

# Finite-Dimensional Davies Interface Lemmas and TFIM Witness Tests

## for the Heisenberg Cut as a Resource Boundary

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December 2025

### Status of claims (read first)

- **Imported result.** The maintenance inequality  $P_{\text{extra}}(\rho) \geq k_{\text{B}}T \dot{C}_{\text{loss}}(\rho)$  is imported from the published preprint [2].
- **Proved here (finite-dimensional).** The exact  $\omega = 0$  Dirichlet identity (Lemma 1), the Bohr-block Dirichlet decomposition (Lemma 2), and the witness lower bound (Proposition 1).
- **Sufficient conditional bound.** The envelope suppression lemma (Lemma 3) is a finite-dimensional sufficient statement under explicit hypotheses (IR exclusion + quasi-local spectral tails).
- **Numerics.** TFIM results are exact diagonalization and stability-checked finite-size diagnostics (here  $N \leq 12$ ); no thermodynamic-limit uniformity is claimed.

### Abstract

We present a finite-dimensional technical core for interpreting the quantum–classical boundary as a control-resource limitation: a state is operationally “classical” (relative to a chosen conditional expectation) whenever sustaining coherence is infeasible under a finite available power budget.

Using battery-assisted thermal operations at temperature  $T$  and energy pinching  $\Delta$ , we quantify coherence by  $C(\rho) = S(\rho \parallel \Delta[\rho])$  and define its instantaneous loss rate  $\dot{C}_{\text{loss}}(\rho)$  under uncontrolled Markovian dynamics. An imported maintenance inequality yields the operational bound  $P_{\text{extra}}(\rho) \geq k_{\text{B}}T \dot{C}_{\text{loss}}(\rho)$ .

We isolate the dynamical hinge required for geometric scaling claims: the relation between geometric separation and effective decoherence rates. In Davies generators we give two interface lemmas: a negative lemma showing that an active  $\omega = 0$  channel yields a witness-based lower bound via a KMS commutator with  $S(0)$ , and a positive sufficient lemma yielding envelope suppression under explicit infrared exclusion plus a quasi-local spectral tail hypothesis.

Finally, we provide TFIM exact-diagonalization witness protocols computing ratios  $R(\epsilon)$  and optimized local witnesses  $R_{\text{opt}}(\epsilon)$ , including scaling tests with  $\epsilon = \epsilon(N)$  and temperature sweeps over  $\beta$ . We also report direct state-side interface micro-tests on weak-coherence families, confirming plateau behavior of  $\dot{C}_{\text{loss}}/C$  and controlled  $\omega = 0$  scaling in  $\gamma(0)$  and  $\beta$ , with stability checks in the Euler step size and PSD clipping threshold.

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# 1 Scope and non-claims

This work does not propose modified quantum dynamics, collapse mechanisms, derivations of the Born rule, or ontological resolutions of the measurement problem. The claims are operational: they concern minimal control resources required to maintain coherence against uncontrolled open-system dynamics under an explicit thermodynamic bookkeeping model.

**Remark 1** (Figures). *This manuscript expects several figure files (e.g. `scaling_Ropt_N_over_4.png`, `interface_ratio-vs-C.png`) to be present in ‘figures/’ (or the working directory). If they are missing, compilation will fail.*

Table 1: Program Map: logical structure of the coherence-maintenance series.

ID	Role	Key Contribution
<i>Foundational Preprints (Published)</i>		
<b>P1</b>	Geometric Foundation	AQFT static clustering bounds and Petz recovery.
<b>P2</b>	Thermodynamic Law	Maintenance inequality: $P_{\text{extra}}(\rho) \geq k_{\text{B}}T \dot{C}_{\text{loss}}(\rho)$ .
<b>P3</b>	Synthesis/Outlook	Operational cut framing and interface discussion.
<i>Core Technical Paper (This Work)</i>		
<b>P7</b>	<b>Finite-Dimensional Core</b>	<b>Davies interface lemmas, TFIM witnesses, and reproducible micro-tests.</b>
<i>Companion Notes (to be posted separately)</i>		
<b>P4</b>	Hypothesis Framing	RIP: weak vs. strong forms, plausibility routes, and falsification criteria.
<b>P6</b>	Stress Tests	Operatorial $\kappa(\epsilon)$ tests and $\omega = 0$ -floor mechanisms.
<b>P5</b>	Logical Closure	Separation of proved results vs. conditional interfaces vs. open dynamical gaps.

## 2 Finite-dimensional operational core

### 2.1 Thermal state, energy pinching, and coherence

Let the system Hamiltonian be

$$H_S = \sum_n E_n \Pi_n,$$

and define at temperature  $T$  the thermal state

$$\gamma_S := \frac{e^{-\beta H_S}}{\text{Tr}(e^{-\beta H_S})}, \quad \beta = \frac{1}{k_{\text{B}}T}.$$

Define energy pinching

$$\Delta[\rho] := \sum_n \Pi_n \rho \Pi_n,$$

and relative-entropy coherence (nats)

$$C(\rho) := S(\rho \parallel \Delta[\rho]), \quad S(\rho \parallel \sigma) := \text{Tr}[\rho(\log \rho - \log \sigma)].$$

## 2.2 Uncontrolled Markovian dynamics and coherence-loss rate

Let  $\rho_t = e^{t\mathcal{L}}(\rho)$  be a Markovian semigroup with a thermal stationary state  $\mathcal{L}(\gamma_S) = 0$ . Define the instantaneous coherence-loss rate

$$\dot{C}_{\text{loss}}(\rho) := - \left. \frac{d}{dt} C(\rho_t) \right|_{t=0}.$$

## 2.3 Control model and extra power

We assume the control primitive of battery-assisted thermal operations at temperature  $T$ , and define incremental (extra) power  $P_{\text{extra}}$  as an infimum over pairs of strategies maintaining  $\rho$  and maintaining  $\Delta[\rho]$ , respectively.

## 2.4 Core maintenance inequality (imported)

**Theorem 1** (Extra-power coherence maintenance bound (imported)). *Assume strategies exist that maintain  $\rho$  and that maintain  $\Delta[\rho]$  under battery-assisted thermal operations at temperature  $T$ . Then*

$$P_{\text{extra}}(\rho) \geq k_{\text{B}}T \dot{C}_{\text{loss}}(\rho).$$

**Corollary 1** (No-go from a uniform loss floor). *Fix  $C_0 > 0$  and a nonempty target family  $\mathcal{F}_{C_0} \subset \{\rho : C(\rho) \geq C_0\}$ . If there exists*

$$g(C_0) := \inf_{\rho \in \mathcal{F}_{C_0}} \dot{C}_{\text{loss}}(\rho) > 0,$$

*then any controller maintaining any  $\rho \in \mathcal{F}_{C_0}$  requires  $P_{\text{extra}}(\rho) \geq k_{\text{B}}T g(C_0)$ . If  $P_{\text{max}} < k_{\text{B}}T g(C_0)$ , sustained maintenance is operationally impossible on  $\mathcal{F}_{C_0}$ .*

# 3 Davies generators, KMS Dirichlet form, and an effective envelope

## 3.1 Bohr components

Consider a Davies generator derived in the weak-coupling–secular limit with a single system coupling operator  $S = S^\dagger$ . Define Bohr components

$$S(\omega) := \sum_{E_m - E_n = \omega} \Pi_m S \Pi_n, \quad S = \sum_{\omega} S(\omega), \quad S(\omega)^\dagger = S(-\omega).$$

