

RIP-U and the $\omega = 0$ obstruction in Davies dynamics

Upper envelopes, witness floors, and falsification protocols for separation-dependent decoherence rates

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Abstract

We isolate the dynamic hinge in typed separation \rightarrow rate \rightarrow power pipelines within the Davies (weak-coupling, Markovian) setting in finite dimension. First, we formulate an upper-envelope statement (RIP-U): under an explicit factorized bath-correlation envelope with separation-dependent amplitude $f(\epsilon)$ and integrable time profile, Fourier-transformed Davies rates inherit an $O(f(\epsilon))$ envelope. Under additional regularity assumptions preventing trivial degeneracies, this yields a worst-case bound $\kappa^\uparrow(\epsilon) \leq Cf(\epsilon)$ for the instantaneous relative loss rate of a coherence-like functional. Second, we isolate a structural obstruction to lower-envelope statements: we prove an exact identity for the $\omega = 0$ contribution to the Davies Dirichlet form,

$$\mathcal{E}_\sigma^{(0)}(O) = \frac{\gamma^{(0)}}{2} \|[S(0), O]\|_{2,\sigma}^2,$$

and derive a witness mechanism showing how $\omega = 0$ channels can enforce a non-vanishing dissipation contribution for suitable observable families. We emphasize directionality: RIP-U (upper) does not imply a positive lower envelope $\kappa^\downarrow(\epsilon)$ without additional family-qualified input. All assumptions are explicit and accompanied by falsification routes.

1 Scope and conventions

What this paper does.

- Proves a Fourier-envelope inheritance lemma for rate coefficients under a factorized correlator bound.
- States a clean upper-envelope implication (RIP-U) under explicit additional bridge assumptions.
- Proves an exact $\omega = 0$ Dirichlet identity in standard Davies form and a witness corollary.

What this paper does not do. We do not claim a Hamiltonian spectral gap and we do not claim RIP-U implies a family-wide true floor $\kappa^\downarrow(\epsilon) \geq \kappa_0$.

Parameter convention. ϵ is geometric separation/buffer width; δ is tolerance/regularization. They are never interchanged.

2 Davies setting and typed envelopes

We work in finite dimension: system Hilbert space \mathcal{H} with $\dim \mathcal{H} = d < \infty$.

Assumption 2.1 (DAVIES: weak-coupling Markovian dynamics). **[ASSUMED]**.

The system dynamics is governed by a CPTP semigroup $\mathcal{T}_t = e^{t\mathcal{L}}$ with a Davies-type GKLS generator \mathcal{L} and a faithful stationary state σ (typically $\sigma \propto e^{-\beta H}$). We assume the standard Bohr-frequency decomposition and GKLS form used in Davies theory.

2.1 A coherence-like functional and instantaneous loss

Definition 2.2 (D1: Δ -track coherence). Fix a dephasing (pinching) CPTP map Δ on $\mathcal{B}(\mathcal{H})$. Define

$$C_\Delta(\rho) := S(\rho \parallel \Delta[\rho]),$$

where $S(\rho \parallel \tau) := \text{Tr}(\rho(\log \rho - \log \tau))$ when well-defined.

Definition 2.3 (D2: instantaneous loss). Define the instantaneous change of C_Δ under \mathcal{T}_t by

$$\dot{C}_{\Delta, \text{loss}}(\rho) := - \left. \frac{d}{dt} \right|_{t=0} C_\Delta(\mathcal{T}_t(\rho)),$$

whenever the derivative exists.

Assumption 2.4 (Δ -MONO: sign convention for “loss”). **[ASSUMED]** (family restriction). For the chosen Δ and semigroup \mathcal{T}_t , we restrict attention to families \mathcal{F}_ϵ such that

$$\dot{C}_{\Delta, \text{loss}}(\rho) \geq 0 \quad \text{for all } \rho \in \mathcal{F}_\epsilon.$$

Definition 2.5 (D3: ratios and envelopes). Whenever $C_\Delta(\rho) > 0$, define

$$r(\rho) := \frac{\dot{C}_{\Delta, \text{loss}}(\rho)}{C_\Delta(\rho)}.$$

For a specified family \mathcal{F}_ϵ define

$$\kappa^\uparrow(\epsilon) := \sup_{\rho \in \mathcal{F}_\epsilon: C_\Delta(\rho) > 0, \dot{C}_{\Delta, \text{loss}}(\rho) \text{ exists}} r(\rho), \quad \kappa^\downarrow(\epsilon) := \inf_{\rho \in \mathcal{F}_\epsilon: C_\Delta(\rho) > 0, \dot{C}_{\Delta, \text{loss}}(\rho) \text{ exists}} r(\rho).$$

Remark 2.6 (Directionality). RIP-U targets $\kappa^\uparrow(\epsilon)$ (worst-case upper envelope). Minimum-maintenance-power lower bounds require $\kappa^\downarrow(\epsilon)$; RIP-U alone does not provide them.

