$$[\Psi, \Psi, \Psi] = 0, \quad (15)$$

forcing the system into associative subalgebras.

## 8 Cosmological Expansion

Residual non-associativity acts as a positive effective energy density. An effective stress-energy contribution may be written as

$$T_{\mu\nu}^{(\text{assoc})} \sim \langle [e_{\mu}, e_{\alpha}, e_{\beta}] [e_{\nu}, e^{\alpha}, e^{\beta}] \rangle. \quad (16)$$

This leads to an effective Friedmann equation,

$$\left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \rho_{\text{assoc}}, \quad (17)$$

implying unavoidable expansion as long as non-associative degrees of freedom relax.

## 9 Discussion

Within this framework:

- time is emergent rather than fundamental,
- consistent dynamics enforce symmetry reduction,
- the Big Bang is an algebraic singularity,
- spacetime arises as an associative subalgebra,
- expansion follows from residual non-associativity.

These features are not independent assumptions but consequences of the octonionic structure.

## 10 Conclusion

Assuming octonionic algebra as a pre-geometric foundation naturally leads to an emergent spacetime and a Big-Bang-like origin. The results follow from algebraic consistency alone. While speculative, the framework provides a coherent foundations-based perspective on the origin of spacetime and cosmology.

**Disclaimer.** This work is purely theoretical. No experimental verification is currently known.

## A Mathematical Appendix

### A.1 Basic Identities of the Octonionic Associator

Let  $\mathbb{O}$  denote the octonionic division algebra. The associator is defined as

$$[a, b, c] := (ab)c - a(bc). \quad (18)$$

The associator satisfies the following fundamental identities:

**Total Antisymmetry.** For all  $a, b, c \in \mathbb{O}$ ,

$$[a, b, c] = -[b, a, c] = -[a, c, b]. \quad (19)$$

**Alternativity.** The octonions are alternative, implying

$$[a, a, b] = [b, a, a] = 0. \quad (20)$$

As a consequence, any subalgebra generated by two elements is associative.

**Norm Compatibility.** Let  $\bar{x}$  denote octonionic conjugation. The associator satisfies

$$\langle [a, b, c], 1 \rangle = 0, \quad (21)$$

and hence is purely imaginary. Moreover,

$$\langle [a, b, c], [a, b, c] \rangle \geq 0. \quad (22)$$

This ensures that quadratic expressions constructed from the associator define positive semi-definite scalar quantities.

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### A.2 Quaternionic Subalgebras

For any triple of imaginary units  $(e_i, e_j, e_k)$  satisfying

$$e_i e_j = e_k, \quad (23)$$

and closed under cyclic permutation, the generated subalgebra

$$\mathbb{H}_{ijk} = \text{span}\{1, e_i, e_j, e_k\} \quad (24)$$

is isomorphic to the quaternion algebra  $\mathbb{H}$ .

**Proposition.** On  $\mathbb{H}_{ijk}$ , the associator vanishes identically:

$$[a, b, c] = 0 \quad \forall a, b, c \in \mathbb{H}_{ijk}. \quad (25)$$

**Proof.** Quaternionic multiplication is associative. Since  $\mathbb{H}_{ijk}$  is closed under multiplication, the associator vanishes by restriction.  $\square$

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### A.3 The Exceptional Lie Group $G_2$

The automorphism group of the octonions is the exceptional Lie group

$$\text{Aut}(\mathbb{O}) = G_2. \quad (26)$$

An automorphism  $\phi \in G_2$  preserves multiplication:

$$\phi(ab) = \phi(a)\phi(b), \quad (27)$$

and therefore preserves both the norm and the associator:

$$\phi([a, b, c]) = [\phi(a), \phi(b), \phi(c)]. \quad (28)$$

**Dimension.** The Lie algebra  $\mathfrak{g}_2$  has dimension 14 and is the stabilizer of the associative 3-form

$$\varphi = \frac{1}{6} C_{ijk} e^i \wedge e^j \wedge e^k. \quad (29)$$

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### A.4 Associative 3-Form and $G_2$ Structure

The octonionic structure constants define a canonical 3-form

$$\varphi_{ijk} := C_{ijk}, \quad (30)$$

which satisfies

$$\nabla\varphi = 0 \quad \text{iff the connection has } G_2 \text{ holonomy.} \quad (31)$$

The associator can be written in index notation as

$$[a, b, c]^m = \varphi^m_{ij} \varphi^j_{kl} a^i b^k c^l, \quad (32)$$

up to normalization.

This expression makes explicit that non-associativity is controlled by the  $G_2$ -invariant tensor  $\varphi$ .

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### A.5 Symmetry Reduction and Stabilizers

Choosing a quaternionic subalgebra  $\mathbb{H} \subset \mathbb{O}$  breaks  $G_2$  to the stabilizer subgroup

$$G_2 \longrightarrow SU(2). \quad (33)$$

More generally, selecting a preferred imaginary unit induces the chain

$$G_2 \supset SU(3) \supset SU(2) \times U(1). \quad (34)$$

This reduction corresponds to suppressing non-associative directions orthogonal to the chosen subalgebra.

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## A.6 Justification of the Action Functional

The scalar

$$\langle [\Psi, \Psi, \Psi], [\Psi, \Psi, \Psi] \rangle \quad (35)$$

is:

- real,
- positive semi-definite,
- invariant under  $G_2$ .

Therefore, it is the lowest-order nontrivial invariant available for constructing a purely algebraic action.

In the associative limit, this term vanishes identically, justifying its interpretation as a non-associative energy density.

