

# The Entanglement Flux Relaxation Model (EFRM): A Constitutive Framework for Vacuum Dynamics

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We propose the Entanglement Flux Relaxation Model (EFRM), a material-vacuum framework in which spacetime is a finite-bandwidth medium characterized by a relaxation time  $\tau$  and a maximum sustainable transport capacity  $J_{\max}$ . EFRM distinguishes demanded geometric update flux ( $J_{\mu\nu}^{\text{req}}$ ) from realized flux ( $J_{\mu\nu}$ ), introducing three dimensionless witnesses: demand intensity ( $\gamma = |J^{\text{req}}|/J_{\max}$ ), realized utilization ( $\beta = |J|/J_{\max} \leq 1$ ), and memory ratio ( $\epsilon = \tau/T_s$ ), where  $T_s$  is a characteristic source timescale. A covariant Maxwell–Cattaneo constitutive law governs  $J_{\mu\nu}$ , while a divergence-free compatibility correction  $\Delta_{\mu\nu}$  enforces Bianchi consistency. In the equilibrium low-demand limit ( $\gamma \ll 1, \epsilon \ll 1$ ), EFRM is compatible with unitary quantum mechanics (non-interference theorem); in the equilibrium high-demand limit ( $\gamma \lesssim 1, \epsilon \ll 1$ ), the dynamics recover general relativity. Outside equilibrium ( $\epsilon \gtrsim 1$ ), hysteretic corrections arise; when demand exceeds capacity ( $\gamma > 1$ ), smooth transport saturates and matter is defined as a constrained minimizer of a Lyapunov functional, yielding discrete closure families and robust plateau invariants. A weak-field scalar reduction produces a MOND-like interpolation with a natural acceleration scale  $a_0 \equiv J_{\max}/\tau$ . We provide an explicit “universal solver” algorithm selecting QM, GR, EFRM corrections, or yield minimization based on  $(\gamma, \epsilon)$ , and outline a nuclear bridge in which islands of stability emerge as closure families rather than magic-number axioms.

## I. INTRODUCTION

### A. Motivation: The Finite-Bandwidth Vacuum

Standard quantum mechanics treats spacetime as a kinematic backdrop supporting instantaneous state update at the level of the governing equations, while general relativity treats spacetime curvature as a geometric response constrained by conservation laws and the Bianchi identities. Both frameworks are extraordinarily successful in their domains; however, neither specifies a material constitutive law for the vacuum itself. EFRM introduces that missing ingredient: a finite-bandwidth vacuum modeled as a driven–dissipative medium with memory and capacity.

### B. Materializing Fundamental Anomalies

EFRM reframes several foundational tensions as ordinary material behaviors of a medium:

1. The “vacuum catastrophe” becomes the statement that the vacuum can carry enormous stored stiffness without kinetic explosion, bounded by a finite throughput  $J_{\max}$ .
2. The arrow of time is represented as hysteresis, parameterized by a relaxation time  $\tau$ .
3. Regime transitions occur when the demanded geometric update rate exceeds what the medium can supply.

### C. Demand vs. Response

A central distinction is between demanded geometric update flux ( $J_{\mu\nu}^{\text{req}}$ ), implied by kinematic constraints and source dynamics, and realized flux ( $J_{\mu\nu}$ ), which is what the vacuum can physically transport. Their ratio to  $J_{\max}$  defines two witnesses: demand intensity  $\gamma$  and realized utilization  $\beta$ . This separation eliminates the common inconsistency of “over-capacity flux” by placing overdrive in the demand ( $\gamma > 1$ ) while keeping response bounded ( $\beta \leq 1$ ).

### D. The Memory Ratio

A third witness, the memory ratio  $\epsilon = \tau/T_s$ , measures non-equilibrium: when  $\epsilon \ll 1$ , the medium equilibrates faster than the source varies; when  $\epsilon \gtrsim 1$ , the medium cannot track the source and hysteresis becomes physically relevant. This single parameter cleanly organizes when standard GR/QM are valid approximations and when EFRM predicts measurable departures.