Bath rates satisfy the KMS relation

$$\gamma(-\omega) = e^{-\beta\omega} \gamma(\omega).$$

## 3.2 KMS inner product and Dirichlet form

Define the KMS (GNS) inner product

$$\langle A, B \rangle_\sigma := \text{Tr} \left( \sigma^{1/2} A^\dagger \sigma^{1/2} B \right), \quad \|A\|_{2,\sigma}^2 := \langle A, A \rangle_\sigma,$$

with  $\sigma = \gamma_S$ . For the Heisenberg-picture generator  $\mathcal{L}^\dagger$ , define the Dirichlet form

$$\mathcal{E}_\sigma(O) := -\text{Re} \langle O, \mathcal{L}^\dagger(O) \rangle_\sigma.$$

### 3.3 Separation and variational envelope

Fix a coupling region (e.g. a site  $j_0$ ) and let  $\mathcal{A}_\epsilon$  denote operators supported at distance at least  $\epsilon$  from that region. Define the operatorial envelope

$$\kappa(\epsilon) := \sup_{O \in \mathcal{A}_\epsilon, O \neq 0} \frac{\mathcal{E}_\sigma(O)}{\|O\|_{2,\sigma}^2}.$$

### 3.4 A concrete tail map on spin chains (Pauli-string conditional expectation)

In lattice spin systems we use a canonical “tail” map built from the Pauli-string expansion. Let  $\Lambda$  be the set of lattice sites and let  $\mathcal{A}(\Omega)$  denote the local operator algebra on region  $\Omega \subset \Lambda$ .

Fix a coupling region  $\Omega_0$  (e.g.  $\Omega_0 = \{j_0\}$ ). For  $\epsilon \geq 1$  define the far region

$$\Omega_\epsilon := \{j \in \Lambda : d(j, \Omega_0) \geq \epsilon\}.$$

Let  $\mathcal{A}_{<\epsilon} := \mathcal{A}(\Lambda \setminus \Omega_\epsilon)$  be the algebra supported within distance  $< \epsilon$  of the coupling region.

Define the *Pauli-string conditional expectation*  $\mathbb{E}_{<\epsilon} : \mathcal{A}(\Lambda) \rightarrow \mathcal{A}_{<\epsilon}$  as the orthogonal projection (with respect to the normalized Hilbert–Schmidt inner product) onto the span of Pauli strings supported in  $\Lambda \setminus \Omega_\epsilon$ . Equivalently, writing any operator as a Pauli expansion  $X = \sum_P x_P P$  with Pauli strings  $P$ ,

$$\mathbb{E}_{<\epsilon}[X] := \sum_{P: \text{supp}(P) \subseteq \Lambda \setminus \Omega_\epsilon} x_P P.$$

We then define the tail at separation  $\epsilon$  by

$$\text{tail}_\epsilon(X) := X - \mathbb{E}_{<\epsilon}[X].$$

**Remark 2.** We define  $\mathbb{E}_{<\epsilon}$  as a normalized Hilbert–Schmidt projection for concreteness and computational convenience. Since the system is finite-dimensional,  $\|\cdot\|_{2,\sigma}$  and the Hilbert–Schmidt norm are equivalent; thus any exponential tail bound in one induces an exponential tail bound in the other up to multiplicative constants depending on  $\sigma$ .

These constants may depend on  $(N, \beta)$  in finite-size tests. Alternatively, one may define  $\mathbb{E}_{<\epsilon}$  as the orthogonal projection with respect to  $\langle \cdot, \cdot \rangle_\sigma$  to align norms more tightly.

## 4 Interface lemmas: rate floors and conditional envelope suppression

### 4.1 Negative lemma: the $\omega = 0$ channel yields a witness lower bound

Define the  $\omega = 0$  Bohr component

$$S(0) := \sum_n \Pi_n S \Pi_n.$$

Assume  $\gamma(0) > 0$  (e.g. after infrared regularization in an Ohmic model).

**Lemma 1** (Exact  $\omega = 0$  Dirichlet identity). *If the  $\omega = 0$  Davies channel is*

$$\mathcal{L}_0^\dagger(O) = \gamma(0) \left( S(0) O S(0) - \frac{1}{2} \{S(0)^2, O\} \right),$$

then for all  $O$ ,

$$\mathcal{E}_\sigma^{(0)}(O) := -\text{Re} \langle O, \mathcal{L}_0^\dagger(O) \rangle_\sigma = \frac{\gamma(0)}{2} \|[S(0), O]\|_{2,\sigma}^2.$$

*Proof.* Let  $\sigma = \gamma_S$  and set  $S_0 := S(0)$ . Since  $S_0$  is diagonal in the energy eigenbasis, we have  $[S_0, H_S] = 0$  and hence  $[S_0, \sigma] = 0$ , i.e.  $\sigma^{1/2}S_0 = S_0\sigma^{1/2}$ .

Write

$$\mathcal{L}_0^\dagger(O) = \gamma(0) \left( S_0 O S_0 - \frac{1}{2} \{S_0^2, O\} \right).$$

Using  $\langle A, B \rangle_\sigma = \text{Tr}(\sigma^{1/2} A^\dagger \sigma^{1/2} B)$  and  $[S_0, \sigma^{1/2}] = 0$ , one expands both  $\langle O, \mathcal{L}_0^\dagger(O) \rangle_\sigma$  and  $\|[S_0, O]\|_{2,\sigma}^2$  and matches terms, obtaining

$$-\text{Re} \langle O, \mathcal{L}_0^\dagger(O) \rangle_\sigma = \frac{\gamma(0)}{2} \|[S_0, O]\|_{2,\sigma}^2.$$

□

**Proposition 1** (Fixed-separation witness lower bound (sufficient)). *Fix  $\epsilon_0 \geq 1$ . If there exists  $c_{\epsilon_0} > 0$  and an operator  $O_{\epsilon_0} \in \mathcal{A}_{\epsilon_0}$  such that*

$$\frac{\|[S(0), O_{\epsilon_0}]\|_{2,\sigma}^2}{\|O_{\epsilon_0}\|_{2,\sigma}^2} \geq c_{\epsilon_0},$$

then

$$\kappa(\epsilon_0) \geq \frac{\gamma(0)}{2} c_{\epsilon_0}.$$

## 4.2 Bohr-block decomposition of the KMS Dirichlet form

**Lemma 2** (Bohr-block decomposition of the KMS Dirichlet form (single-channel Davies)). *Assume a single Hermitian coupling operator  $S = S^\dagger$  and a Davies generator in the Heisenberg picture of the form*

$$\mathcal{L}^\dagger(O) = \sum_\omega \gamma(\omega) \left( S(\omega)^\dagger O S(\omega) - \frac{1}{2} \{S(\omega)^\dagger S(\omega), O\} \right), \quad S(\omega)^\dagger = S(-\omega),$$

with rates satisfying the KMS relation  $\gamma(-\omega) = e^{-\beta\omega} \gamma(\omega)$  and reference state  $\sigma = \gamma_S$ . Let

$$\langle A, B \rangle_\sigma := \text{Tr}(\sigma^{1/2} A^\dagger \sigma^{1/2} B), \quad \mathcal{E}_\sigma(O) := -\text{Re} \langle O, \mathcal{L}^\dagger(O) \rangle_\sigma.$$

Then

$$\mathcal{E}_\sigma(O) = \frac{1}{2} \sum_\omega \gamma(\omega) \|[S(\omega), O]\|_{2,\sigma}^2.$$

*Proof.* Fix  $\omega$  and set  $L := S(\omega)$  and  $\gamma := \gamma(\omega)$ . Define

$$\mathcal{L}_\omega^\dagger(O) := \gamma \left( L^\dagger O L - \frac{1}{2} \{L^\dagger L, O\} \right).$$