3 Correlator envelope \Rightarrow Fourier rate envelope

Assumption 3.1 (CORR: factorized correlator envelope). **[ASSUMED]**.

For each ϵ , the relevant bath correlators satisfy

$$|G_{\alpha\beta}(t; \epsilon)| \leq g_{\alpha\beta}(t) f(\epsilon),$$

where $g_{\alpha\beta} \in L^1(\mathbb{R})$ and $f(\epsilon) \geq 0$ captures separation dependence.

Definition 3.2 (D4: Fourier rate coefficients). Define the (formal) Fourier rate coefficients by

$$\gamma_{\alpha\beta}(\omega; \epsilon) := \int_{\mathbb{R}} dt e^{i\omega t} G_{\alpha\beta}(t; \epsilon),$$

whenever the integral exists (or under the regularity hypotheses used in Davies theory).

Lemma 3.3 (Fourier envelope inheritance). **[PROVED]** (under [Definition 3.1](#) and integrability). Under [Definition 3.1](#), for all ω and ϵ for which [Definition 3.2](#) is well-defined,

$$|\gamma_{\alpha\beta}(\omega; \epsilon)| \leq \|g_{\alpha\beta}\|_{L^1} f(\epsilon).$$

Proof. By the triangle inequality and [Definition 3.1](#),

$$|\gamma_{\alpha\beta}(\omega; \epsilon)| \leq \int_{\mathbb{R}} dt |G_{\alpha\beta}(t; \epsilon)| \leq \int_{\mathbb{R}} dt |g_{\alpha\beta}(t)| f(\epsilon) = \|g_{\alpha\beta}\|_{L^1} f(\epsilon).$$

□

4 RIP-U: an upper envelope on instantaneous loss

Assumption 4.1 (REG: non-degeneracy for relative-entropy ratios). **[ASSUMED]** (family restriction).

The family \mathcal{F}_ϵ is restricted so that:

- (i) $C_\Delta(\rho) \geq c_{\min}(\delta) > 0$ uniformly over $\rho \in \mathcal{F}_\epsilon$;
- (ii) $\Delta[\rho] \succeq p_{\min}(\delta)\mathbf{1}$ uniformly over $\rho \in \mathcal{F}_\epsilon$;
- (iii) when needed, ρ stays a fixed distance (within tolerance δ) away from $\text{Fix}(\mathcal{T}_t)$.

Assumption 4.2 (BRIDGE: rate envelope \Rightarrow loss envelope). **[ASSUMED]**.

Assume [Definitions 2.1](#), [3.1](#) and [4.1](#). Suppose the Davies generator is built from Fourier rate coefficients $\gamma_{\alpha\beta}(\omega; \epsilon)$ in such a way that the envelope from [Definition 3.3](#) implies: there exists a constant $C_{\text{RIP}} > 0$ such that for all $\rho \in \mathcal{F}_\epsilon$,

$$\dot{C}_{\Delta, \text{loss}}(\rho) \leq C_{\text{RIP}} f(\epsilon) C_\Delta(\rho).$$

Theorem 4.3 (T1: RIP-U (upper envelope)). **[PROVED]** (given [Definitions 2.1](#), [2.4](#), [3.1](#), [4.1](#) and [4.2](#)).

Under [Definitions 2.1](#), [2.4](#), [3.1](#), [4.1](#) and [4.2](#), for all $\rho \in \mathcal{F}_\epsilon$,

$$r(\rho) \leq C_{\text{RIP}} f(\epsilon),$$

and hence

$$\kappa^\uparrow(\epsilon) \leq C_{\text{RIP}} f(\epsilon).$$

Proof. By [Definition 4.2](#), $\dot{C}_{\Delta, \text{loss}}(\rho) \leq C_{\text{RIP}} f(\epsilon) C_\Delta(\rho)$ for all $\rho \in \mathcal{F}_\epsilon$ with $C_\Delta(\rho) > 0$, hence $r(\rho) \leq C_{\text{RIP}} f(\epsilon)$. Taking the supremum yields the stated envelope. □

Remark 4.4 (RIP-U does not imply a floor). RIP-U controls worst-case *upper* loss rates. It is compatible with $\kappa^\downarrow(\epsilon)$ being arbitrarily small on a large family. Thus RIP-U alone does not imply minimum maintenance power lower bounds in pipelines requiring a positive lower envelope.