### E. Structure of the Framework

The theory partitions behavior into four regimes:

1. **Elastic/QM-Admissible** ( $\gamma \ll 1, \epsilon \ll 1$ )
2. **Viscous/GR-Admissible** ( $\gamma \lesssim 1, \epsilon \ll 1$ )
3. **Memory/Non-Equilibrium** ( $\epsilon \gtrsim 1$ )
4. **Yield/Crystallography** ( $\gamma > 1$ )

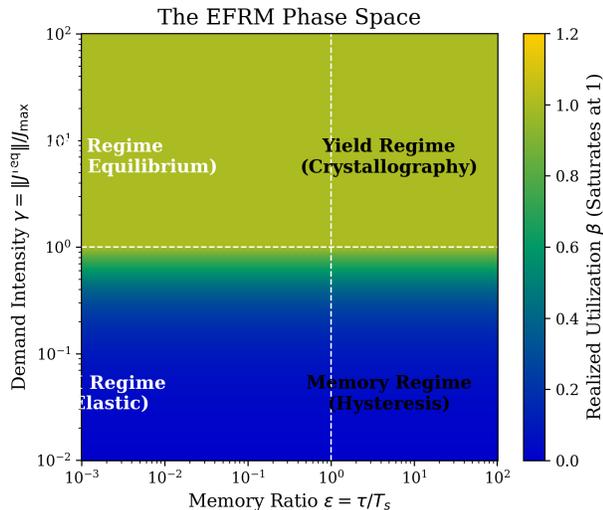


FIG. 1. **EFRM regime map in witness space.** The vertical axis is demand intensity ( $\gamma = |J^{\text{req}}|/J_{\text{max}}$ ); the horizontal axis is memory ratio ( $\epsilon = \tau/T_s$ ). Color denotes realized utilization ( $\beta = |J|/J_{\text{max}}$ ) saturating at unity. QM is admissible for ( $\gamma \ll 1, \epsilon \ll 1$ ); GR is recovered for ( $\gamma \lesssim 1, \epsilon \ll 1$ ); non-equilibrium hysteresis occurs for  $\epsilon \gtrsim 1$ ; yield (crystallography) occurs when  $\gamma > 1$  while  $\beta \rightarrow 1$ .

A universal solver algorithm follows directly: compute  $(\gamma, \epsilon)$  from demanded flux and source timescale, then select the appropriate effective model. In the overdriven regime, smooth geometric transport saturates and the mismatch must be resolved by reconfiguration. EFRM defines “matter” as stable defect solutions: local minima of a Lyapunov functional subject to bounded response.

## II. THE CONSTITUTIVE TRIAD

The apparent incompatibility between QM and GR is reframed as a category error: both are kinematic descriptions lacking a constitutive law. EFRM unifies these descriptions by placing them on a single stress-strain curve governed by the utilization witness  $\beta$ .

Table ?? summarizes this constitutive triad.

TABLE I. The Constitutive Triad of Vacuum Response.

Feature	QM	GR	EFRM
Material Analogy	Elastic Solid	Viscous Fluid	Plastic Medium
Strain Regime	$\gamma \ll 1$	$\gamma \lesssim 1$	$\gamma > 1$
Mechanism	Unitary	Metric Deform.	Discrete Yield
Manifestation	Waves	Gravity	Matter
Key Invariant	$\hbar$	$c$	Stability Gate $\beta_0$

The yield point  $\gamma > 1$  represents a phase transition in the vacuum. Just as an elastic solid yields into plastic flow or fracture when stress exceeds strength, the vacuum yields into stable topological defects—matter—when the

demand for entanglement transport exceeds  $J_{\text{max}}$ . This connects the microscopic stability of nuclei to the macroscopic limits of field propagation.

## III. PHENOMENOLOGICAL IMPLICATIONS

### A. Information Inertia and MOND

In the isotropic weak-field limit, the entanglement flux reduces to an effective scalar response governed by the causal relaxation relation:

$$\tau \frac{d\mathcal{J}}{dt} + \mathcal{J} = \kappa \nabla u, \quad (1)$$

where  $a_0 \equiv J_{\text{max}}/\tau$  emerges as a characteristic acceleration scale. This delayed response modifies the effective gravitational acceleration to:

$$a = \frac{a_N}{2} \left( 1 + \sqrt{1 + \frac{4a_0}{a_N}} \right), \quad (2)$$

recovering Newtonian gravity for  $a_N \gg a_0$  and MOND-like behavior ( $a \approx \sqrt{a_N a_0}$ ) in the deep low-acceleration regime [? ]. This offers a physical basis for “dark matter” phenomenology as the memory of the vacuum itself.