Since  $\sigma = \gamma_S \propto e^{-\beta H_S}$  and  $L$  has Bohr frequency  $\omega$ , the intertwining relations hold:

$$\sigma^{1/2} L = e^{-\beta\omega/2} L \sigma^{1/2}, \quad \sigma^{1/2} L^\dagger = e^{+\beta\omega/2} L^\dagger \sigma^{1/2}.$$

Using these relations and cyclicity of trace, one rewrites  $-\text{Re} \langle O, \mathcal{L}_\omega^\dagger(O) \rangle_\sigma$  as  $\frac{\gamma}{2} \|[L, O]\|_{2,\sigma}^2$ . Summing over  $\omega$  gives the claim. □

### 4.3 Positive lemma: envelope suppression under IR exclusion and quasi-local spectral tails

**Definition 1** (Infrared exclusion). *We say the coupling satisfies IR exclusion at scale  $\omega_* > 0$  if  $S(\omega) = 0$  for all  $|\omega| < \omega_*$  (or alternatively  $\gamma(\omega) = 0$  on that band).*

**Lemma 3** (Envelope suppression under IR exclusion and quasi-local spectral tails (sufficient)). *Assume IR exclusion and suppose there exist  $A > 0$ ,  $\xi > 0$ , and  $p \geq 0$  such that for all  $|\omega| \geq \omega_*$ ,*

$$\|\text{tail}_\epsilon(S(\omega))\|_{2,\sigma} \leq A(1 + \epsilon)^p e^{-\epsilon/\xi},$$

and  $\sup_{|\omega| \geq \omega_*} \gamma(\omega) < \infty$ . Then there exists  $C > 0$  such that

$$\kappa(\epsilon) \leq C(1 + \epsilon)^{2p} e^{-2\epsilon/\xi}.$$

*Proof (sketch).* By Lemma 2,

$$\mathcal{E}_\sigma(O) = \frac{1}{2} \sum_{\omega} \gamma(\omega) \|[S(\omega), O]\|_{2,\sigma}^2.$$

By IR exclusion, only  $|\omega| \geq \omega_*$  contribute. For  $O \in \mathcal{A}_\epsilon$ , write

$$S(\omega) = \mathbb{E}_{<\epsilon}[S(\omega)] + \text{tail}_\epsilon(S(\omega)).$$

Since  $\mathbb{E}_{<\epsilon}[S(\omega)]$  is supported within distance  $< \epsilon$  while  $O$  is supported at distance at least  $\epsilon$ ,

$$[\mathbb{E}_{<\epsilon}[S(\omega)], O] = 0, \quad \Rightarrow \quad [S(\omega), O] = [\text{tail}_\epsilon(S(\omega)), O].$$

Then

$$\|[A, O]\|_{2,\sigma} \leq \|AO\|_{2,\sigma} + \|OA\|_{2,\sigma} \leq 2\|A\|_\infty \|O\|_{2,\sigma},$$

and hence

$$\frac{\|[S(\omega), O]\|_{2,\sigma}^2}{\|O\|_{2,\sigma}^2} \leq 4 \|\text{tail}_\epsilon(S(\omega))\|_\infty^2.$$

Using  $\|X\|_\infty \leq \|X\|_2 \leq \lambda_{\min}(\sigma)^{-1/2} \|X\|_{2,\sigma}$  yields

$$\|\text{tail}_\epsilon(S(\omega))\|_\infty^2 \leq \frac{1}{\lambda_{\min}(\sigma)} \|\text{tail}_\epsilon(S(\omega))\|_{2,\sigma}^2.$$

Combining the bounded-rate assumption with the tail hypothesis gives the stated envelope for  $\kappa(\epsilon)$  after taking the supremum over  $O \in \mathcal{A}_\epsilon$ .  $\square$

**Remark 3** (Finite-size scaling of the prefactor). *The constant  $C$  contains a factor  $\lambda_{\min}(\sigma)^{-1}$  arising from norm comparisons and may grow with  $(N, \beta)$  since  $\sigma = \gamma_S \propto e^{-\beta H_S}$ . Accordingly, Lemma 3 is stated as a sufficient finite-dimensional bound at fixed  $(N, \beta)$  and is not claimed to be uniform in any thermodynamic limit.*

## 5 TFIM witness computations (exact diagonalization)

### 5.1 Model

We consider the transverse-field Ising model on  $N$  spins with open boundary conditions:

$$H_S = -J \sum_{i=1}^{N-1} \sigma_i^x \sigma_{i+1}^x - h \sum_{i=1}^N \sigma_i^z, \quad h > J.$$

We choose a coupling site  $j_0 = \lfloor N/2 \rfloor$  and take  $S = \sigma_{j_0}^z$ . We consider thermal states  $\sigma = \gamma_S$  at inverse temperatures  $\beta \in \{0.5, 1, 2\}$ .

### 5.2 Witness ratios

For a distant site  $k = j_0 \pm \epsilon$  define

$$R(\epsilon; O) := \frac{\| [S(0), O] \|_{2,\sigma}^2}{\| O \|_{2,\sigma}^2},$$

and define the optimized local witness ratio

$$R_{\text{opt}}(\epsilon) := \max_{O \in \text{span}\{\sigma_k^x, \sigma_k^y, \sigma_k^z\}} R(\epsilon; O).$$

By Lemma 1,

$$\kappa(\epsilon) \geq \frac{\gamma(0)}{2} R_{\text{opt}}(\epsilon).$$

### 5.3 Fixed $\epsilon$ results (at $\beta = 1$ )

For  $J = 1$ ,  $h = 1.5$ ,  $\beta = 1$ :

$N$	$R_{\text{opt}}(1)$	$R_{\text{opt}}(2)$	$R_{\text{opt}}(3)$
8	0.134648	0.0885503	0.0684289
10	0.0999406	0.0696450	0.0550748
12	0.0793451	0.0574518	0.0464596

Table 2: Optimized local KMS commutator witness  $R_{\text{opt}}(\epsilon)$  for TFIM with  $S = \sigma_{j_0}^z$ , center coupling,  $\beta = 1$ ,  $h = 1.5$ ,  $J = 1$ .

### 5.4 Scaling tests with $\epsilon = \epsilon(N)$ (at $\beta = 1$ )

We computed  $R_{\text{opt}}(\epsilon(N))$  for four rules:

$$\epsilon(N) \in \left\{ \left\lfloor \frac{N}{4} \right\rfloor, \left\lfloor \frac{N}{3} \right\rfloor, \epsilon_{\text{max}}(N), 3 \right\}, \quad \epsilon_{\text{max}}(N) = \min(j_0, N - 1 - j_0).$$

### 5.5 Temperature sweep tests over $\beta$

We swept  $\beta \in \{0.5, 1, 2\}$  for the same  $\epsilon(N)$  rules and  $N \in \{8, 10, 12\}$ .

Rule	$N = 8$	$N = 10$	$N = 12$
$\lfloor N/4 \rfloor$	0.0885503 ( $\epsilon = 2$ )	0.0696450 ( $\epsilon = 2$ )	0.0464596 ( $\epsilon = 3$ )
$\lfloor N/3 \rfloor$	0.0885503 ( $\epsilon = 2$ )	0.0550748 ( $\epsilon = 3$ )	0.0396591 ( $\epsilon = 4$ )
$\epsilon_{\max}(N)$	0.0684289 ( $\epsilon = 3$ )	0.0461050 ( $\epsilon = 4$ )	0.0330550 ( $\epsilon = 5$ )
3 (fixed)	0.0684289 ( $\epsilon = 3$ )	0.0550748 ( $\epsilon = 3$ )	0.0464596 ( $\epsilon = 3$ )

Table 3: Scaling test values  $R_{\text{opt}}(\epsilon(N))$  for TFIM with  $S = \sigma_{j_0}^z$ ,  $\beta = 1$ ,  $h = 1.5$ ,  $J = 1$ .