5 The $\omega = 0$ obstruction: an exact Dirichlet identity

Definition 5.1 (D5: σ -weighted 2-norm). For a faithful state σ , define

$$\|X\|_{2,\sigma} := \sqrt{\text{Tr}(\sigma^{1/2} X^\dagger \sigma^{1/2} X)}.$$

Definition 5.2 (D6: Dirichlet form (one convenient convention)). Let \mathcal{L} be a GKLS generator with faithful stationary state σ . Define the Dirichlet form on observables by

$$\mathcal{E}_\sigma(O) := -\text{Tr}\left(\sigma^{1/2} O^\dagger \sigma^{1/2} \mathcal{L}^\sharp(O)\right),$$

where \mathcal{L}^\sharp denotes the Heisenberg-picture action of the generator on observables.

Assumption 5.3 (BOHR: frequency decomposition). **[ASSUMED]** (standard in Davies theory). Let H be the system Hamiltonian with spectral projectors $\{\Pi_E\}$. For a system coupling operator S , define its Bohr components

$$S(\omega) := \sum_{E'-E=\omega} \Pi_E S \Pi_{E'}.$$

Assume the Davies generator uses dissipators built from $S(\omega)$ with scalar rates $\gamma(\omega) \geq 0$.

Lemma 5.4 (Commutation at zero frequency). **[PROVED]**.

If $\sigma \propto e^{-\beta H}$ and $S(0)$ is the $\omega = 0$ Bohr component, then

$$[S(0), H] = 0 \quad \Rightarrow \quad [S(0), \sigma] = 0 \quad \text{and} \quad [S(0), \sigma^{1/2}] = 0.$$

Proof. By construction, $S(0) = \sum_E \Pi_E S \Pi_E$ commutes with $H = \sum_E E \Pi_E$, hence with any function of H , in particular σ and $\sigma^{1/2}$. \square

Theorem 5.5 (T2: $\omega = 0$ Dirichlet identity). **[PROVED]** (under *Definitions 2.1 and 5.3* in standard Davies form).

The zero-frequency contribution $\mathcal{E}_\sigma^{(0)}$ to the Dirichlet form satisfies

$$\mathcal{E}_\sigma^{(0)}(O) = \frac{\gamma(0)}{2} \|[S(0), O]\|_{2,\sigma}^2.$$

Proof. We give a direct expansion in Appendix A. \square

Remark 5.6 (The $\gamma(0) = 0$ case). If $\gamma(0) = 0$ (e.g. baths with vanishing zero-frequency spectral density), then $\mathcal{E}_\sigma^{(0)} \equiv 0$ and this particular $\omega = 0$ obstruction mechanism disappears. Any lower-envelope floor must then come from other frequency sectors or from structural restrictions defining the family.

Corollary 5.7 (C1: $\omega = 0$ witness mechanism). **[PROVED]**.

Assume $\gamma(0) > 0$. If there exists an observable family O_ϵ such that

$$\frac{\|[S(0), O_\epsilon]\|_{2,\sigma}^2}{\|O_\epsilon\|_{2,\sigma}^2} \geq c > 0,$$

then

$$\mathcal{E}_\sigma^{(0)}(O_\epsilon) \geq \frac{\gamma(0)}{2} c \|O_\epsilon\|_{2,\sigma}^2.$$

Proof. Immediate from [Definition 5.5](#). □

Remark 5.8 (C1 is not RIP-L). C1 is a mechanism for a chosen observable family. Converting it into a uniform lower-envelope bound $\kappa^\downarrow(\epsilon) \geq \kappa_0$ over a state family requires additional typed bridges.

Assumption 5.9 (RIP-L: true floor (family-qualified)). **[CONJECTURED]**.

There exist $\epsilon_* > 0$ and $\kappa_0 > 0$ such that

$$\kappa^\downarrow(\epsilon) \geq \kappa_0 \quad \text{for all } \epsilon \geq \epsilon_*,$$

for a specified family \mathcal{F}_ϵ (with any regularity restrictions stated explicitly).