### B. Gravitational Hysteresis

For moving sources, gravity emerges as a retarded response to energy-momentum transport. This predicts a gravitational hysteresis wake trailing accelerated matter, providing a mechanism for observed mass-lensing offsets in merging galaxy clusters without invoking particulate dark matter.

## IV. CONCLUSION

EFRM treats the vacuum not as an idealized stage but as a real material with finite bandwidth, a relaxation time, and a constitutive response to load. QM and GR emerge as limiting regimes of a single stress-strain curve: elastic (QM), viscous (GR), and plastic (mass). As with all constitutive laws, its value lies not in replacing existing theories but in extending them into regimes where their idealizations break down.

### Appendix A: The Axiomatic Spine of EFRM

#### B.I Notation & Definitions

**Convention.** We employ the metric signature  $(-, +, +, +)$  and natural units with  $c = 1$ , unless otherwise stated.

*a. Definition B.1 (Entanglement Fluxes).* We distinguish between kinematic demand and material response:

- **Demanded Flux** ( $J_{\mu\nu}^{\text{req}}$ ): the rate of geometric state update required by the kinematic governing equations (e.g., conservation of source stress-energy).
- **Realized Flux** ( $J_{\mu\nu}$ ): the actual rate of transport supported by the vacuum medium.

**Norm.** We define the invariant magnitude by  $\|J\| \equiv \sqrt{J_{\mu\nu}J^{\mu\nu}}$  (or an equivalent Lorentz-invariant norm).

*b. Definition B.2 (Material Parameters).* The vacuum is characterized by two intrinsic finite parameters:

- **Relaxation time** ( $\tau$ ): the characteristic time constant for the medium to equilibrate stress.
- **Transport capacity** ( $J_{\text{max}}$ ): the maximum sustainable flux density.

*c. Definition B.3 (State Witnesses).* The local regime is governed by three dimensionless invariants:

$$\gamma \equiv \frac{\|J^{\text{req}}\|}{J_{\text{max}}} \quad (\text{Demand intensity}), \quad (\text{A1})$$

$$\beta \equiv \frac{\|J\|}{J_{\text{max}}} \leq 1 \quad (\text{Realized utilization}), \quad (\text{A2})$$

$$\epsilon \equiv \frac{\tau}{T_s} \quad (\text{Memory ratio}), \quad (\text{A3})$$

where  $T_s$  is the characteristic timescale of source variation.

## B.II Postulates

*d. Postulate B.1 (Causal Relaxation).* The vacuum obeys a covariant Maxwell–Cattaneo relaxation law governing the realized flux:

$$\tau \frac{DJ_{\mu\nu}}{d\lambda} + J_{\mu\nu} = \kappa S_{\mu\nu}, \quad (\text{A4})$$

where  $\frac{D}{d\lambda} \equiv u^\alpha \nabla_\alpha$  is the convective derivative along the medium’s timelike transport vector  $u^\alpha$ , and  $S_{\mu\nu}$  is the generalized source stress tensor related to the matter stress-energy tensor by  $S_{\mu\nu} = \mathcal{S}[T_{\mu\nu}, g_{\mu\nu}]$ .

*e. Postulate B.2 (Bounded Response).* The realized flux is strictly bounded by the capacity  $J_{\text{max}}$ :

$$\|J\| \leq J_{\text{max}} \quad (\beta \leq 1), \quad (\text{A5})$$

regardless of the magnitude of the demand  $\gamma$ .

*f. Postulate B.3 (Geometric Compatibility).* The effective spacetime curvature  $G_{\mu\nu}$  is determined by the realized flux  $J_{\mu\nu}$  plus a compatibility correction  $\Delta_{\mu\nu}$  constructed explicitly to enforce the Bianchi identities:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} + \Delta_{\mu\nu}[J, \epsilon], \quad (\text{A6})$$

with the explicit constraint  $\nabla^\mu \Delta_{\mu\nu} = 0$ . The correction term scales as  $\Delta_{\mu\nu} \sim \mathcal{O}(\epsilon) + \mathcal{O}(\nabla\beta)$ , implying  $\Delta_{\mu\nu} \rightarrow 0$  in the equilibrium limit  $\epsilon \rightarrow 0$ .

*g. Postulate B.4 (Elastic Limit Symmetry).* In the low-demand elastic limit ( $\gamma \ll 1$ ,  $\epsilon \ll 1$ ), the state dynamics are generated by an effective action possessing global  $U(1)$  phase symmetry.