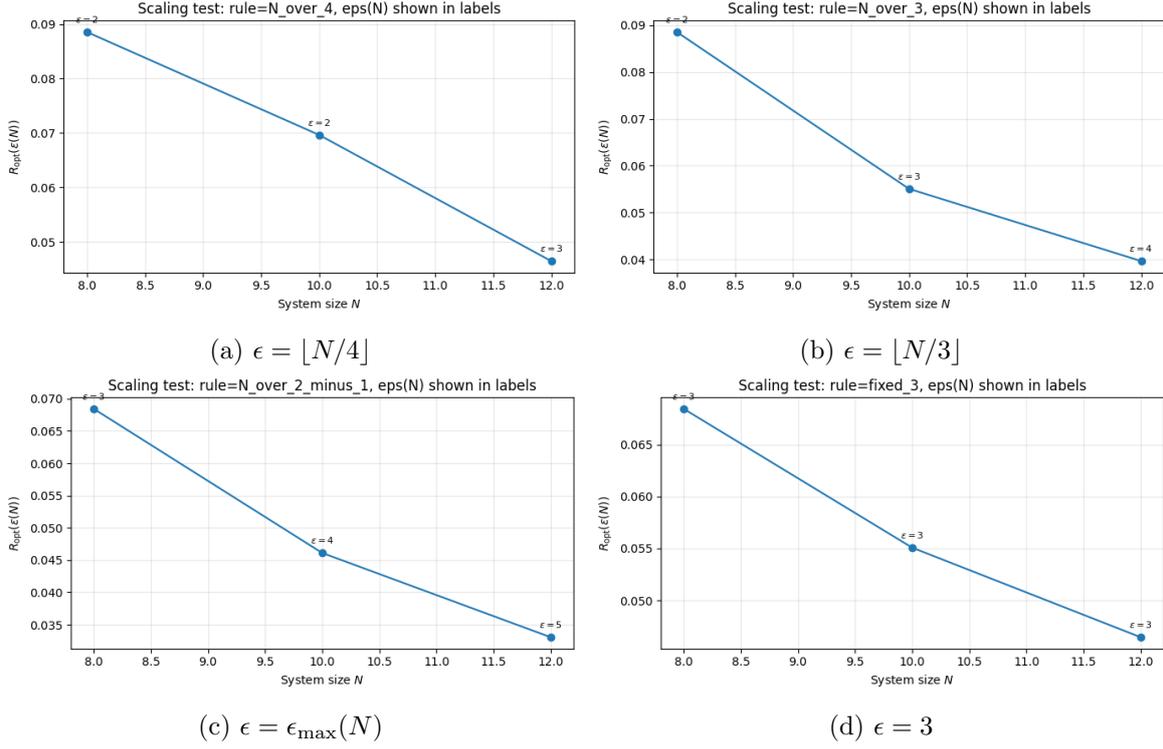


Figure 1: Scaling plots of  $R_{\text{opt}}(\epsilon(N))$  under four  $\epsilon(N)$  rules (at  $\beta = 1$ ).

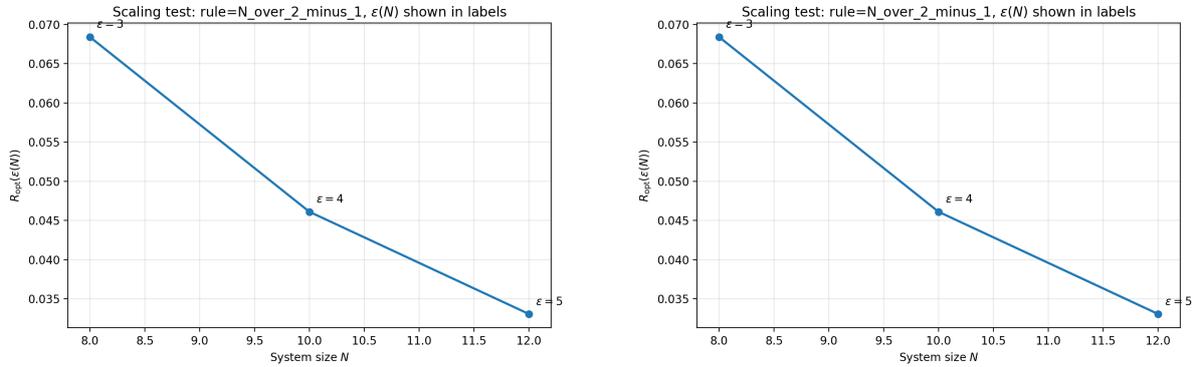


Figure 2: Scaling tests where  $\epsilon(N)$  is printed on each data point.

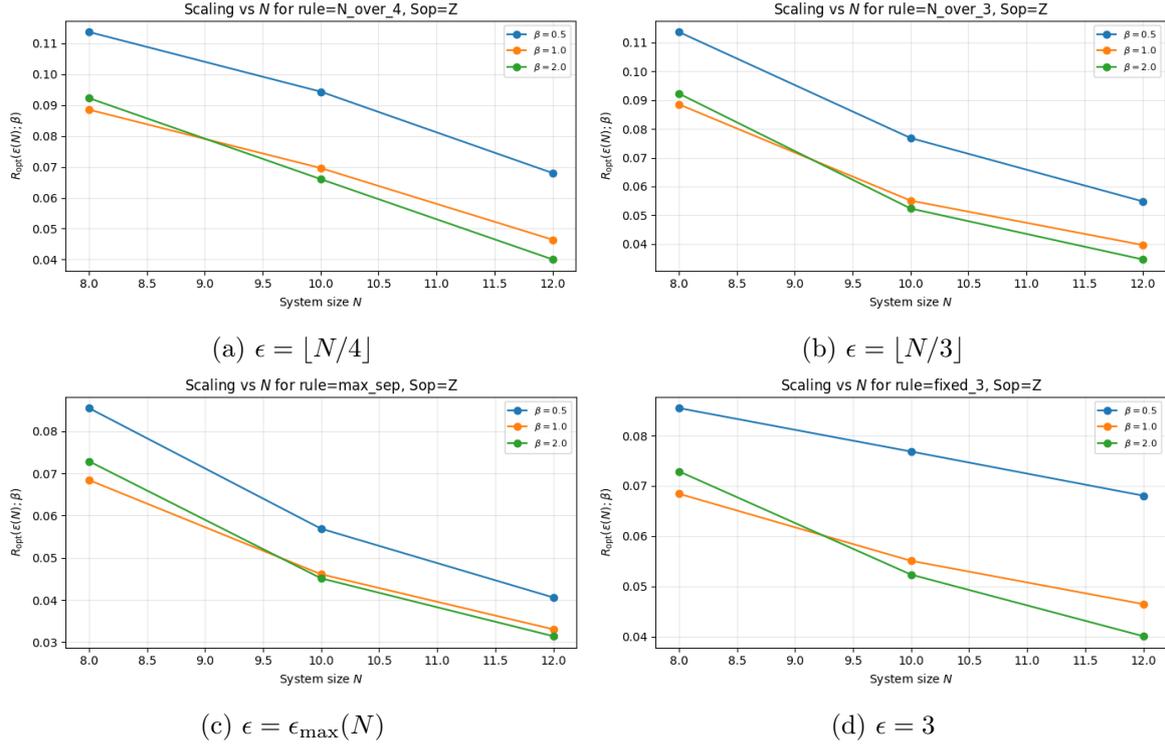


Figure 3: Optimized witness  $R_{\text{opt}}(\epsilon(N); \beta)$  versus  $N$  for multiple  $\beta$  values.