6 Falsification matrix

| Item | Status | Verification proxy | Concrete falsifier |
|------------------------|----------------------|--|--|
| CORR | [ASSUMED] | compute $G_{\alpha\beta}(t; \epsilon)$ | separation dependence not factorizable as $g(t)f(\epsilon)$ |
| Fourier envelope lemma | [PROVED] | compute $\ g\ _{L^1}$ | violate integrability/envelope inequality |
| Δ -MONO | [ASSUMED] | evaluate $\dot{C}_{\Delta, \text{loss}}(\rho)$ | find ρ with $\dot{C}_{\Delta, \text{loss}}(\rho) < 0$ in family |
| REG | [ASSUMED] | lower-bound spectrum of $\Delta[\rho]$ | sequence with $\Delta[\rho_n]$ nearly singular |
| BRIDGE | [ASSUMED] | compute $r(\rho)$ on the family | show $r(\rho)$ violates $O(f(\epsilon))$ despite CORR |
| RIP-U | [PROVED] | apply Definition 4.3 | refute any upstream assumption |
| $\omega = 0$ identity | [PROVED] | direct Dirichlet-form expansion | refute Davies form/hypotheses |
| RIP-L | [CONJECTURED] | estimate $\kappa^\downarrow(\epsilon)$ | construct ρ_ϵ with $r(\rho_\epsilon) \rightarrow 0$ |

Table 1: Falsification matrix for the dynamic layer.

A Line-by-line proof of the $\omega = 0$ identity

We give a direct expansion showing

$$\mathcal{E}_\sigma^{(0)}(O) = \frac{\gamma(0)}{2} \|[S(0), O]\|_{2, \sigma}^2.$$

For clarity we present the Hermitian case $S(0)^\dagger = S(0)$; the general case replaces $S(0)^2$ by $S(0)^\dagger S(0)$ and yields the same commutator structure.

The $\omega = 0$ Heisenberg dissipator has the standard GKLS form

$$\mathcal{L}^{\#, (0)}(O) = \gamma(0) \left(S(0)OS(0) - \frac{1}{2}\{S(0)^2, O\} \right).$$

Using the Dirichlet form convention from [Definition 5.2](#),

$$\mathcal{E}_\sigma^{(0)}(O) = -\mathrm{Tr}\left(\sigma^{1/2}O^\dagger\sigma^{1/2}\mathcal{L}^{\sharp,(0)}(O)\right).$$

Expand:

$$\mathcal{E}_\sigma^{(0)}(O) = -\gamma(0)\mathrm{Tr}\left(\sigma^{1/2}O^\dagger\sigma^{1/2}S(0)OS(0)\right) + \frac{\gamma(0)}{2}\mathrm{Tr}\left(\sigma^{1/2}O^\dagger\sigma^{1/2}S(0)^2O\right) + \frac{\gamma(0)}{2}\mathrm{Tr}\left(\sigma^{1/2}O^\dagger\sigma^{1/2}OS(0)^2\right).$$

On the other hand,

$$\begin{aligned}\|[S(0), O]\|_{2,\sigma}^2 &= \mathrm{Tr}\left(\sigma^{1/2}(S(0)O - OS(0))^\dagger\sigma^{1/2}(S(0)O - OS(0))\right) \\ &= \mathrm{Tr}\left(\sigma^{1/2}(O^\dagger S(0) - S(0)O^\dagger)\sigma^{1/2}(S(0)O - OS(0))\right).\end{aligned}$$

Using [Definition 5.4](#) to commute $S(0)$ through $\sigma^{1/2}$ and expanding the product yields

$$\|[S(0), O]\|_{2,\sigma}^2 = \mathrm{Tr}\left(\sigma^{1/2}O^\dagger\sigma^{1/2}S(0)^2O\right) - 2\mathrm{Tr}\left(\sigma^{1/2}O^\dagger\sigma^{1/2}S(0)OS(0)\right) + \mathrm{Tr}\left(\sigma^{1/2}O^\dagger\sigma^{1/2}OS(0)^2\right).$$

Comparing with the expression for $\mathcal{E}_\sigma^{(0)}(O)$ gives

$$\mathcal{E}_\sigma^{(0)}(O) = \frac{\gamma(0)}{2} \|[S(0), O]\|_{2,\sigma}^2,$$

as claimed.

References

- [1] E. B. Davies, Markovian master equations, *Commun. Math. Phys.* **39**, 91–110 (1974).
- [2] H.-P. Breuer and F. Petruccione, *The Theory of Open Quantum Systems*, Oxford University Press, 2002.
- [3] H. Spohn and J. L. Lebowitz, Irreversible thermodynamics for quantum systems weakly coupled to thermal reservoirs, *Adv. Chem. Phys.* **38**, 109–142 (1978).