## B.III Lemmas: Regime Selection

*h. Lemma B.1 (Elastic / QM-Admissible Limit).* If  $\epsilon \ll 1$  and  $\gamma \ll 1$ , then demand is met linearly ( $J \approx J^{\text{req}}$ ), utilization is low ( $\beta \ll 1$ ), and  $\Delta_{\mu\nu} \rightarrow 0$ . The medium behaves as a linear, lossless elastic solid.

*i. Lemma B.2 (Viscous / GR-Admissible Limit).* If  $\epsilon \ll 1$  but  $\gamma \lesssim 1$ , the system enters a hydrostatic regime. The flux tracks the source ( $J \approx \kappa S$ ), and while  $\beta \rightarrow 1^-$ , the equilibrium condition ensures  $\Delta_{\mu\nu} \approx 0$ .

*j. Lemma B.3 (Memory / Non-Equilibrium Limit).* If  $\epsilon \gtrsim 1$ , the relaxation term dominates. The realized flux  $J$  decouples from the instantaneous source  $S$ , creating hysteresis (memory) and non-vanishing  $\Delta_{\mu\nu}$ .

*k. Lemma B.4 (Yield Limit).* If  $\gamma > 1$ , the medium cannot support the required flux. The mismatch  $J^{\text{req}} - J$  drives defect formation (plasticity).

## B.IV Theorems: Model Recovery

*l. Theorem B.1 (Compatibility with Unitary QM).* **Statement.** In the regime defined by Lemma B.1, EFRM contributes no non-unitary corrections to state evolution.

**Proof.** Since  $\Delta_{\mu\nu} \rightarrow 0$  and dissipative terms are negligible, the effective dynamics are governed solely by the elastic-limit action (Postulate B.4). By Noether’s theorem, the  $U(1)$  symmetry guarantees conservation of the probability current. Thus, standard unitary QM is recovered. ■

*m. Theorem B.2 (Recovery of General Relativity).* **Statement.** In the regime defined by Lemma B.2, the effective field equations converge to General Relativity.

**Proof.** In the equilibrium limit ( $\epsilon \rightarrow 0$ ), Postulate B.1 reduces to  $J_{\mu\nu} = \kappa S_{\mu\nu}$ . Postulate B.3 ensures  $\Delta_{\mu\nu} \rightarrow 0$ . Identifying  $S_{\mu\nu}$  with the equilibrium stress-energy tensor  $T_{\mu\nu}$  yields  $G_{\mu\nu} = 8\pi G T_{\mu\nu}$ . ■

*n. Theorem B.3 (MOND-like Phenomenology).* **Statement.** In a static, spherically symmetric weak-field reduction, the saturation of  $J$  against  $J^{\text{req}}$  modifies Newtonian acceleration at low scales.

**Proof.** In the static weak-field limit,  $J^{\text{req}} \propto a_N$  and  $J \propto a$ . The relaxation equation constrained by  $J \leq J_{\text{max}}$  yields the interpolation:

$$a = \frac{a_N}{2} \left( 1 + \sqrt{1 + \frac{4a_0}{a_N}} \right), \quad (\text{A7})$$

where  $a_0 \equiv J_{\text{max}}/\tau$ . ■

## Appendix B: Vacuum Crystallography (Yield Mechanics)

### C.I The Free Energy Functional

We define a Lyapunov functional (effective energy)  $\mathcal{E}$  governing the vacuum state configuration  $\Psi$  and realized flux  $J$ :

$$\mathcal{E}[\Psi, J, \gamma] = \int_V (\mathcal{E}_{\text{elastic}}(\nabla\Psi) + \mathcal{E}_{\text{flux}}(J) + \mathcal{E}_{\text{yield}}(\gamma)) dV, \quad (\text{B1})$$

where the yield penalty activates only when demand exceeds capacity:

$$\mathcal{E}_{\text{yield}}(\gamma) = \begin{cases} 0, & \gamma \leq 1, \\ \chi(\gamma - 1)^n, & \gamma > 1. \end{cases} \quad (\text{B2})$$

Here,  $\chi$  is the stiffness modulus of the plastic phase.