## 6 State-side interface micro-tests: weak-coherence families and $\omega = 0$ scaling

The witness computations above are observable-side tests based on the exact  $\omega = 0$  Dirichlet identity. To directly probe the state-side quantity appearing in the maintenance inequality, we also test the interface numerically on explicit weak-coherence families.

**Weak-coherence family.** Fix an energy eigenbasis  $\{n\}$  of  $H_S$ , thermal weights  $p_n \propto e^{-\beta E_n}$ , and a pair  $(m, n)$  with  $m \neq n$ . Define in the energy basis

$$\rho_E(\alpha) := \text{diag}(p) + \alpha \sqrt{p_m p_n} (nm + mn),$$

and set  $\rho(\alpha) = V \rho_E(\alpha) V^\dagger$  where  $V$  diagonalizes  $H_S$ . For small  $\alpha$ , this is positive and satisfies  $\Delta[\rho(\alpha)] = \text{diag}(p)$  by construction. We evaluate  $C(\rho) = S(\rho \| \Delta[\rho])$  and approximate

$$\dot{C}_{\text{loss}}(\rho) \approx -\frac{C(\rho_t) - C(\rho)}{t}, \quad \rho_t \approx \Pi(\rho + t \mathcal{L}(\rho)),$$

where  $\Pi$  denotes numerical projection to a valid density matrix (Hermitian, PSD with eigenvalue clipping, trace one). We report the empirical ratio

$$r(\rho) := \frac{\dot{C}_{\text{loss}}(\rho)}{C(\rho)}.$$

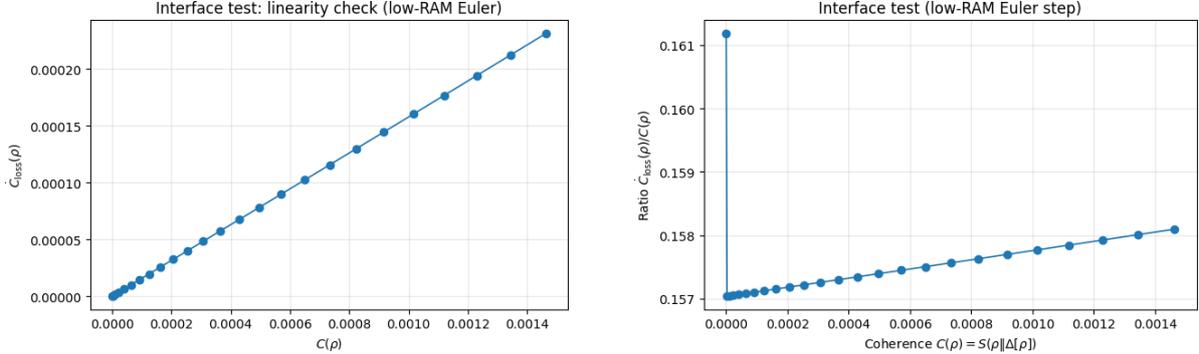
**Observed plateau (truncated Davies blocks).** For  $N = 6$ ,  $J = 1$ ,  $h = 1.5$ ,  $\beta = 1$ , coupling  $S = \sigma_{j_0}^z$  at the center site, and a transition-selected pair optimized for a nonzero Bohr block, we observed that  $r(\rho)$  is approximately constant across several decades of  $C(\rho)$  (excluding ultra-small  $C$  where PSD projection dominates), with a plateau around

$$r(\rho) \approx 0.157\text{--}0.158$$

for the  $\{\pm\omega\}$  truncation, and a slightly larger plateau

$$r(\rho) \approx 0.160\text{--}0.161$$

when an additional  $\omega = 0$  channel with  $\gamma(0) = 0.2$  is included.



(a) Linearity check:  $\dot{C}_{\text{loss}}(\rho)$  versus  $C(\rho)$ .

(b) Ratio  $r(\rho) = \dot{C}_{\text{loss}}(\rho)/C(\rho)$  versus  $C(\rho)$ .

Figure 4: State-side interface micro-test (Euler step with PSD projection). Ultra-small  $C(\rho)$  points can be affected by eigenvalue clipping; the plateau region is stable.

**$\omega = 0$ -only regime: linear scaling in  $\gamma(0)$ .** To isolate the commutator-squared mechanism of the negative lemma, we ran an  $\omega = 0$ -only truncation and chose  $(m, n)$  to maximize  $|s_m - s_n| \sqrt{p_m p_n}$  where  $s_k := kSk$  in the energy basis. With  $\beta = 1$  and numerical prescription  $\gamma(0) = \gamma_0/\beta$ , the measured plateau ratio  $r(\alpha = 0.15)$  scales linearly with  $\gamma(0)$ :

$\gamma_0$	$\gamma(0) = \gamma_0/\beta$	$r(\alpha = 0.15)$	$r/\gamma(0)$
0.05	0.05	0.0101957	0.2039
0.10	0.10	0.0203914	0.2039
0.40	0.40	0.0815656	0.2039

Table 4: TFIM  $\omega = 0$ -only micro-test ( $N = 6$ ,  $J = 1$ ,  $h = 1.5$ ,  $\beta = 1$ ). The empirical ratio  $r = \dot{C}_{\text{loss}}/C$  scales linearly in  $\gamma(0) = \gamma_0/\beta$ .

**Fixed-pair temperature scaling:  $r(\beta) \propto 1/\beta$  at fixed  $\gamma_0$ .** To avoid conflating temperature dependence with changes in the coherence sector, we fixed  $(m, n) = (0, 1)$  and kept  $\gamma_0 = 0.10$  fixed (so  $\gamma(0) = \gamma_0/\beta$ ). The measured plateau satisfies  $r(\beta) \approx k/\beta$ :

$\beta$	$\gamma(0) = \gamma_0/\beta$ ( $\gamma_0 = 0.10$ )	$r(\alpha = 0.15)$	$\beta r(\alpha = 0.15)$
0.5	0.20	0.004317824675	0.002158912338
1.0	0.10	0.002158560224	0.002158560224
2.0	0.05	0.001078731632	0.002157463263
4.0	0.025	0.000538771757	0.002155087026

Table 5: Fixed-pair  $\omega = 0$ -only temperature sweep ( $N = 6$ ,  $J = 1$ ,  $h = 1.5$ ,  $(m, n) = (0, 1)$ ,  $\gamma_0 = 0.10$ ). The near-constancy of  $\beta r$  confirms  $r(\beta) \propto 1/\beta$  at fixed  $\gamma_0$ .

## 7 State–observable interface (explicit hypothesis)

The maintenance inequality bounds power in terms of the state functional  $\dot{C}_{\text{loss}}(\rho)$ , whereas  $\kappa(\epsilon)$  is defined via a supremum over observables in  $\mathcal{A}_\epsilon$ . To avoid a quantifier mismatch, we state the necessary bridge as an explicit hypothesis.