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## A.7 Summary of Mathematical Structure

- Non-associativity is encoded in the octonionic associator.
- $G_2$  preserves both the norm and the associator.
- Quaternionic subalgebras are maximally associative sectors.
- Symmetry reduction corresponds to selecting associative directions.
- The action principle is uniquely fixed at lowest order.

These results support the interpretation of the Big-Bang-like transition as an algebraic singularity driven by symmetry reduction.

# B Projection to Effective Einstein Equations

## B.1 Associative Projection Map

Let  $\mathbb{O}$  denote the octonionic algebra and  $\mathbb{H} \subset \mathbb{O}$  a selected quaternionic subalgebra. We define a projection operator

$$\mathcal{P}_{\mathbb{H}} : \mathbb{O} \longrightarrow \mathbb{H}, \quad (36)$$

such that for any  $x \in \mathbb{O}$ ,

$$x = x_{\parallel} + x_{\perp}, \quad x_{\parallel} := \mathcal{P}_{\mathbb{H}}(x), \quad x_{\perp} := x - x_{\parallel}. \quad (37)$$

By construction,

$$[a, b, c] = 0 \quad \forall a, b, c \in \mathbb{H}. \quad (38)$$

All non-associative contributions are therefore contained in  $x_{\perp}$ .

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## B.2 Decomposition of the Octonionic Field

We decompose the fundamental octonionic field as

$$\Psi = \psi + \chi, \quad (39)$$

where

$$\psi := \mathcal{P}_{\mathbb{H}}(\Psi), \quad \chi := \Psi - \psi. \quad (40)$$

Here  $\psi \in \mathbb{H}$  describes the emergent spacetime sector, while  $\chi$  parametrizes non-associative fluctuations.

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## B.3 Projected Action

Inserting the decomposition into the fundamental action,

$$S_{\mathbb{O}} = \int d\tau \left[ \frac{1}{2} \langle D_{\tau} \Psi, D_{\tau} \Psi \rangle - \frac{\lambda}{4} \langle [\Psi, \Psi, \Psi], [\Psi, \Psi, \Psi] \rangle \right], \quad (41)$$

and projecting onto  $\mathbb{H}$  yields

$$S_{\text{eff}} = \int d\tau \left[ \frac{1}{2} \langle D_{\tau} \psi, D_{\tau} \psi \rangle - V_{\text{eff}}(\chi) \right], \quad (42)$$

where the effective potential is

$$V_{\text{eff}}(\chi) = \frac{\lambda}{4} \langle [\chi, \chi, \chi], [\chi, \chi, \chi] \rangle + \mathcal{O}(\chi^4). \quad (43)$$

To lowest order, the associative sector  $\psi$  is dynamically decoupled from non-associative fluctuations.

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## B.4 Emergent Metric and Covariant Derivative

On the associative subalgebra  $\mathbb{H}$ , the octonionic norm induces a Lorentzian metric

$$g_{\mu\nu} = \langle e_{\mu}, e_{\nu} \rangle, \quad \mu, \nu = 0, 1, 2, 3. \quad (44)$$

The derivative  $D_{\tau}$  induces an effective spacetime covariant derivative

$$D_{\mu} \psi = \partial_{\mu} \psi + \Gamma_{\mu} \psi, \quad (45)$$

where  $\Gamma_{\mu}$  arises from variations of the projection  $\mathcal{P}_{\mathbb{H}}$  along  $\tau$ .

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## B.5 Effective Gravitational Action

At leading order, the projected action can be rewritten as

$$S_{\text{grav}} = \frac{1}{16\pi G_{\text{eff}}} \int d^4x \sqrt{-g} R + \int d^4x \sqrt{-g} \mathcal{L}_{\chi}, \quad (46)$$

where:

- $R$  is the Ricci scalar constructed from  $g_{\mu\nu}$ ,
- $G_{\text{eff}}$  is an emergent gravitational coupling,
- $\mathcal{L}_{\chi}$  encodes non-associative corrections.

The Einstein–Hilbert structure arises from the kinetic term of the associative projection  $\psi$ .

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## B.6 Emergent Einstein Equations

Variation with respect to the metric yields

$$G_{\mu\nu} = 8\pi G_{\text{eff}} \left( T_{\mu\nu}^{(\chi)} + T_{\mu\nu}^{(\text{matter})} \right), \quad (47)$$

where the non-associative contribution is given by

$$T_{\mu\nu}^{(\chi)} = -\frac{2}{\sqrt{-g}} \frac{\delta}{\delta g^{\mu\nu}} \int d^4x \sqrt{-g} \mathcal{L}_\chi. \quad (48)$$

To leading order,

$$T_{\mu\nu}^{(\chi)} \sim \langle [\chi, \chi, \chi]^2 \rangle g_{\mu\nu}, \quad (49)$$

which acts as an effective cosmological energy density.

## B.7 Cosmological Limit

In a homogeneous and isotropic sector, the projected Einstein equations reduce to

$$\left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi G_{\text{eff}}}{3} \rho_\chi, \quad (50)$$

where

$$\rho_\chi \propto \langle [\chi, \chi, \chi]^2 \rangle. \quad (51)$$

As  $\chi \rightarrow 0$ , the universe approaches a standard Einstein regime.

## B.8 Interpretational Remarks

- Einstein gravity emerges as a low-energy, associative limit of octonionic dynamics.
- The Big-Bang-like transition corresponds to  $\chi$ -dominated dynamics.
- Non-associativity provides a natural source for early-universe expansion.

No additional geometric postulates are required.

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