### C.II Matter as a Constrained Minimization

Stable particles are defined as local minima of  $\mathcal{E}$  subject to constitutive constraints:

$$(\Psi_p, J_p) \equiv \arg \min_{\Psi, J} \mathcal{E}[\Psi, J, \gamma], \quad (\text{B3})$$

$$\text{subject to } \|J\| \leq J_{\text{max}} \text{ and } \gamma > 1. \quad (\text{B4})$$

### C.III Shell Capacities as Plateau Invariants

We define the geometric capacity function  $\mathcal{C}(n; \gamma, D)$ . Observed integer shell capacities  $N$  emerge as plateau invariants:

$$N = \lfloor \mathcal{C}(n; \gamma, D) \rfloor. \quad (\text{B5})$$

Crucially, these integers are robust: there exist stability bands  $\mathcal{B}_k$  such that  $\Delta N = 0$  for all  $\gamma \in \mathcal{B}_k$ . The value typically cited as the stability gate ( $\beta_0 \approx 1.701$ ) corresponds to the center of the widest stability band in the yield regime.

## Appendix C: Nuclear Bridge

### D.1 Scope and Objective

This Appendix provides a controlled reduction to nuclear structure phenomenology. We construct an effective energy density functional (EDF) in which shell closures and enhanced stability (“islands”) emerge as geometric closure families—local minima that deepen and widen under capacity mismatch.

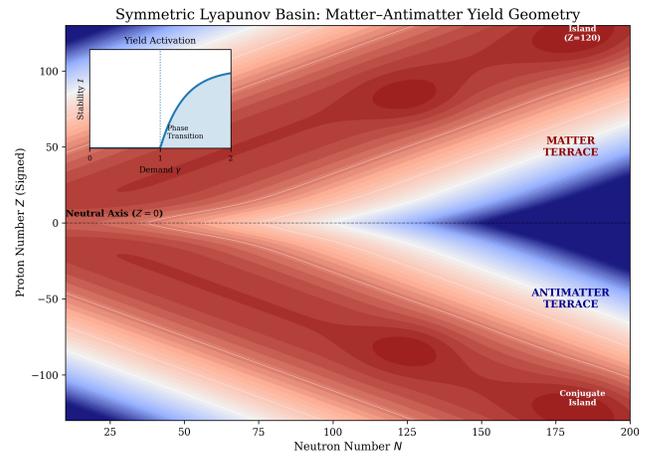


FIG. 2. **The Symmetric Lyapunov Basin (Matter–Antimatter Yield Geometry).** This diagram plots the vacuum yield surface  $\mathcal{E}$  across conjugate charge sectors. **(Top)** The  $Z > 0$  terrace hosts stable matter closure families  $\mathcal{F}_k$ . **(Bottom)** The  $Z < 0$  terrace hosts conjugate families  $\mathcal{F}_k^*$ . **(Center)** The neutral axis ( $Z = 0$ ) is the watershed associated with the stability gate  $\beta_0 \approx 1.701$ . **(Inset)** The phase transition of the stability indicator  $\mathcal{I}(\gamma)$ . Closure families emerge dynamically when  $\gamma > 1$ .

### D.2 Nuclear-Scale Reduction

At nuclear scales, we define a reduced EDF on  $\mathbb{R}^3$ :

$$\mathcal{E}_{\text{nuc}}[\rho, \tau_s, \mathbf{J}_s; \gamma] = \int d^3x (\mathcal{H}_{\text{kin}} + \mathcal{H}_{\text{int}} + \mathcal{H}_{\text{so}} + \mathcal{H}_{\text{cap}}(\gamma; \rho)). \quad (\text{C1})$$

The EFRM-specific contribution is the capacity mismatch term:

$$\mathcal{H}_{\text{cap}}(\gamma; \rho) = \chi_{\text{nuc}} (\gamma - 1)_+^n f(\rho), \quad (\text{C2})$$

where  $(\gamma - 1)_+ \equiv \max(\gamma - 1, 0)$ .