**Definition 2** (Weak-coherence regime on a family). *A family of states  $\mathcal{F}_\epsilon$  lies in the weak-coherence regime if there exists  $\delta_\epsilon > 0$  such that  $C(\rho) \leq \delta_\epsilon$  for all  $\rho \in \mathcal{F}_\epsilon$ .*

**Definition 3** (Population non-extremality on a family). *A family  $\mathcal{F}_\epsilon$  is population-nonextremal if there exists  $p_{\min} > 0$  such that for all  $\rho \in \mathcal{F}_\epsilon$ ,*

$$\Delta[\rho] \geq p_{\min} \mathbf{1}.$$

**Remark 4.** *This assumption rules out diagonals arbitrarily close to the boundary of the simplex, ensuring uniform local expansions of relative entropy around  $\Delta[\rho]$  and preventing degenerate small-denominator effects in weak-coherence estimates.*

**Definition 4** (Effective coherence sector). *Fix a set  $\Omega$  of Bohr frequencies. We say  $\rho$  has coherence supported in  $\Omega$  if all its off-diagonal matrix elements in the energy basis lie in frequency blocks  $\omega \in \Omega$ .*

**Definition 5** (Sector rate lower bound). *For a Davies generator, let  $\Gamma(\omega)$  denote the Kosakowski rate matrix in the Bohr block  $\omega$ . In the single-channel case,  $\Gamma(\omega)$  is a scalar and we identify  $\lambda_{\min}(\Gamma(\omega)) = \Gamma(\omega)$ . Define*

$$\underline{\Gamma}_\Omega := \min_{\omega \in \Omega} \lambda_{\min}(\Gamma(\omega)).$$

**Hypothesis 1** (State–observable interface (lower proportionality on a family)). *Fix  $\epsilon$  and a target family  $\mathcal{F}_\epsilon$ . Assume:*

1.  $\mathcal{F}_\epsilon$  lies in a weak-coherence regime;
2.  $\mathcal{F}_\epsilon$  is population-nonextremal;
3. all  $\rho \in \mathcal{F}_\epsilon$  have coherence supported in a fixed sector set  $\Omega$ ;
4. the Davies generator is validated in the regime of interest.

*Then there exists  $\eta_\epsilon \rightarrow 0$  as the coherence scale  $\delta_\epsilon \rightarrow 0$  such that for all  $\rho \in \mathcal{F}_\epsilon$ ,*

$$\dot{C}_{\text{loss}}(\rho) \geq (1 - \eta_\epsilon) \underline{\Gamma}_\Omega C(\rho).$$

**Remark 5.** *This interface is falsifiable: one can numerically test whether  $\dot{C}_{\text{loss}}(\rho)/C(\rho)$  is bounded below on the specified family.*

## 8 Operational consequence (conditional, well-typed)

Combining Theorem 1 with Hypothesis 1 yields, for  $\rho \in \mathcal{F}_\epsilon$ ,

$$P_{\text{extra}}(\rho) \geq k_{\text{B}}T \dot{C}_{\text{loss}}(\rho) \geq k_{\text{B}}T(1 - \eta_\epsilon) \underline{\Gamma}_\Omega C(\rho).$$

In particular, if  $C(\rho) \geq C_0$  on  $\mathcal{F}_\epsilon$ , then

$$P_{\text{extra}}(\rho) \geq k_{\text{B}}T(1 - \eta_\epsilon) \underline{\Gamma}_\Omega C_0.$$

## 9 Conclusion

We separated a proven finite-dimensional thermodynamic maintenance bound from conditional dynamical interface claims. In Davies dynamics, an exact  $\omega = 0$  identity yields operational witness-based lower bounds. In TFIM witness computations,  $R_{\text{opt}}(\epsilon)$  is robust and non-negligible at fixed separations, while decreasing as  $\epsilon$  scales with  $N$  over the tested sizes.

We emphasize that the TFIM numerics ( $N \leq 12$ ) are presented as finite-size witness diagnostics, not as definitive thermodynamic-limit scaling evidence.

## A Numerical stability checks: time step and PSD clipping

We estimate  $\dot{C}_{\text{loss}}(\rho)$  via a forward Euler step  $\rho_t \approx \Pi(\rho + t\mathcal{L}(\rho))$  with PSD projection  $\Pi$  (eigenvalue clipping) and a finite-difference quotient.

$dt$	$r$ (last sample)
$1e - 4$	0.1613363453
$5e - 5$	0.1613377388
$1e - 5$	0.1613385899

Table 6: Sensitivity of the plateau ratio  $r = \dot{C}_{\text{loss}}/C$  to the Euler step size  $dt$  (interface micro-test).

psd_clip	$r$ (last sample)
$1e - 12$	0.1613377388
$1e - 14$	0.1613377388
$1e - 16$	0.1613377388

Table 7: Sensitivity of the plateau ratio  $r = \dot{C}_{\text{loss}}/C$  to the PSD projection threshold (interface micro-test).

## B Imported maintenance inequality: precise pointer map

The main maintenance inequality used in this manuscript is

$$P_{\text{extra}}(\rho; \mathcal{L}, T) \geq k_{\text{B}}T \dot{C}_{\text{loss}}(\rho), \quad \dot{C}_{\text{loss}}(\rho) := - \left. \frac{d}{dt} S(\rho_t \| \Delta[\rho_t]) \right|_{t=0}, \quad \rho_t = e^{t\mathcal{L}}(\rho).$$

A complete finite-dimensional proof is given in the published preprint [2].

## C Reproducibility: Colab-friendly scripts (included inline)

This appendix includes scripts used to generate the state-side interface micro-tests (including figures) and the  $\omega = 0$  fixed-pair  $\beta$  sweep.

### C.1 Low-RAM interface micro-test script (Davies blocks; optional $\omega = 0$ ; pair selection)

```
# file: test_interface_hypothesis_tfim_lowram.py
import argparse
import numpy as np
import scipy.linalg as la
import matplotlib.pyplot as plt

def parse_args():
    ap = argparse.ArgumentParser()
    ap.add_argument("--N", type=int, default=6)
    ap.add_argument("--J", type=float, default=1.0)
    ap.add_argument("--h", type=float, default=1.5)
    ap.add_argument("--beta", type=float, default=1.0)

    ap.add_argument("--Sop", type=str, default="Z", choices=["X", "Y", "Z"])
    ap.add_argument("--center", action="store_true")
    ap.add_argument("--j0", type=int, default=None)

    ap.add_argument("--dt", type=float, default=5e-5)
    ap.add_argument("--nsamples", type=int, default=25)
    ap.add_argument("--alpha_max", type=float, default=0.15)

    ap.add_argument("--bin_tol", type=float, default=1e-10)
    ap.add_argument("--gamma0", type=float, default=0.2)
    ap.add_argument("--s", type=float, default=1.0)

    ap.add_argument("--use_omega0", action="store_true")
    ap.add_argument("--only_omega0", action="store_true")
    ap.add_argument("--only_pm_omega", action="store_true")

    ap.add_argument("--pair_mode", type=str, default="pmomega", choices=["pmomega", "omega0"])
    ap.add_argument("--psd_clip", type=float, default=1e-14)
    ap.add_argument("--seed", type=int, default=0)

    ap.add_argument("--save_prefix", type=str, default="interface")
    ap.add_argument("--no_plots", action="store_true")

    args, unknown = ap.parse_known_args()
    if unknown:
        print(f"Ignoring unknown args: {unknown}")

    if args.only_omega0 and args.only_pm_omega:
        raise ValueError("Choose at most one of --only_omega0 or --only_pm_omega.")
    if args.only_omega0 and (not args.use_omega0):
        raise ValueError("--only_omega0 requires --use_omega0.")

    if not args.center and args.j0 is None:
        args.center = True
```

```

return args

def paulis():
    I = np.array([[1, 0], [0, 1]], dtype=complex)
    X = np.array([[0, 1], [1, 0]], dtype=complex)
    Y = np.array([[0, -1j], [1j, 0]], dtype=complex)
    Z = np.array([[1, 0], [0, -1]], dtype=complex)
    return I, X, Y, Z

def kron_all(ops):
    out = ops[0]
    for op in ops[1:]:
        out = np.kron(out, op)
    return out

def local_pauli(N, site, which):
    I, X, Y, Z = paulis()
    P = {"I": I, "X": X, "Y": Y, "Z": Z}[which.upper()]
    ops = [I] * N
    ops[site] = P
    return kron_all(ops)

def two_site_xx(N, i):
    I, X, Y, Z = paulis()
    ops = [I] * N
    ops[i] = X
    ops[i + 1] = X
    return kron_all(ops)

def tfim_hamiltonian(N, J, h):
    dim = 2**N
    H = np.zeros((dim, dim), dtype=complex)
    for i in range(N - 1):
        H += -J * two_site_xx(N, i)
    for i in range(N):
        H += -h * local_pauli(N, i, "Z")
    return (H + H.conj().T) / 2

def thermal_probs(E, beta):
    Emin = np.min(E.real)
    w = np.exp(-beta * (E.real - Emin))
    return w / np.sum(w)

def gamma_model(omega, beta, gamma0=0.2, s=1.0):
    if omega > 0:
        denom = 1.0 - np.exp(-beta * omega)
        denom = denom if denom > 1e-15 else 1e-15
        return float(gamma0 * (omega**s) / denom)
    if omega < 0:
        op = abs(omega)
        return float(np.exp(-beta * op) * gamma_model(op, beta, gamma0=gamma0, s=s))

```

```

return 0.0

def project_to_density(rho, eps=1e-14):
    rho = 0.5 * (rho + rho.conj().T)
    evals, evecs = la.eigh(rho)
    evals = np.maximum(evals.real, eps)
    rho_psd = (evecs * evals) @ evecs.conj().T
    rho_psd = 0.5 * (rho_psd + rho_psd.conj().T)
    rho_psd = rho_psd / np.trace(rho_psd)
    return rho_psd

def rel_entropy(rho, sigma):
    rho = 0.5 * (rho + rho.conj().T)
    sigma = 0.5 * (sigma + sigma.conj().T)
    log_rho = la.logm(rho)
    log_sig = la.logm(sigma)
    return float(np.trace(rho @ (log_rho - log_sig)).real)

def lindblad_L_of_rho(rho, L_ops, rates):
    out = np.zeros_like(rho)
    for Lk, rk in zip(L_ops, rates):
        if rk == 0:
            continue
        LdL = Lk.conj().T @ Lk
        out += rk * (Lk @ rho @ Lk.conj().T - 0.5 * (LdL @ rho + rho @ LdL))
    return out

def build_Somega_bin(S_E, E, omega_target, bin_tol):
    d = S_E.shape[0]
    Sbin = np.zeros((d, d), dtype=complex)
    for m in range(d):
        for n in range(d):
            omega = float((E[m] - E[n]).real)
            if abs(omega - omega_target) <= bin_tol:
                Sbin[m, n] = S_E[m, n]
    return Sbin

def choose_pair_pmomega(S_E, p, E, bin_tol):
    d = S_E.shape[0]
    best = (-1.0, None)
    for m in range(d):
        for n in range(d):
            if m == n:
                continue
            omega = float((E[m] - E[n]).real)
            if omega <= bin_tol:
                continue
            score = float(abs(S_E[m, n]) * np.sqrt(p[m] * p[n]))
            if score > best[0]:
                best = (score, (m, n, omega))
    if best[1] is None:
        raise RuntimeError("Could not find a positive-frequency pair.")
    return best[1]

```

```

def choose_pair_omega0(S_E, p, E, bin_tol):
    d = S_E.shape[0]
    sdiag = np.diag(S_E).real
    best = (-1.0, None)
    for m in range(d):
        for n in range(d):
            if m == n:
                continue
            omega = float((E[m] - E[n]).real)
            if abs(omega) <= bin_tol:
                continue
            score = float(abs(sdiag[m] - sdiag[n]) * np.sqrt(p[m] * p[n]))
            if score > best[0]:
                best = (score, (m, n, abs(omega)))
    if best[1] is None:
        raise RuntimeError("Could not find a pair for omega0 mode.")
    return best[1]

def make_rho_E(p, m, n, alpha):
    rho_E = np.diag(p.astype(complex))
    amp = alpha * np.sqrt(p[m] * p[n])
    rho_E[m, n] += amp
    rho_E[n, m] += amp
    return rho_E

def main():
    args = parse_args()
    np.random.seed(args.seed)

    N, J, h, beta = args.N, args.J, args.h, args.beta
    j0 = args.j0 if args.j0 is not None else (N // 2 if args.center else 0)

    H = tfim_hamiltonian(N, J, h)
    E, V = la.eigh(H)
    Vdag = V.conj().T
    p = thermal_probs(E, beta)

    S = local_pauli(N, j0, args.Sop)
    S_E = Vdag @ S @ V

    if args.pair_mode == "pmomega":
        m, n, omega = choose_pair_pmomega(S_E, p, E, args.bin_tol)
        print("Pair mode: pmomega")
    else:
        m, n, omega = choose_pair_omega0(S_E, p, E, args.bin_tol)
        print("Pair mode: omega0")

    print(f"Selected pair (m,n)={({m},{n}), omega≈{omega:.12g}")

    Somega_E = build_Somega_bin(S_E, E, omega, args.bin_tol)
    Sminus_E = Somega_E.conj().T
    Somega = V @ Somega_E @ Vdag
    Sminus = V @ Sminus_E @ Vdag

```

```

g_plus = gamma_model(omega, beta, gamma0=args.gamma0, s=args.s)
g_minus = gamma_model(-omega, beta, gamma0=args.gamma0, s=args.s)

omega0_ops, omega0_rates = [], []
pm_ops, pm_rates = [Somega, Sminus], [g_plus, g_minus]

if args.use_omega0:
    S0_E = np.diag(np.diag(S_E))
    S0 = V @ S0_E @ Vdag
    gamma0_val = float(args.gamma0 / beta)
    omega0_ops, omega0_rates = [S0], [gamma0_val]
    print(f"Included omega=0 channel with gamma(0)=gamma0/beta={gamma0_val:.6g}")

if args.only_omega0:
    L_ops, rates = omega0_ops, omega0_rates
    print("Mode: ONLY omega=0")
elif args.only_pm_omega:
    L_ops, rates = pm_ops, pm_rates
    print("Mode: ONLY ±omega")
else:
    L_ops, rates = omega0_ops + pm_ops, omega0_rates + pm_rates
    print("Mode: omega=0 (if enabled) PLUS ±omega")

alphas = np.linspace(1e-4, args.alpha_max, args.nsamples)
dt = args.dt

Cs, dCs, ratios = [], [], []
for alpha in alphas:
    rho_E = make_rho_E(p, m, n, alpha)
    rho = project_to_density(V @ rho_E @ Vdag, eps=args.psd_clip)

    rho_E_chk = Vdag @ rho @ V
    Delta_rho = project_to_density(V @ np.diag(np.diag(rho_E_chk)) @ Vdag, eps=args.psd_clip)
    C = rel_entropy(rho, Delta_rho)

    Lrho = lindblad_L_of_rho(rho, L_ops, rates)
    rho_t = project_to_density(rho + dt * Lrho, eps=args.psd_clip)

    rho_t_E = Vdag @ rho_t @ V
    Delta_rho_t = project_to_density(V @ np.diag(np.diag(rho_t_E)) @ Vdag, eps=args.psd_clip)
    C_t = rel_entropy(rho_t, Delta_rho_t)

    dotC_loss = -(C_t - C) / dt

    Cs.append(C)
    dCs.append(dotC_loss)
    ratios.append(dotC_loss / C if C > 0 else np.nan)

Cs = np.array(Cs, dtype=float)
dCs = np.array(dCs, dtype=float)
ratios = np.array(ratios, dtype=float)

alpha_last = float(alphas[-1])
r_last = float(ratios[-1])
C_last = float(Cs[-1])