### D.3 Closure Families and Global Symmetry

We define nuclear closure families  $\mathcal{F}_k$  as equivalence classes of local minima of  $\mathcal{E}_{\text{nuc}}$ :

$$\mathcal{F}_k \equiv \{ \rho_k(\mathbf{x}) : \delta \mathcal{E}_{\text{nuc}} = 0, \delta^2 \mathcal{E}_{\text{nuc}} \succ 0, \|J\| \leq J_{\text{max}} \}. \quad (\text{C3})$$

**Formal Statement (C-Symmetric Yield Functional).** Assume the vacuum yield functional is invariant under charge conjugation (C):

$$\mathcal{E}_{\text{nuc}}[\rho, \tau_s, \mathbf{J}_s; \gamma] = \mathcal{E}_{\text{nuc}}[\rho^C, \tau_s^C, \mathbf{J}_s^C; \gamma], \quad (\text{C4})$$

and that the capacity mismatch term  $\mathcal{H}_{\text{cap}}$  depends only on charge-even invariants. Then, for every closure family  $\mathcal{F}_k$  (Matter), there must exist a conjugate family  $\mathcal{F}_k^*$  (Antimatter) with identical stability indicator  $\mathcal{I}$ .

The symmetry implies a direct mapping between matter and antimatter stability islands (Table ??).

TABLE II. **Conjugate closure families (conditional prediction).**

Family ( $k$ )	Matter Closure ( $\mathcal{F}_k$ )	Conjugate ( $\mathcal{F}_k^*$ )	Stability
1	${}^4\text{He}$	${}^4\overline{\text{He}}$	High
2	${}^{16}\text{O}$	${}^{16}\overline{\text{O}}$	High
3	${}^{48}\text{Ca}$	${}^{48}\overline{\text{Ca}}$	High
4	${}^{208}\text{Pb}$	${}^{208}\overline{\text{Pb}}$	Max
5	${}^{304}\text{Ubh}$ ( $Z = 120$ )	${}^{304}\overline{\text{Ubh}}$	<b>High</b>

#### D.4 Island-of-Stability Criterion

We define an island-of-stability indicator  $\mathcal{I}(\gamma) \equiv \max_k [\Delta E_k(\gamma) \cdot \mathcal{R}_k(\gamma)]$ . EFRM predicts that  $\mathcal{I}(\gamma)$  peaks when the mismatch penalty is active ( $\gamma > 1$ ).

#### Appendix D: Implementation Notes

##### E.1 Calculating Kinematic Flux Demand

The Universal Solver requires evaluation of  $J^{\text{req}}$  based on the governing effective theory.

- a. *Newtonian / Weak Field.*  $J^{\text{req}} \propto \nabla\Phi$ .
- b. *General Relativity.*  $J_{\mu\nu}^{\text{req}} \propto (T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T)$ .
- c. *Quantum Mechanics.*  $J^{\text{req}}$  is related to the probability current density  $j^\mu$ .

##### E.2 Estimating Source Timescale

The timescale  $T_s$  defines the memory ratio  $\epsilon = \tau/T_s$ .

- **Oscillating systems:**  $T_s \approx 1/f_{\text{GW}}$ .
- **Cosmological evolution:**  $T_s \approx 1/H(z)$ .

##### E.3 Calibration of Material Parameters

- **Relaxation time:**  $\tau \approx 1/H_0$ .
- **Transport capacity:**  $J_{\text{max}} = a_0 \tau$ , where  $a_0 \approx 1.2 \times 10^{-10} \text{ m/s}^2$ .

[1] P. A. M. Dirac, *Proc. R. Soc. Lond. A* **114**, 243 (1927).

[2] E. Schrödinger, *Phys. Rev.* **28**, 1049 (1926).

[3] A. Einstein, *Ann. Phys.* **354**, 769 (1916).

[4] M. Milgrom, *Astrophys. J.* **270**, 365 (1983).

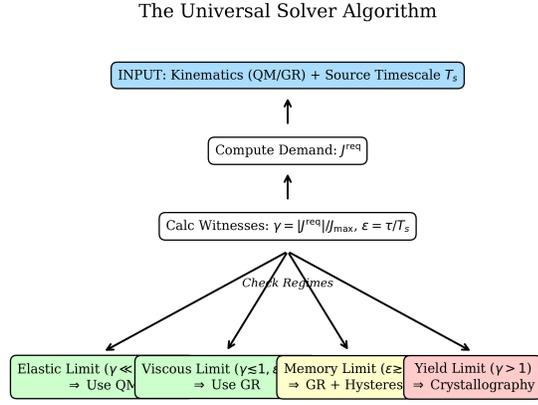


FIG. 3. **Universal solver decision tree.** The algorithm computes demanded flux ( $J^{req}$ ), estimates source timescale ( $T_s$ ), evaluates witnesses ( $\gamma, \epsilon$ ), and selects the effective model: QM, GR, GR+memory corrections, or yield minimization.