```

```

print(f"RESULT alpha_last={alpha_last:.12g} C_last={C_last:.12g} r_last={r_last
      :.12g}")

if not args.no_plots:
    prefix = args.save_prefix

    plt.figure(figsize=(6.5, 4.2))
    plt.plot(Cs, ratios, marker="o", linewidth=1.2)
    plt.grid(True, alpha=0.3)
    plt.xlabel("Coherence  $C(\rho)=S(\rho)\Delta[\rho]$ ")
    plt.ylabel("Ratio  $\dot{C}_{\mathrm{loss}}(\rho)/C(\rho)$ ")
    plt.title("Interface test (low-RAM Euler step)")
    plt.tight_layout()
    plt.savefig(f"{prefix}_ratio_vs_C.png", dpi=200)

    plt.figure(figsize=(6.5, 4.2))
    plt.plot(Cs, dCs, marker="o", linewidth=1.2)
    plt.grid(True, alpha=0.3)
    plt.xlabel(" $C(\rho)$ ")
    plt.ylabel(" $\dot{C}_{\mathrm{loss}}(\rho)$ ")
    plt.title("Interface test: linearity check (low-RAM Euler)")
    plt.tight_layout()
    plt.savefig(f"{prefix}_dotC_vs_C.png", dpi=200)

    print(f"Saved -> {prefix}_ratio_vs_C.png")
    print(f"Saved -> {prefix}_dotC_vs_C.png")

if __name__ == "__main__":
    main()

```

## C.2 $\omega = 0$ -only fixed-pair sweep script (minimal, for the $\beta$ table)

```

# file: tfim_omega0_fixedpair.py
import argparse
import numpy as np
import scipy.linalg as la

def parse_args():
    ap = argparse.ArgumentParser()
    ap.add_argument("--N", type=int, default=6)
    ap.add_argument("--J", type=float, default=1.0)
    ap.add_argument("--h", type=float, default=1.5)
    ap.add_argument("--beta", type=float, default=1.0)
    ap.add_argument("--Sop", type=str, default="Z", choices=["X", "Y", "Z"])
    ap.add_argument("--center", action="store_true")
    ap.add_argument("--j0", type=int, default=None)

    ap.add_argument("--m", type=int, required=True)
    ap.add_argument("--n", type=int, required=True)

    ap.add_argument("--gamma0", type=float, default=0.10, help="Implements gamma(0)=
                    gamma0/beta")
    ap.add_argument("--dt", type=float, default=5e-5)
    ap.add_argument("--alpha_max", type=float, default=0.15)
    ap.add_argument("--nsamples", type=int, default=25)

```

```

ap.add_argument("--psd_clip", type=float, default=1e-14)
ap.add_argument("--seed", type=int, default=0)
return ap.parse_args()

def paulis():
    I = np.array([[1, 0], [0, 1]], dtype=complex)
    X = np.array([[0, 1], [1, 0]], dtype=complex)
    Y = np.array([[0, -1j], [1j, 0]], dtype=complex)
    Z = np.array([[1, 0], [0, -1]], dtype=complex)
    return I, X, Y, Z

def kron_all(ops):
    out = ops[0]
    for op in ops[1:]:
        out = np.kron(out, op)
    return out

def local_pauli(N, site, which):
    I, X, Y, Z = paulis()
    P = {"X": X, "Y": Y, "Z": Z}[which.upper()]
    ops = [I] * N
    ops[site] = P
    return kron_all(ops)

def two_site_xx(N, i):
    I, X, Y, Z = paulis()
    ops = [I] * N
    ops[i] = X
    ops[i + 1] = X
    return kron_all(ops)

def tfim_hamiltonian(N, J, h):
    dim = 2**N
    H = np.zeros((dim, dim), dtype=complex)
    for i in range(N - 1):
        H += -J * two_site_xx(N, i)
    for i in range(N):
        H += -h * local_pauli(N, i, "Z")
    return (H + H.conj().T) / 2

def thermal_probs(E, beta):
    Emin = np.min(E.real)
    w = np.exp(-beta * (E.real - Emin))
    return w / np.sum(w)

def project_to_density(rho, eps=1e-14):
    rho = 0.5 * (rho + rho.conj().T)
    evals, evecs = la.eigh(rho)
    evals = np.maximum(evals.real, eps)
    rho_psd = (evecs * evals) @ evecs.conj().T
    rho_psd = 0.5 * (rho_psd + rho_psd.conj().T)

```

```

rho_psd = rho_psd / np.trace(rho_psd)
return rho_psd

def rel_entropy(rho, sigma):
    rho = 0.5 * (rho + rho.conj().T)
    sigma = 0.5 * (sigma + sigma.conj().T)
    return float(np.trace(rho @ (la.logm(rho) - la.logm(sigma))))

def lindblad_omega0_L(rho, S0, gamma0_val):
    S0sq = S0 @ S0
    return gamma0_val * (S0 @ rho @ S0 - 0.5 * (S0sq @ rho + rho @ S0sq))

def main():
    args = parse_args()
    np.random.seed(args.seed)

    N, J, h, beta = args.N, args.J, args.h, args.beta
    j0 = args.j0 if args.j0 is not None else (N // 2 if args.center else 0)

    H = tfim_hamiltonian(N, J, h)
    E, V = la.eigh(H)
    Vdag = V.conj().T
    p = thermal_probs(E, beta)

    S = local_pauli(N, j0, args.Sop)
    S_E = Vdag @ S @ V

    S0_E = np.diag(np.diag(S_E))
    S0 = V @ S0_E @ Vdag

    m, n = args.m, args.n
    omega = abs(float((E[m] - E[n]).real))
    smsn = abs((S_E[m, m] - S_E[n, n]).real)
    gamma0_val = float(args.gamma0 / beta)

    print(f"Fixed pair (m,n)={m},{n}, omega≈{omega:.12g}")
    print(f"beta={beta}, gamma0={args.gamma0}, gamma(0)=gamma0/beta={gamma0_val}")
    print(f"p_m={p[m]:.6g}, p_n={p[n]:.6g}, |s_m-s_n|={smsn:.6g}")

    alphas = np.linspace(1e-4, args.alpha_max, args.nsamples)
    dt = args.dt

    diag_E = np.diag(p.astype(complex))

    print("\nalpha C(rho) dotC_loss dotC_loss/C")
    last_ratio = None
    for alpha in alphas:
        rho_E = diag_E.copy()
        amp = alpha * np.sqrt(p[m] * p[n])
        rho_E[m, n] += amp
        rho_E[n, m] += amp

        rho = project_to_density(V @ rho_E @ Vdag, eps=args.psd_clip)

        rho_E_chk = Vdag @ rho @ V

```

```

Delta_rho_E = np.diag(np.diag(rho_E_chk))
Delta_rho = project_to_density(V @ Delta_rho_E @ Vdag, eps=args.psd_clip)

C = rel_entropy(rho, Delta_rho)

rho_t = project_to_density(
    rho + dt * lindblad_omega0_L(rho, S0, gamma0_val),
    eps=args.psd_clip
)

rho_t_E = Vdag @ rho_t @ V
Delta_rho_t_E = np.diag(np.diag(rho_t_E))
Delta_rho_t = project_to_density(V @ Delta_rho_t_E @ Vdag, eps=args.psd_clip)

C_t = rel_entropy(rho_t, Delta_rho_t)
dotC_loss = -(C_t - C) / dt
ratio = dotC_loss / C if C > 0 else float("nan")
last_ratio = ratio
print(f"{alpha: .6f} {C: .6e} {dotC_loss: .6e} {ratio: .6g}")

print(f"\nRESULT r(alpha={args.alpha_max}) = {last_ratio}")

if __name__ == "__main__":
    main()

```

## References

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- [3] L. Eriksson, *The Heisenberg Cut as a Resource Boundary: An Operational Outlook from Coherence Maintenance Costs*, ai.vixra:2512.0064 (2025). <https://ai.vixra.org/abs/2512.0064>
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